

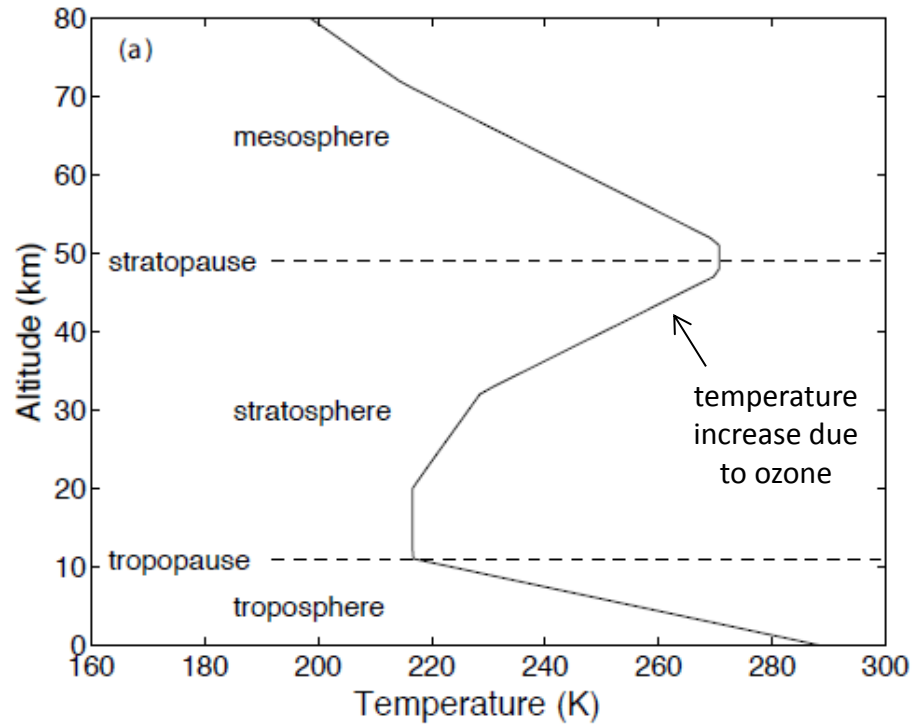
Overview of this week's lectures:

- Global circulation and transport
- Satellite observations of stratospheric temperature and water vapor
- Global upper troposphere - lower stratosphere (UTLS)
- Circulation and transport near the tropical tropopause layer (TTL)
- UTLS monsoon circulations and water vapor isotopes
- Research seminar: tropical dynamics with GPS radio occultation data

Lecture 1: Global atmospheric circulation and satellite observations

- large-scale circulation of the troposphere and stratosphere
 - Climatology and variability of the stratosphere
 - wave forcing of the zonal mean flow
 - Rossby waves: propagation and dissipation (critical layers)
 - Tropospheric baroclinic wave life cycles
 - large-scale tropical circulations
 - QBO and ENSO
- Large-scale transport

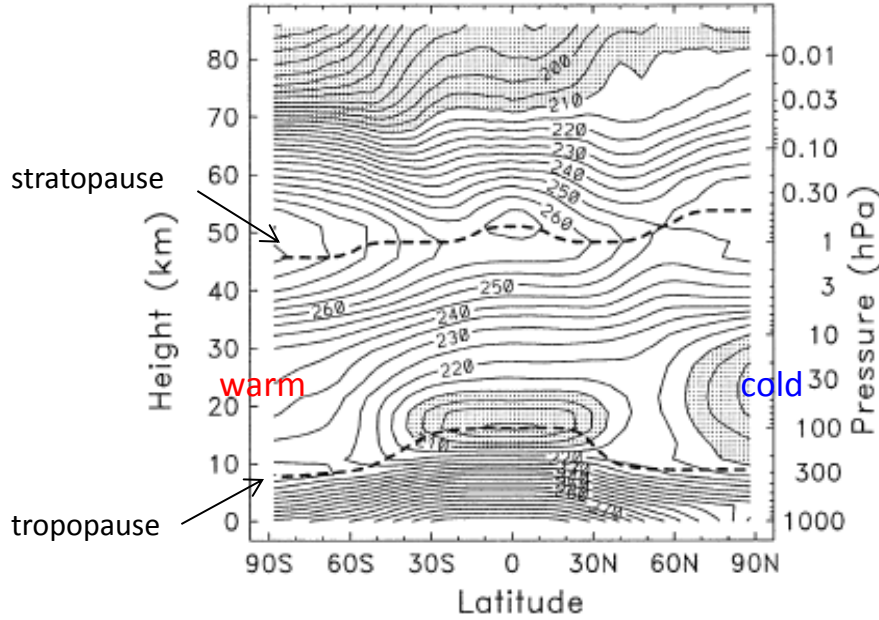
The Standard Atmosphere



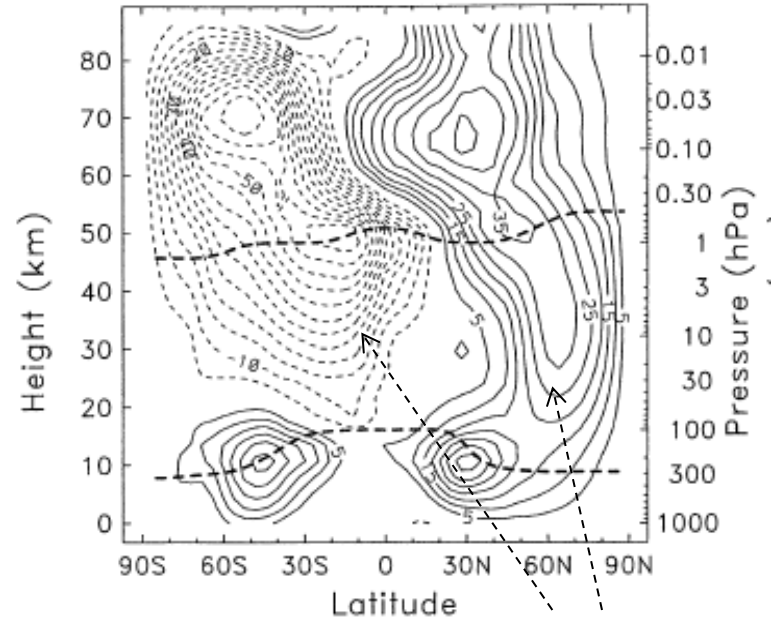
Climatological temperatures and zonal winds in January



January temperature

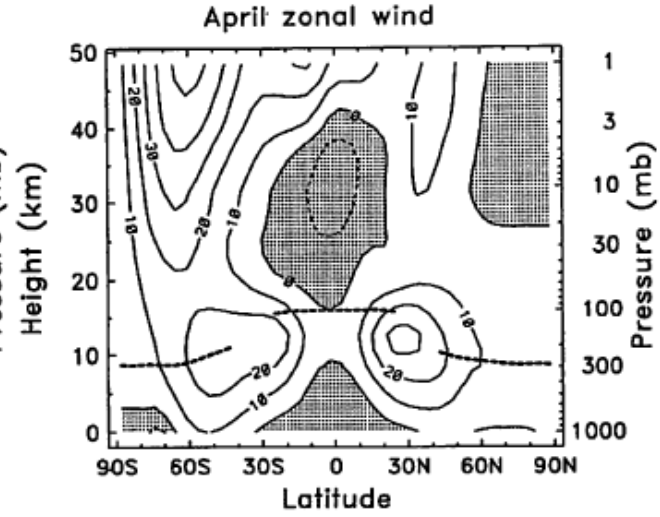
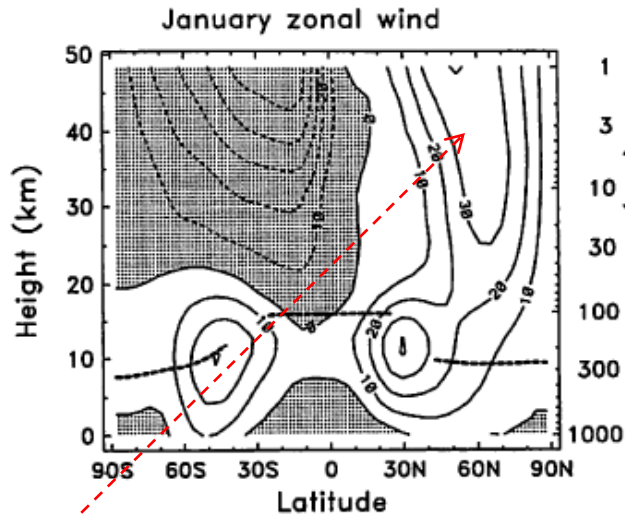


January zonal wind

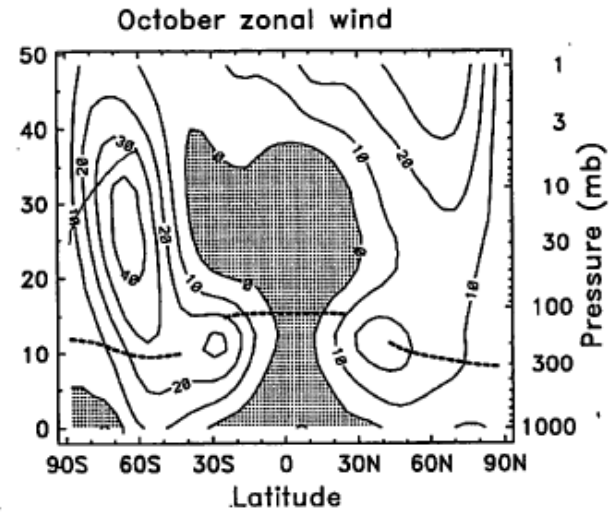
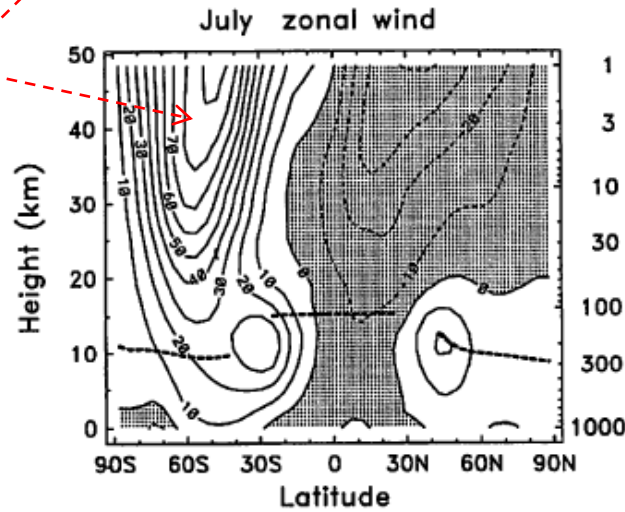


thermal wind balance

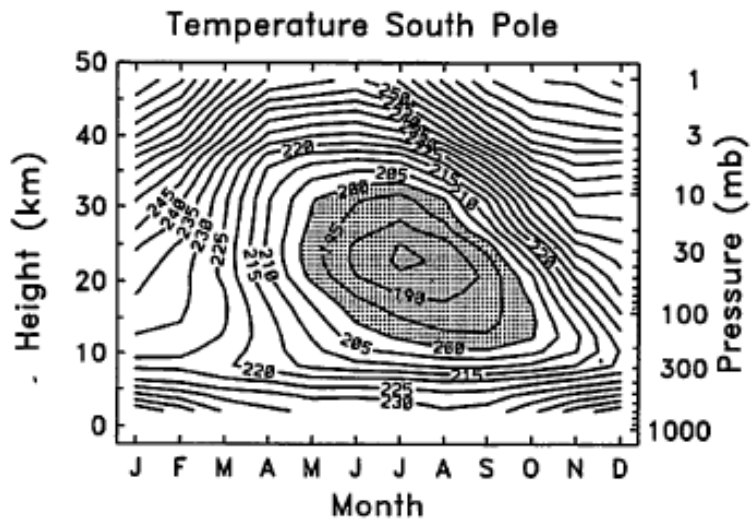
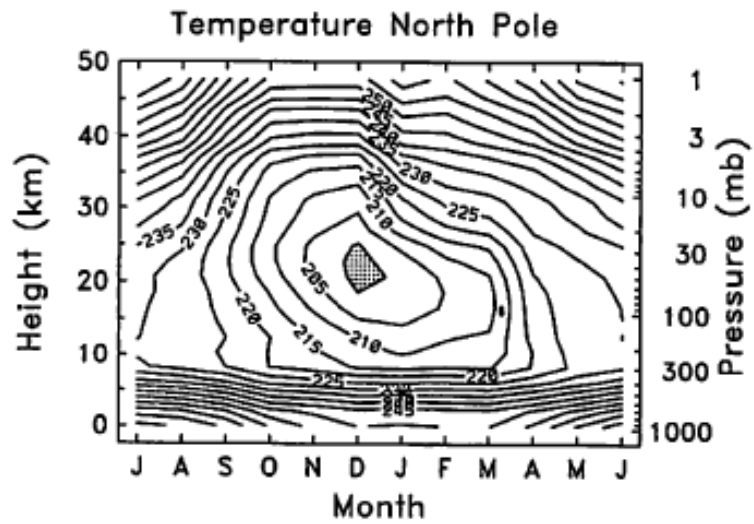
Seasonal cycle of zonal mean zonal winds



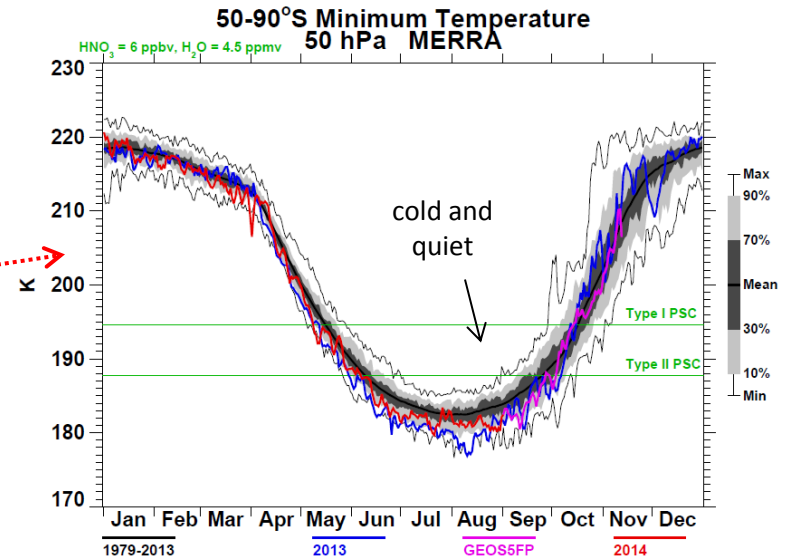
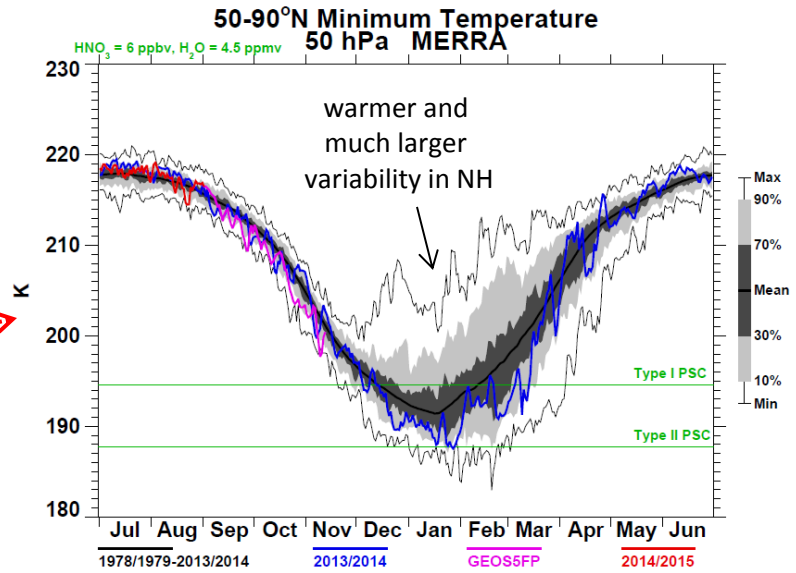
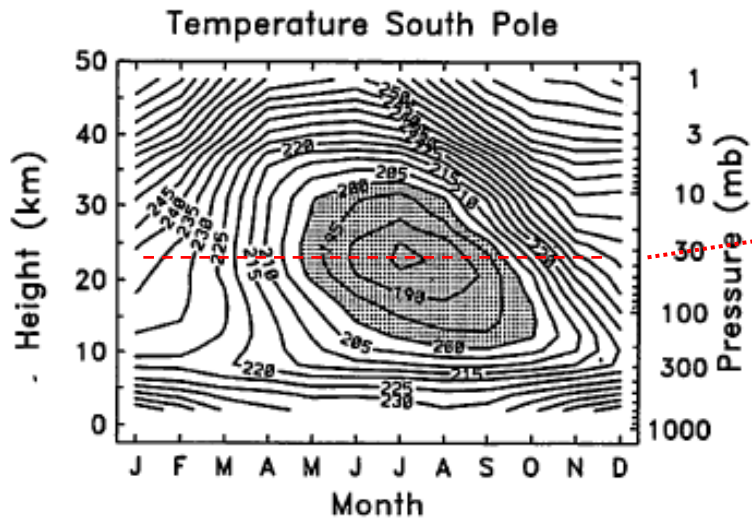
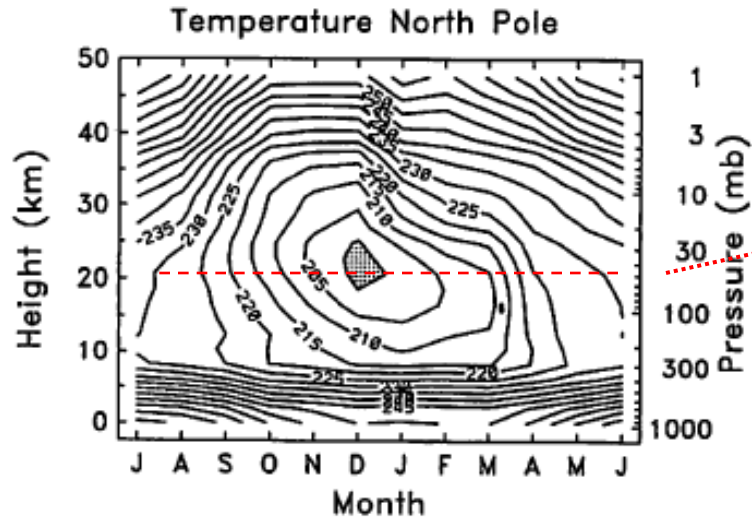
polar night jet
stronger in SH
than in NH

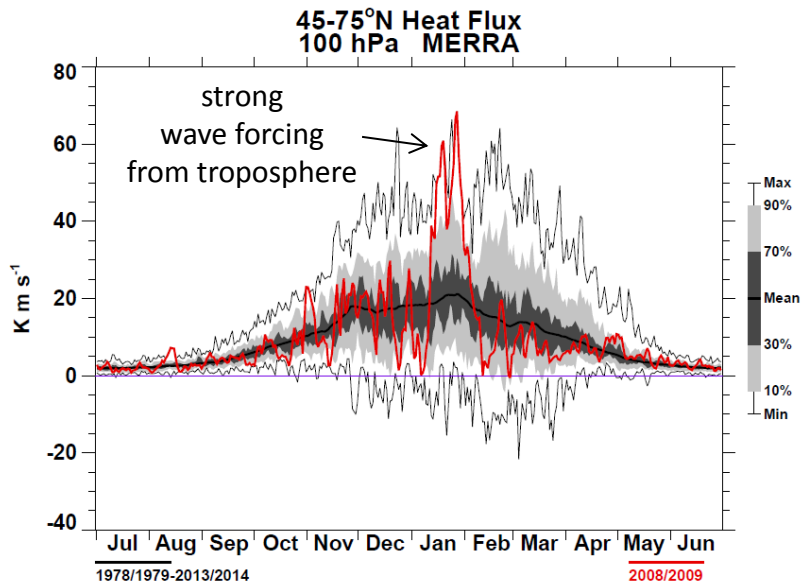
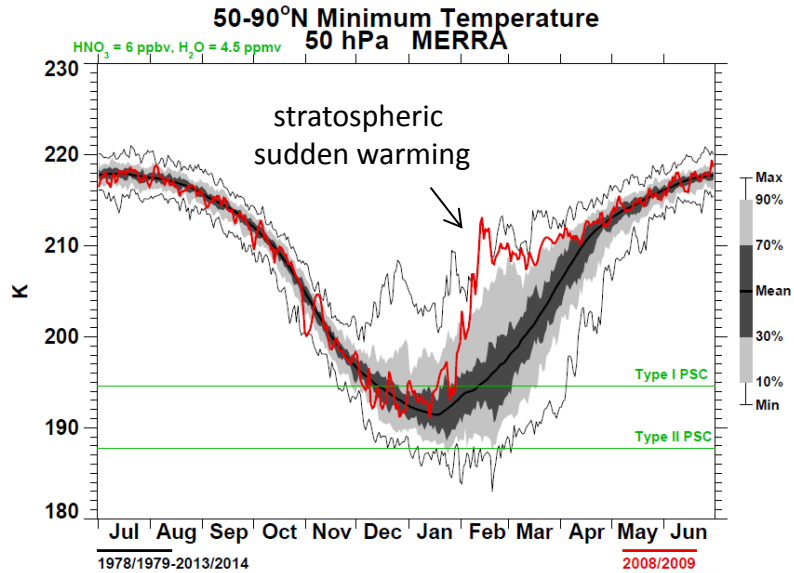


climatological polar temperature

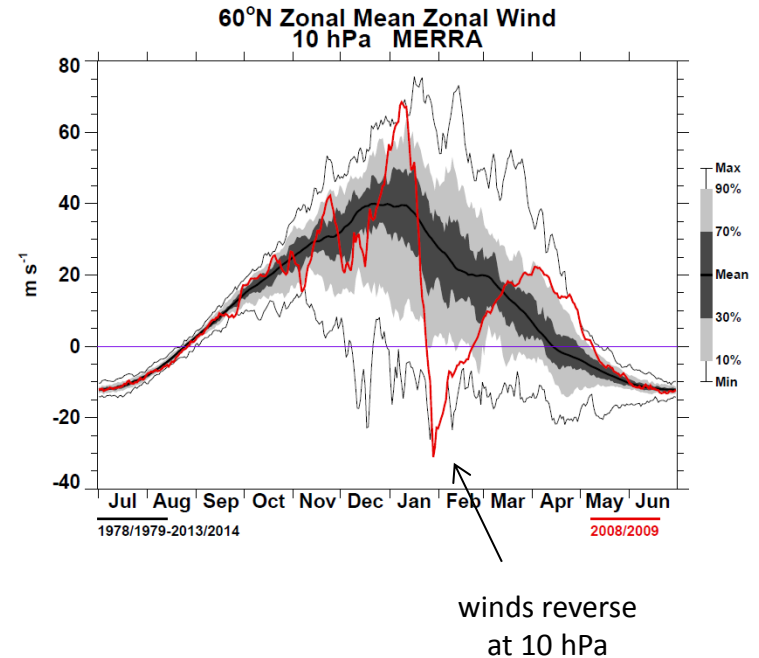


climatological polar temperature





- Variability in NH winter stratosphere tied to large-scale forcing from troposphere.
- Episodic forcing produces 'stratospheric sudden warming' events.
- Largest observed stratosphere sudden warming in January 2009



A Major Stratospheric Sudden Warming Event in January 2009

YAYOI HARADA, ATSUSHI GOTO, HIROSHI HASEGAWA, AND NORIHISA FUJIKAWA

Climate Prediction Division, Japan Meteorological Agency, Tokyo, Japan

JAS 2010

EP flux



(I) 18-20Jan. 2009, U & EPF

(Z at 10 hPa)

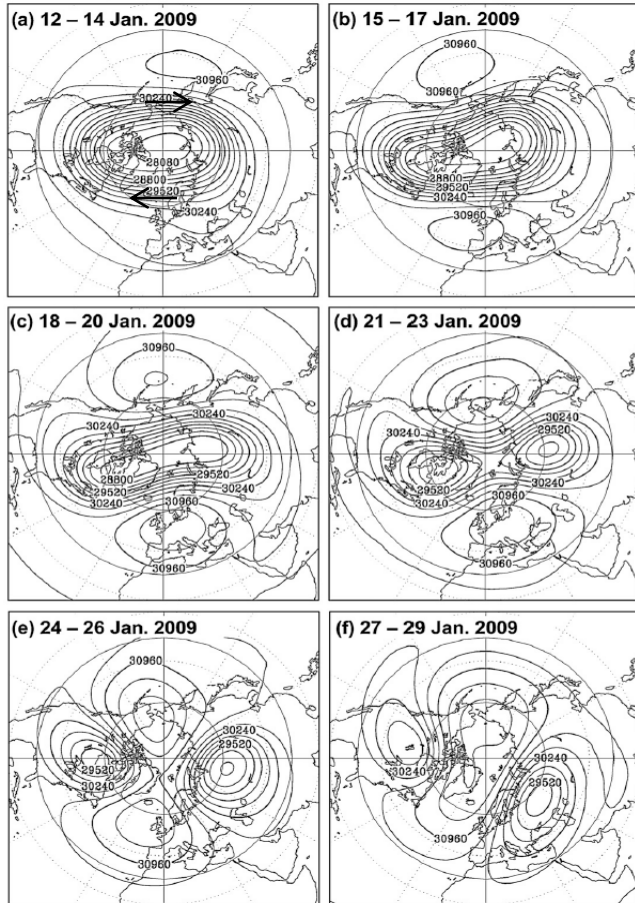
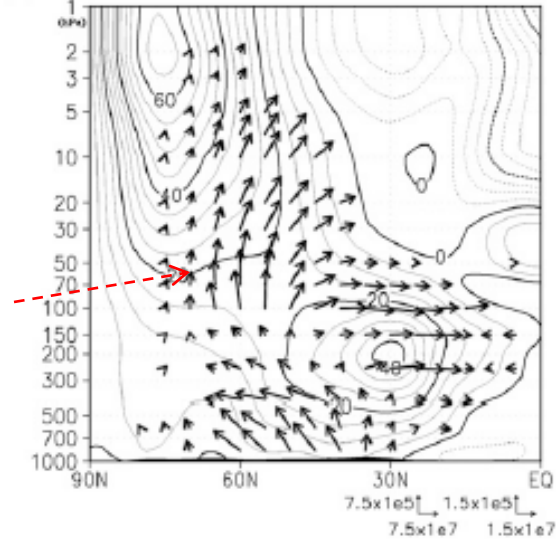


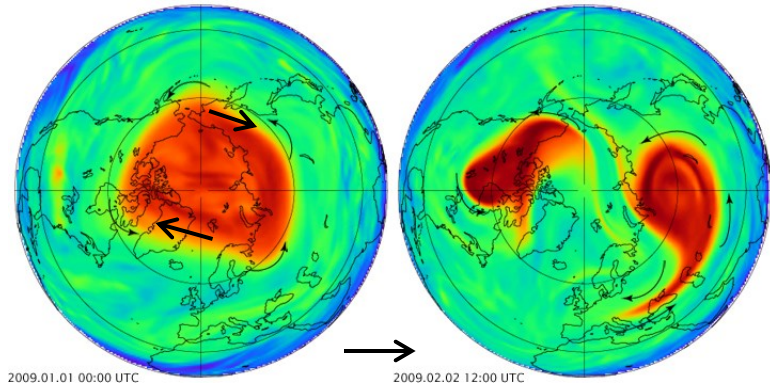
FIG. 3. The 10-hPa geopotential heights for six successive 3-day means in January 2009. The contour interval is 240 m.

wave forcing from troposphere



polar vortex near 30 km

potential vorticity



split of polar vortex

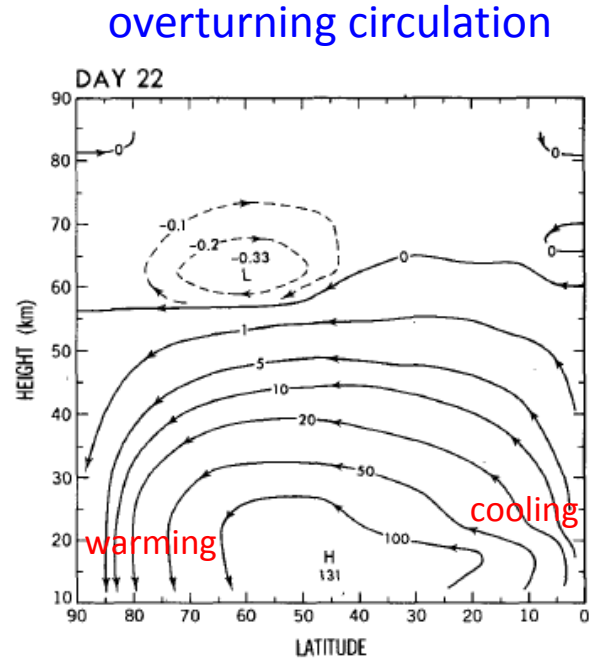
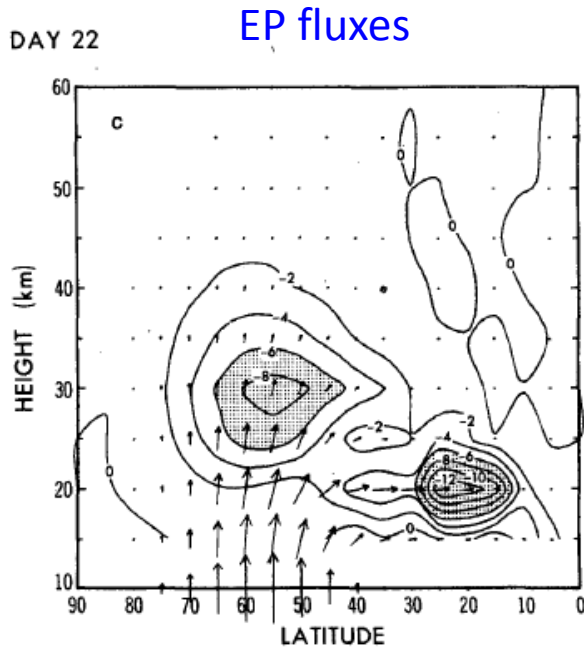
A Dynamical Model of the Stratospheric Sudden Warming

TAROH MATSUNO¹

Geophysical Fluid Dynamics Laboratory, NOAA, Princeton University, Princeton, N. J.

(Manuscript received 29 March 1971, in revised form 16 August 1971)

← solution to puzzle of stratospheric warmings



Some Eulerian and Lagrangian Diagnostics for a Model Stratospheric Warming¹

T. DUNKERTON, C.-P. F. HSU² AND M. E. MCINTYRE³

Department of Atmospheric Sciences, University of Washington, Seattle 98195

(Manuscript received 30 May 1980, in final form 11 December 1980)

Governing equations for the zonal mean flow (Transformed Eulerian mean)

EP flux divergence (wave forcing)

zonal momentum balance

$$\frac{\partial \bar{u}}{\partial t} - \hat{f} \bar{v}^* = \text{DF}$$

thermodynamic balance

$$\frac{\partial \bar{T}}{\partial t} + \bar{v}^* \frac{1}{a} \frac{\partial \bar{T}}{\partial \phi} + \bar{w}^* S = \bar{Q},$$

diabatic forcing

continuity equation

$$(a \cos \phi)^{-1} \frac{\partial}{\partial \phi} (\bar{v}^* \cos \phi) + e^{z/H} \frac{\partial}{\partial z} (\bar{w}^* e^{-z/H}) = 0,$$

geostrophic thermal wind

$$f \frac{\partial \bar{u}}{\partial z} + \frac{R}{aH} \frac{\partial \bar{T}}{\partial \phi} = 0.$$

Andrews et al, 1987

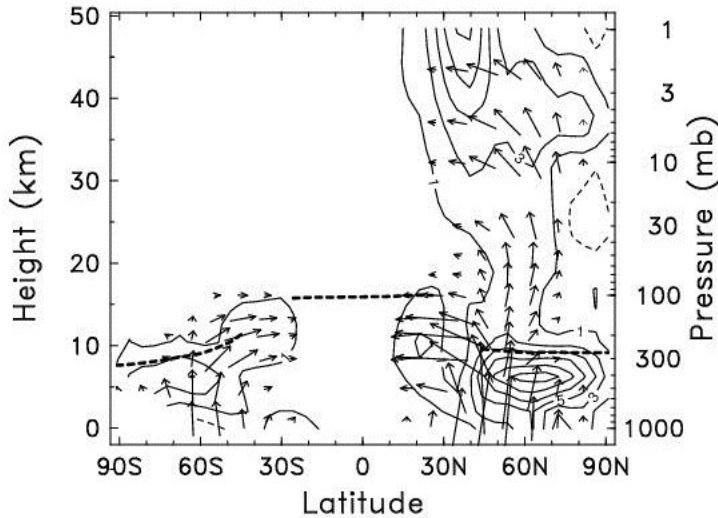
Eliassen-Palm fluxes:

EP flux divergence (wave forcing)

$$\frac{\partial \bar{u}}{\partial t} - \hat{f} \bar{v}^* = DF \quad DF = \frac{\exp(z/H)}{a \cos \phi} \nabla \cdot \mathbf{F}$$

climatology

January EP flux



components:

latitudinal flux

$$F_{\phi} = \exp(-z/H) a \cos \phi \left[-\overline{u'v'} \right]$$

momentum flux

vertical flux

$$F_z = \exp(-z/H) a \cos \phi \left[\overline{v'T'} \right]$$

heat flux

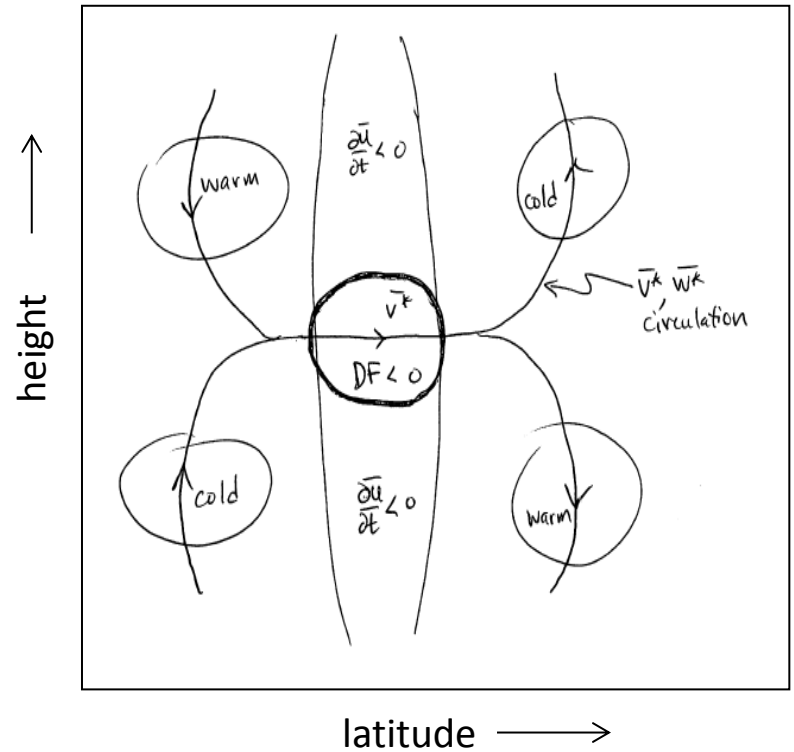
Important points:

- DF quantifies zonal momentum forcing
- \mathbf{F} proportional to 'wave activity' flux (DF shows sources and sinks of waves)
- F_{ϕ} and F_z indicate direction of wave propagation

Response of a balanced vortex to localized EP flux forcing (DF)

$$\frac{\partial \bar{u}}{\partial t} - f \bar{v}^* = DF$$

- response is balanced between $\frac{\partial \bar{u}}{\partial t}$ and $f \bar{v}^*$
- \bar{v}^* , \bar{w}^* and $\frac{\partial \bar{u}}{\partial t}$ act to extend DF forcing non-locally
- overall circulation maintains thermal wind balance



Circulation response depends on frequency of forcing:

Combine equations:

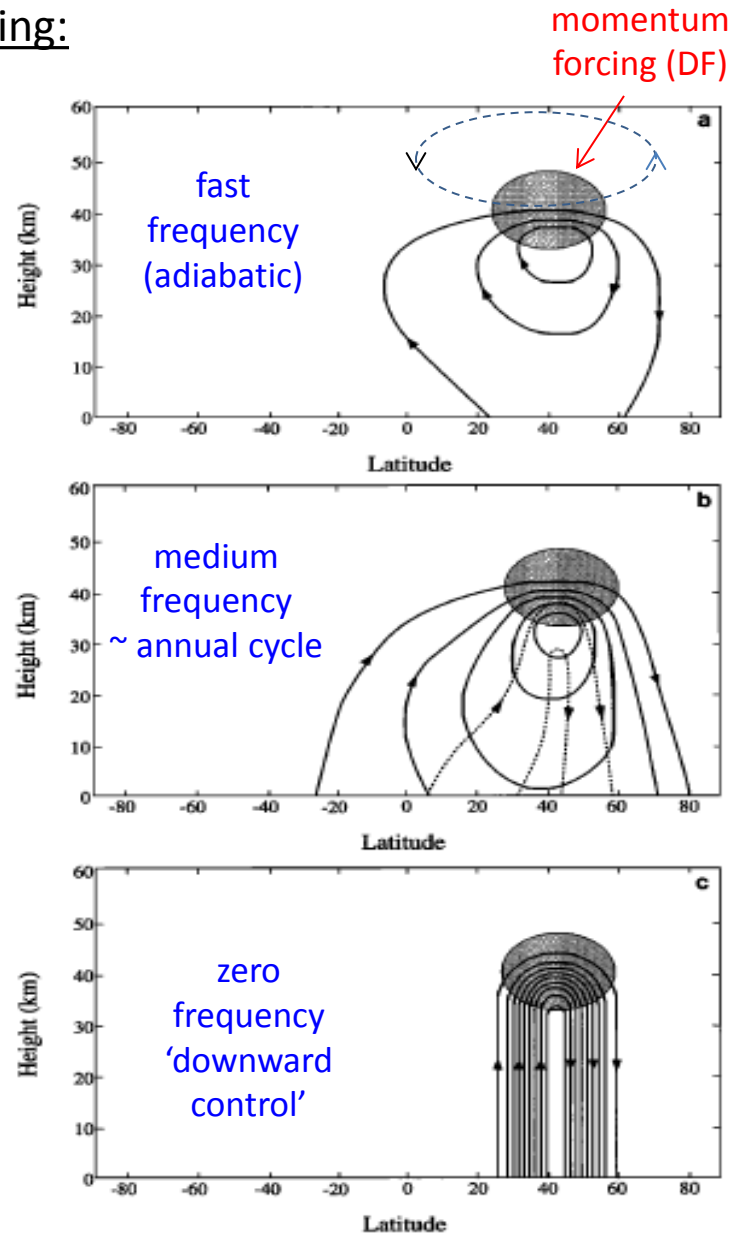
$$\frac{\partial}{\partial z} \left(\frac{1}{\rho_0} \frac{\partial(\rho_0 \hat{w})}{\partial z} \right) \quad \text{time dependence}$$

$$+ \left(\frac{i\sigma}{i\sigma + \alpha} \right) \frac{N^2}{4\Omega^2 a^2 \cos \phi} \frac{\partial}{\partial \phi} \left(\frac{\cos \phi}{\sin^2 \phi} \frac{\partial \hat{w}}{\partial \phi} \right)$$

$$= \left(\frac{i\sigma}{i\sigma + \alpha} \right) \frac{(R/H)}{4\Omega^2 a^2 \cos \phi} \frac{\partial}{\partial \phi} \left(\frac{\cos \phi}{\sin^2 \phi} \frac{\partial \hat{Q}}{\partial \phi} \right) \quad \leftarrow \text{diabatic heating}$$

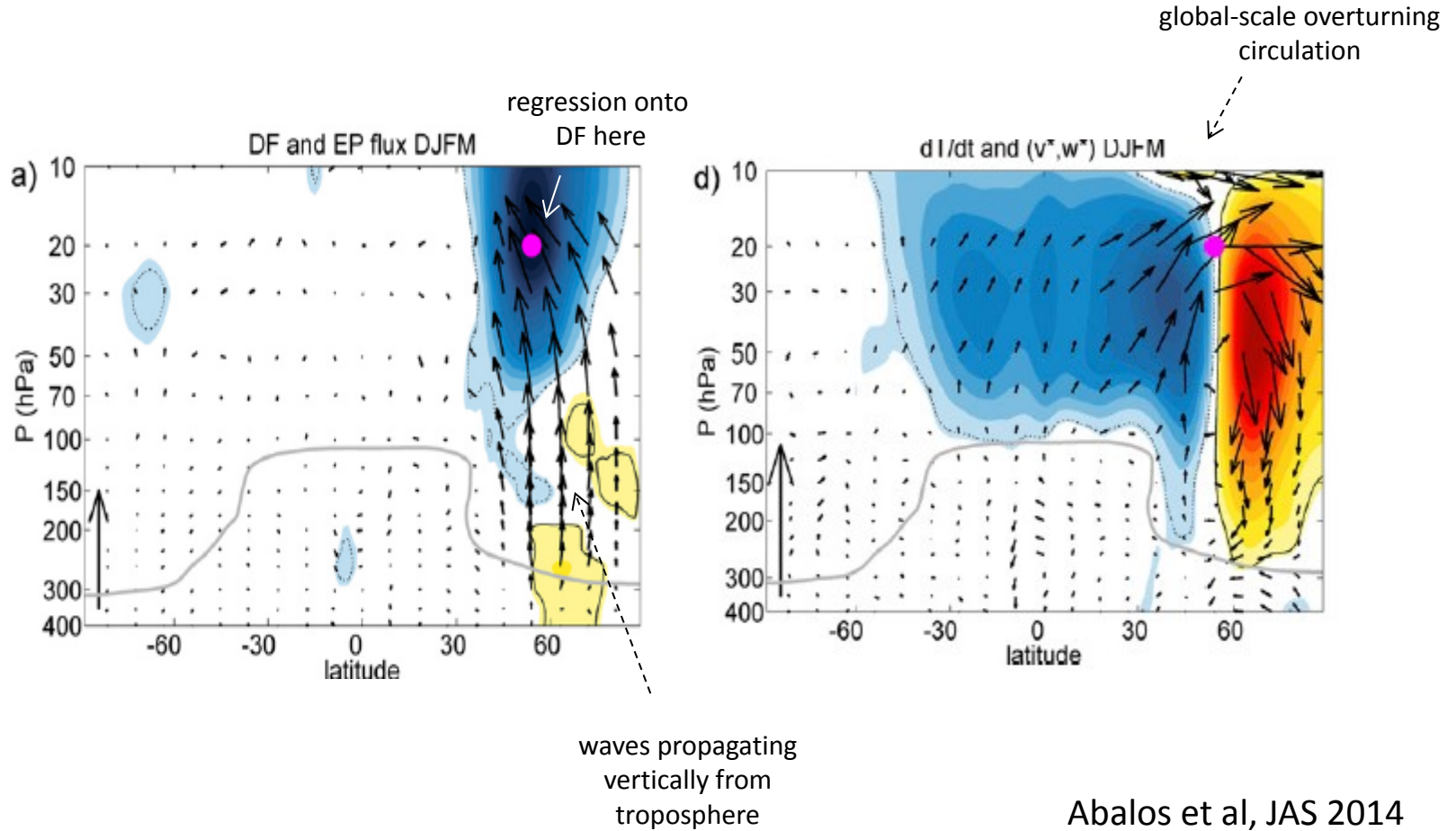
$$+ \frac{1}{2\Omega a \cos \phi} \frac{\partial}{\partial \phi} \left(\frac{\cos \phi}{\sin \phi} \frac{\partial \hat{G}}{\partial z} \right) \quad \leftarrow \text{momentum forcing (DF)}$$

In general both Q and DF drive the mean circulation. These plots show the response to isolated forcing from Rossby wave EP flux divergence. The lower cell becomes more important for slower forcing.

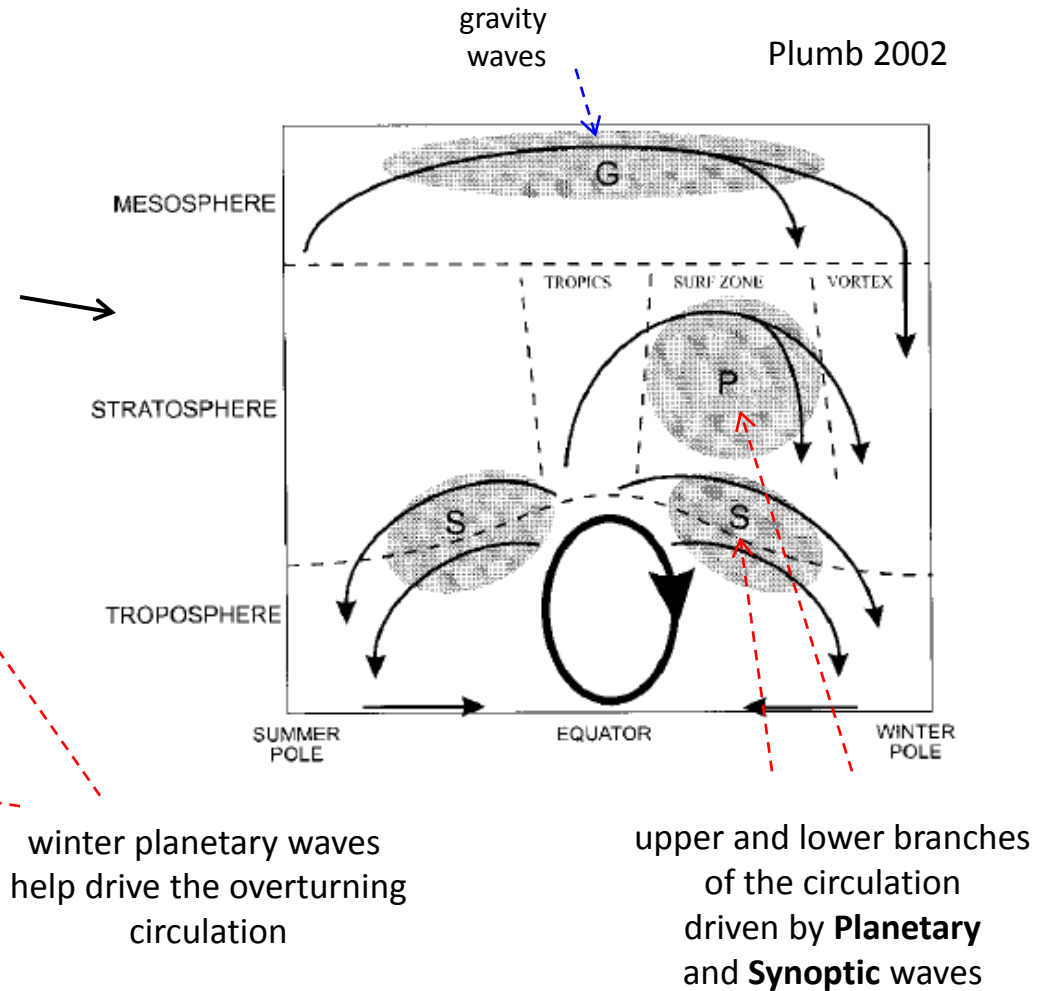
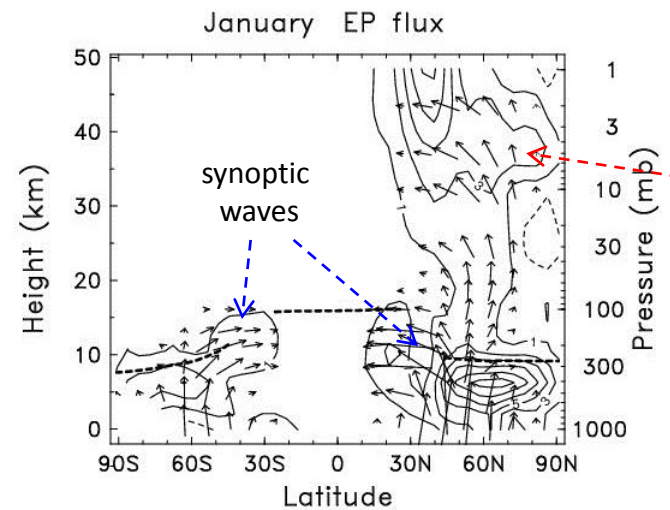
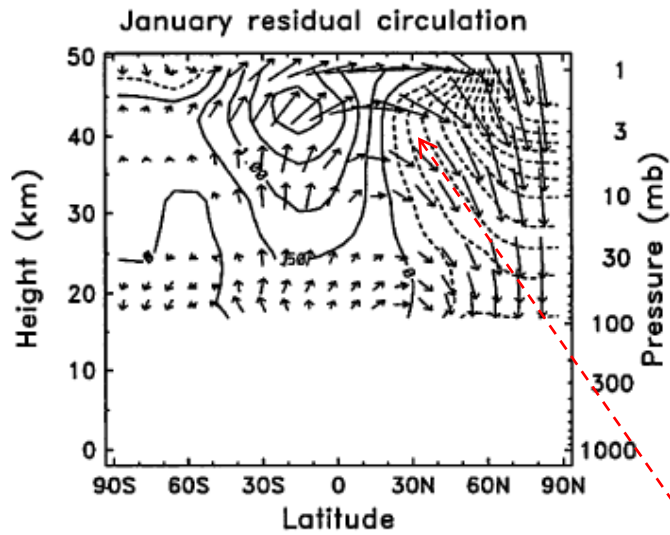


Haynes et al 1991
Holton et al 1995

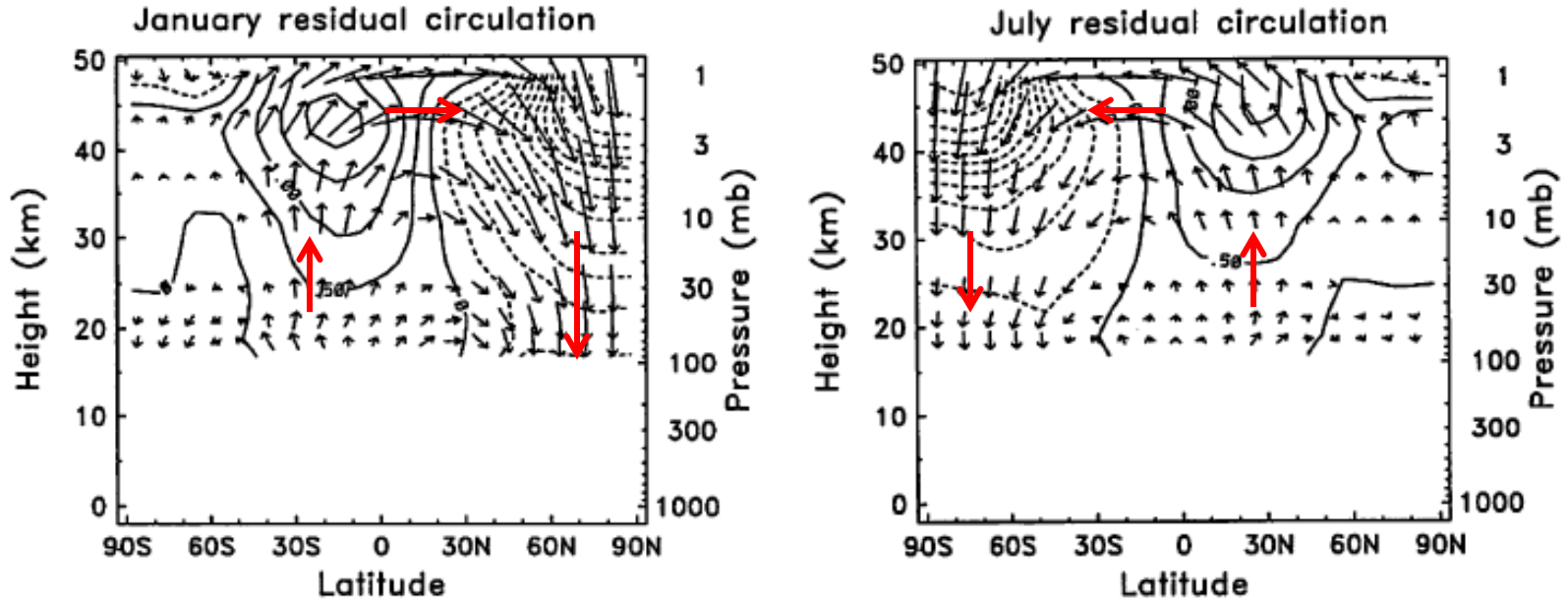
Wave – mean flow patterns from observations (reanalysis data)



Climatology of stratospheric overturning circulation



The overturning circulation reverses between solstice seasons



circulation is stronger during NH winter, related to stronger wave forcing from troposphere

The stratospheric overturning circulation is often termed the Brewer-Dobson circulation (closely related to the Lagrangian or transport circulation)

deduced by Brewer (1949) studying stratospheric water vapor and Dobson (1956) studying stratospheric ozone

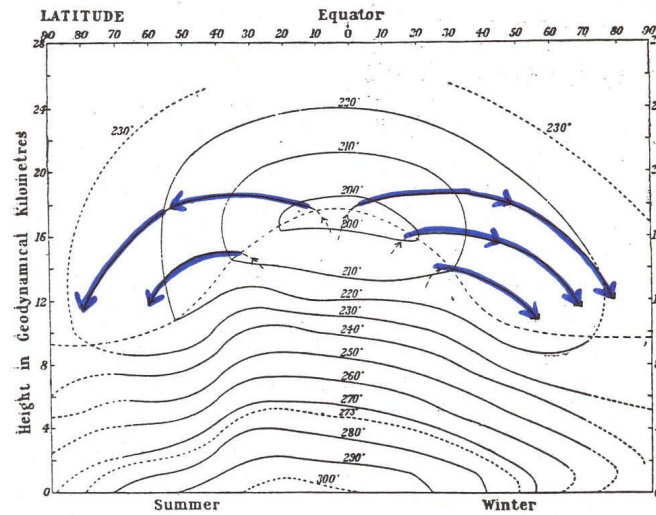


FIG. 5. Isotherms over the Globe
A supply of dry air is maintained by a slow mean circulation from the equatorial tropopause.

see recent review by Butchart 2014

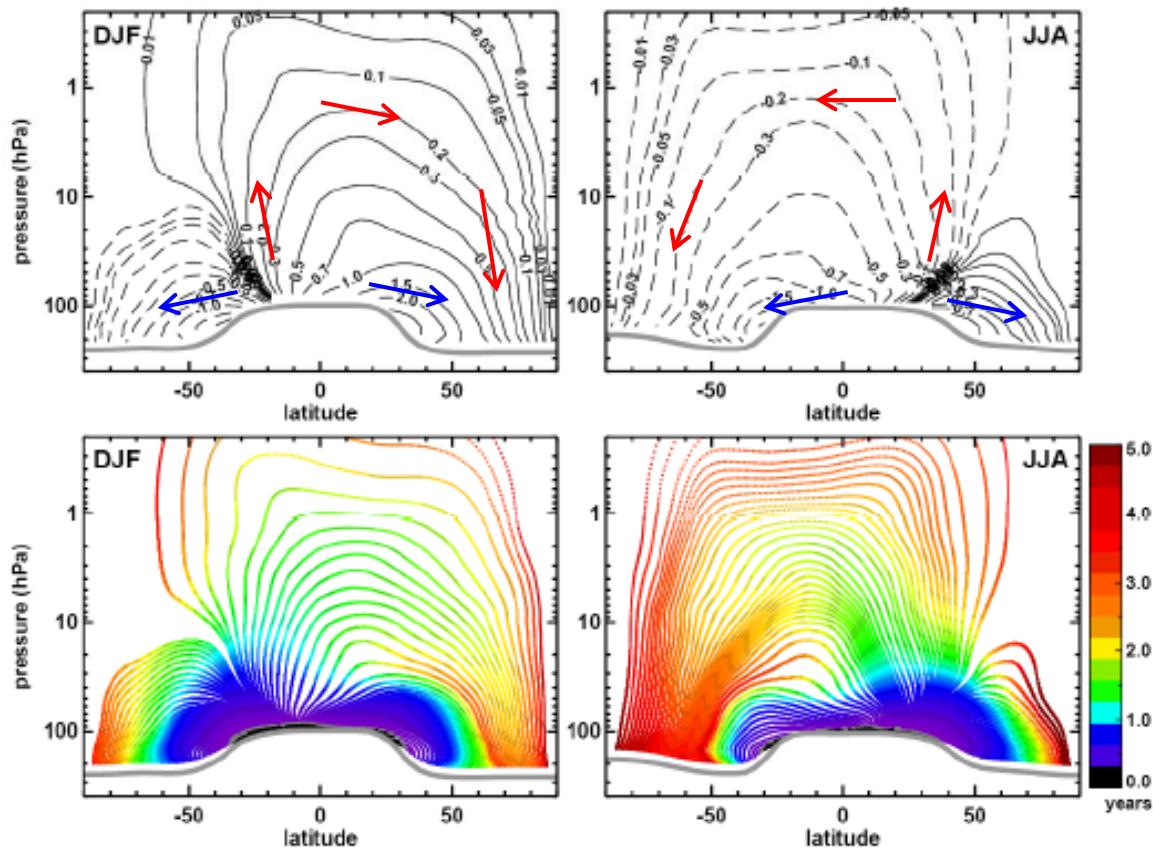
Residual circulation trajectories and transit times into the extratropical lowermost stratosphere

T. Birner¹ and H. Bönisch²

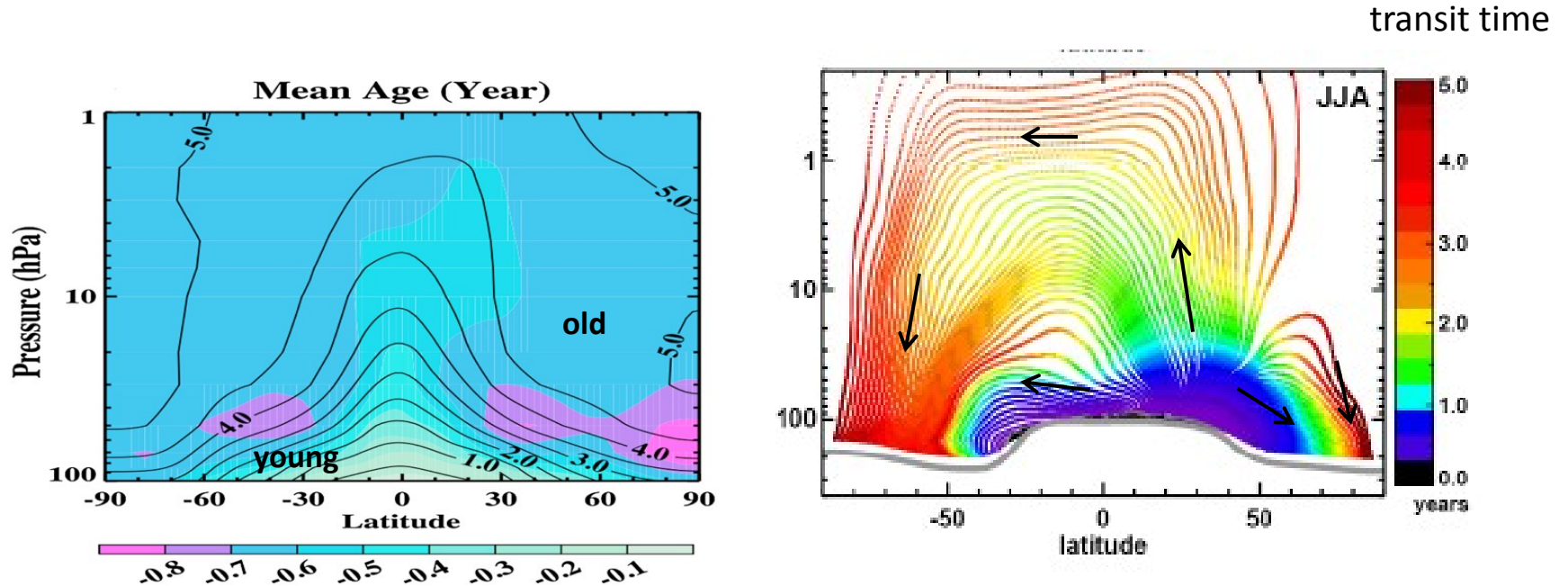
ACP 2011

renewed appreciation that there are upper and lower branches of the BDC

Red: deep branch (slow)
Blue: shallow branch (fast)



Transit time is closely related to 'mean age' (time since air entered stratosphere)

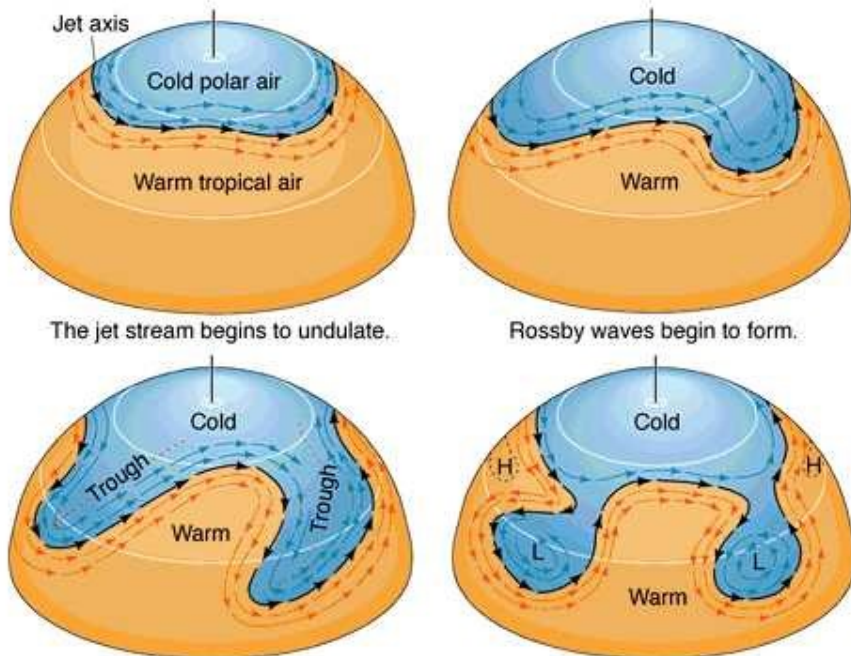


Air at any particular location is characterized by a distribution of transit times and ages (so-called age spectrum)

Key points:

- Asymmetry in winter stratosphere circulations: more disturbed in the NH, cold and quiet in the SH
- The stratosphere is forced by waves from the troposphere (stronger forcing in NH; episodic stratospheric sudden warmings)
- Dynamical response of balanced vortex to wave forcing (non-local temperature and wind changes)
- Eliassen-Palm (EP) fluxes quantify wave forcing
- Brewer-Dobson transport circulation (deep and shallow branches)

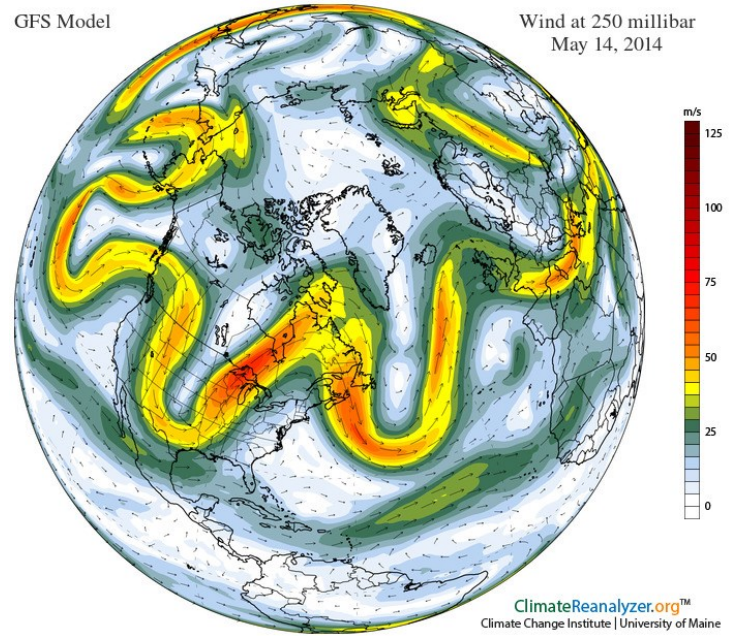
Rossby waves



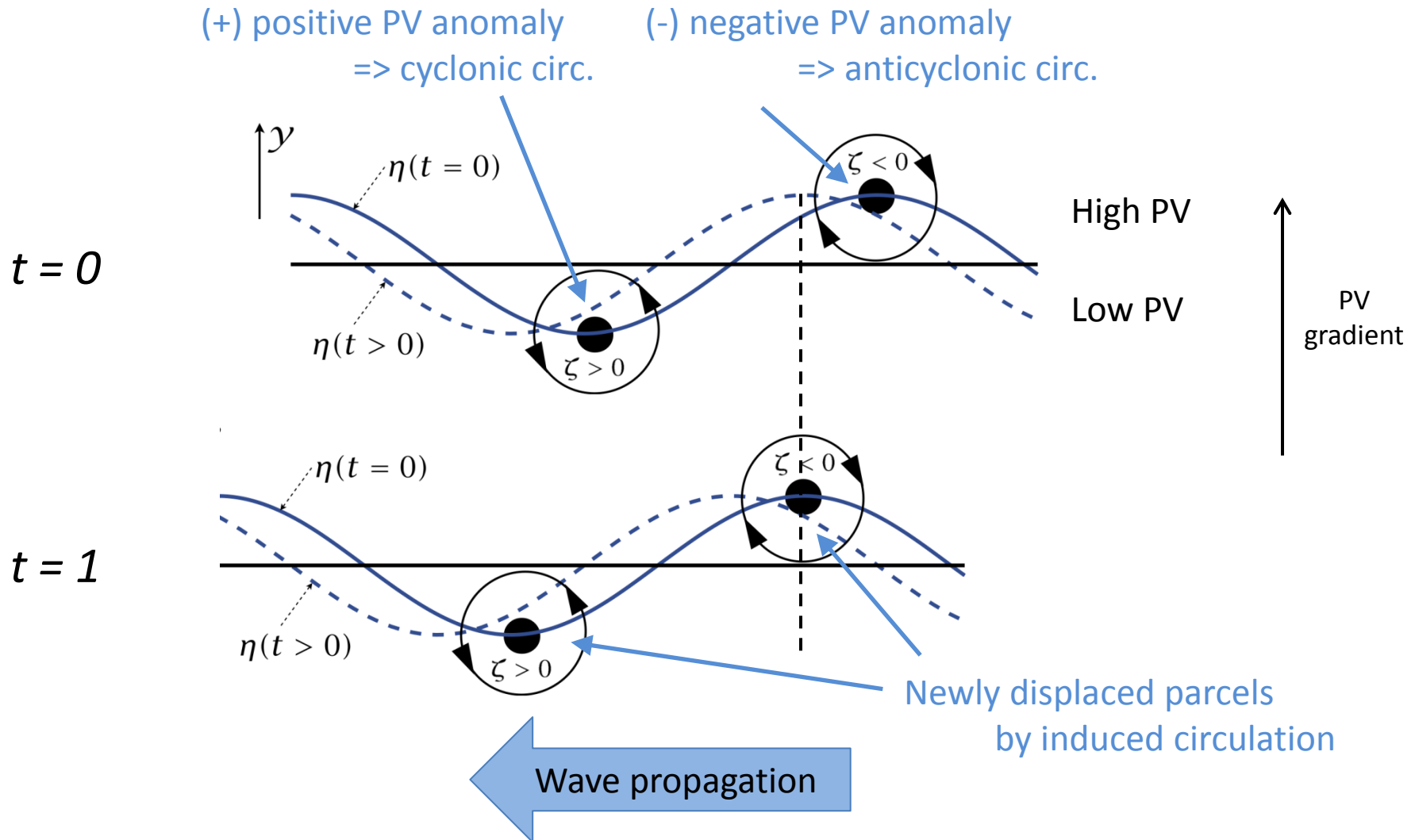
Copyright © A.N. Strahler.

GFS Model

Wind at 250 millibar
May 14, 2014



Rossby wave (potential vorticity wave)



Rossby wave propagation: quasi-geostrophic linearized PV equation

$$\left(\frac{\partial}{\partial t} + \frac{\bar{u}}{a \cos \phi} \frac{\partial}{\partial \lambda} \right) q'_{(M)} + a^{-1} \bar{q}_\phi v' = 0,$$

↑
eddy PV

↑
background PV gradient

wave solution:

$$\Phi' = e^{z/2H} \operatorname{Re} \Psi(\phi, z) e^{is\lambda}$$

$$\bar{q}_\phi = 2\Omega \cos \phi - \left[\frac{(\bar{u} \cos \phi)_\phi}{a \cos \phi} \right]_\phi - \frac{a}{\rho_0} \left(\frac{\rho_0 f^2}{N^2} \bar{u}_z \right)_z.$$

$$\frac{f^2}{a^2 \cos \phi} \left(\frac{\cos \phi}{f^2} \Psi_\phi \right)_\phi + \frac{f^2}{N^2} \Psi_{zz} + n_s^2 \Psi = 0$$

wave equation:
propagation for
 $n_s^2 > 0$

$$n_s^2 = \frac{\bar{q}_\phi}{a\bar{u}} - \frac{s^2}{a^2 \cos^2 \phi} - \frac{f^2}{4N^2 H^2}$$

refractive index

Propagation of Planetary-Scale Disturbances from the Lower into the Upper Atmosphere

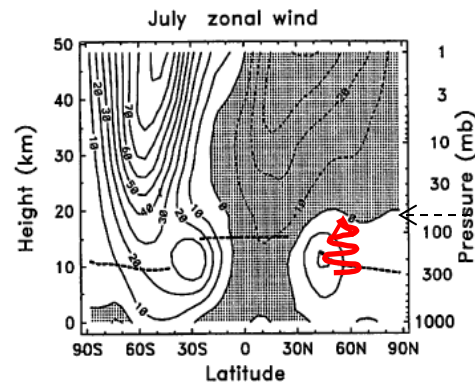
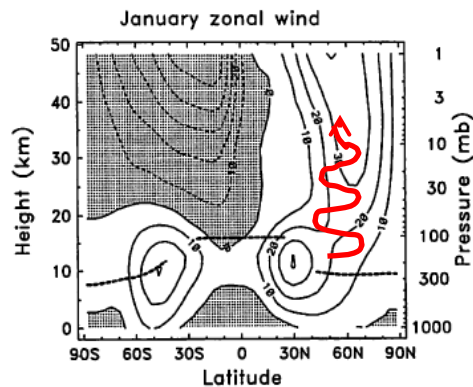
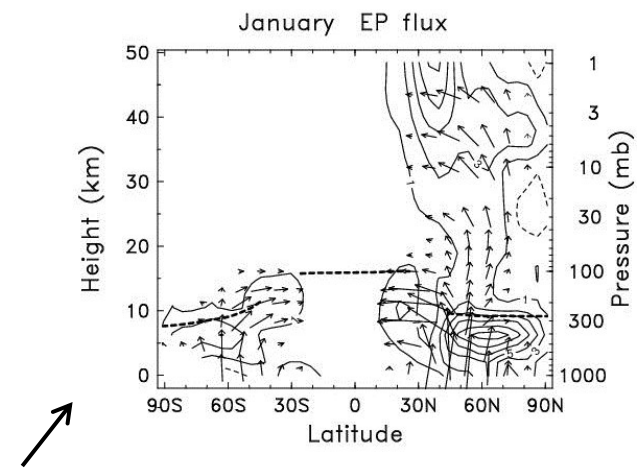
J. G. CHARNEY AND P. G. DRAZIN¹

JGR 1961

$$n_s^2 = \frac{\bar{q}_\phi}{a\bar{u}} \sim \cos(\text{lat}) + U \text{ terms} - \frac{s^2}{a^2 \cos^2 \phi} - \frac{f^2}{4N^2 H^2}$$

2 key points:

- n_s^2 proportional to $\sim \cos(\text{lat})$ (Rossby wave refraction towards low latitudes)
- vertical propagation for $U > 0$ and small zonal wavenumbers (planetary waves propagate to stratosphere only during winter)

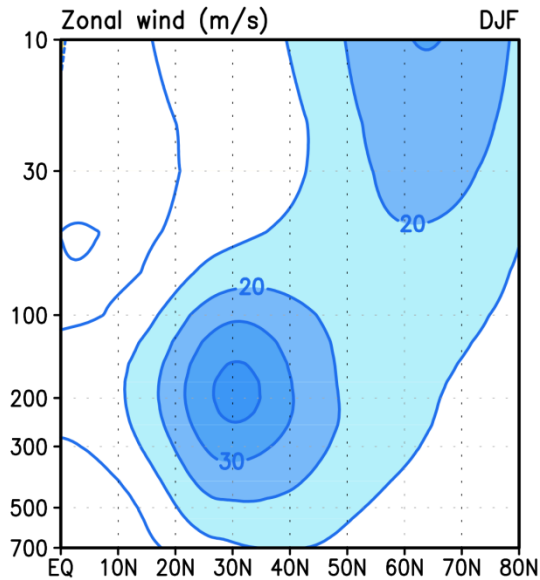


no propagation in summer easterlies

Refractive index squared (Matsuno 1970)

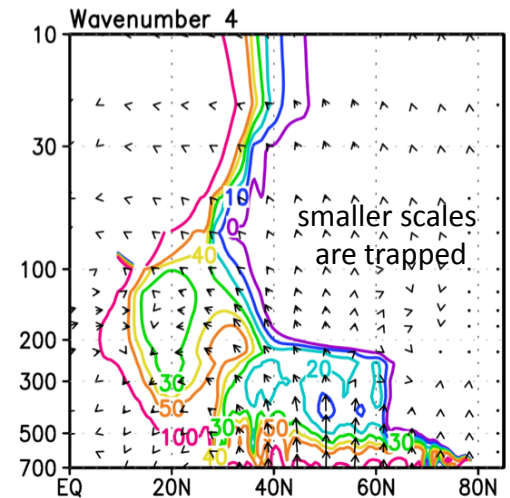
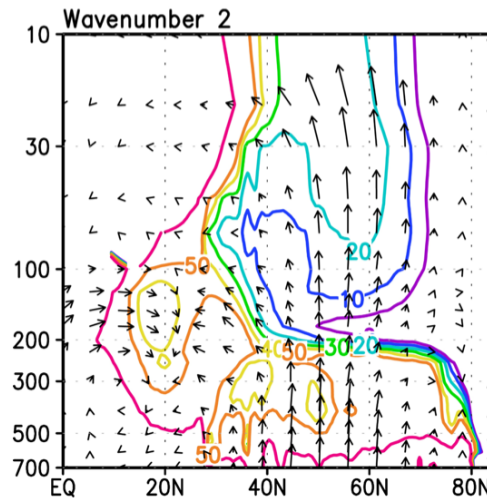
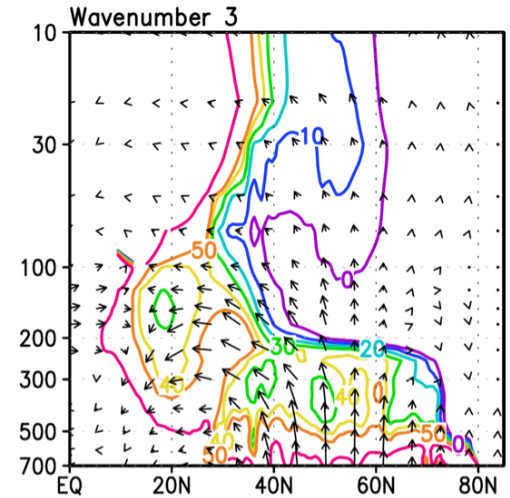
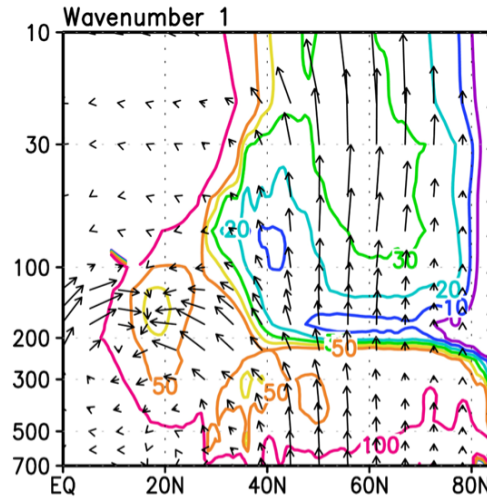
calculated n_s^2 and
observed EP fluxes

$$n_s^2 = \frac{\bar{q}_\phi}{a\bar{u}} - \frac{f^2}{4H^2N^2} - \frac{s^2}{a^2 \cos^2 \phi}$$

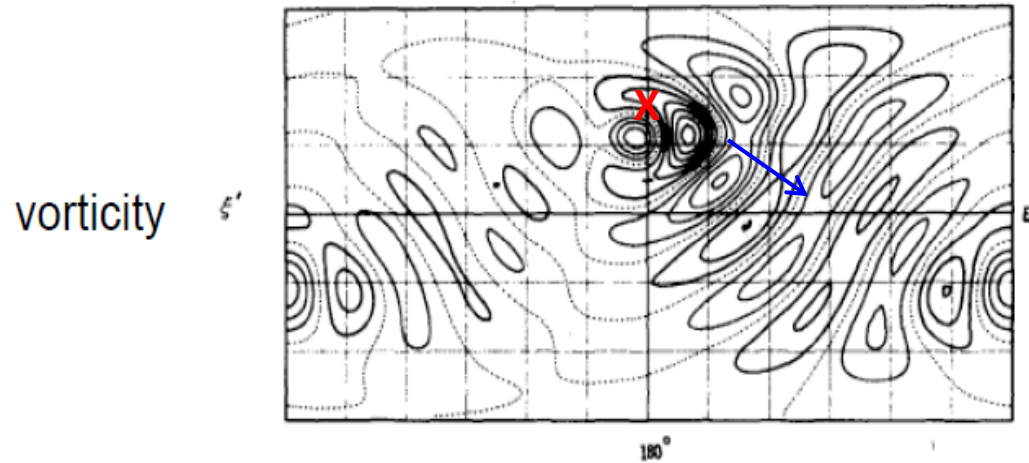


thanks to Joowan Kim

zonal wave 1 and 2 can propagate
vertically in climatological basic state

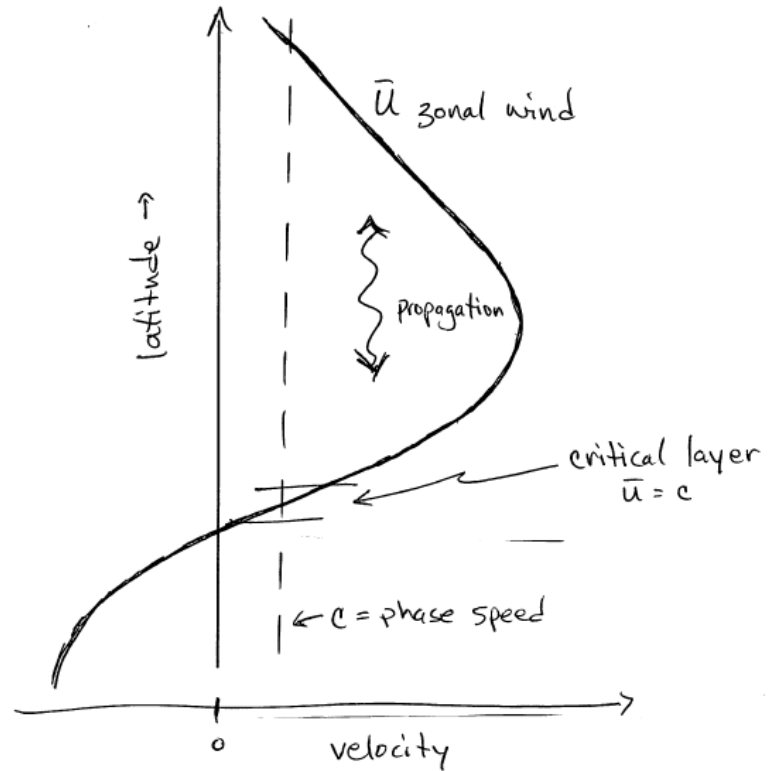
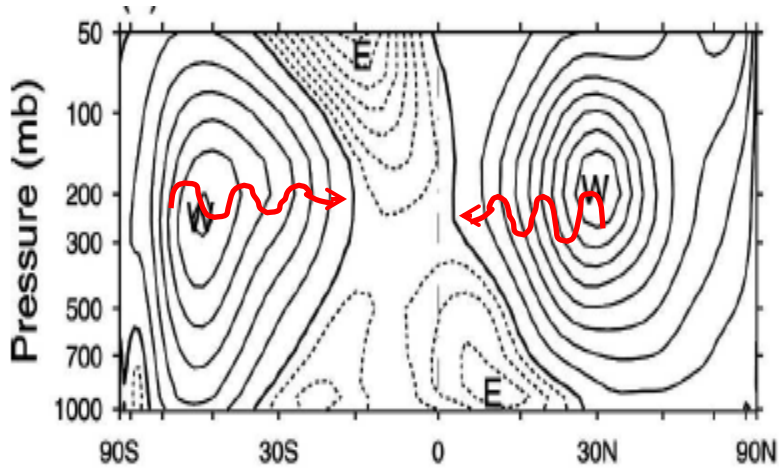


Rossby wave dispersion on the sphere: response to isolated topography

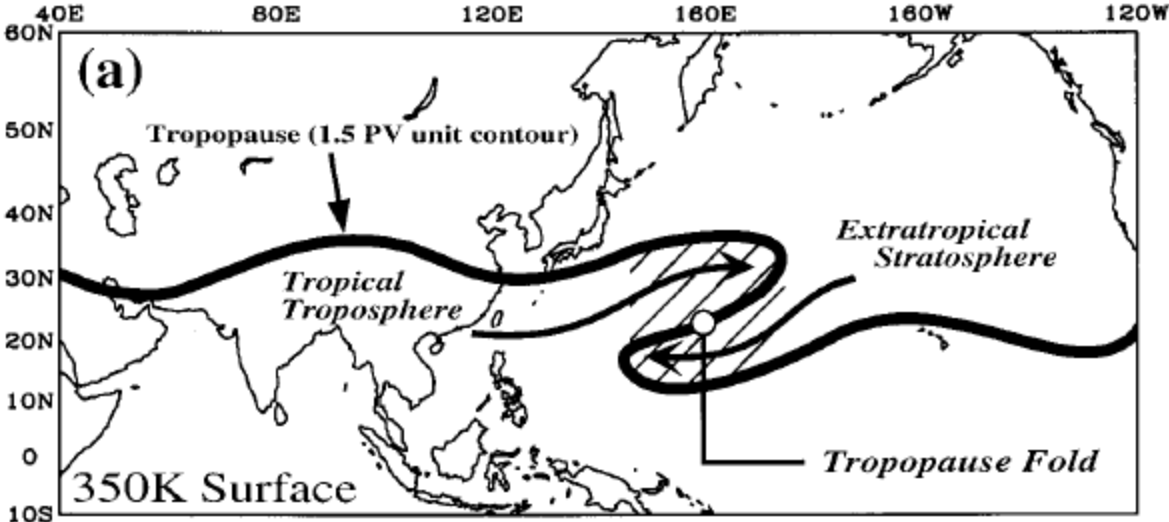


Steady state solution of the shallow water equations linearized about a constant angular velocity flow and perturbed by a circular mountain at $30^\circ N, 180^\circ E$. The model includes linear drag. After B. J. Hoskins, *Horizontal Wave Propagation on the Sphere*, in *The General Circulation. Theory, Modeling and Observations*. NCAR Summer Colloquium 1978.

Rossby waves cannot propagate into regions where $U < 0$ (i.e. across equator)



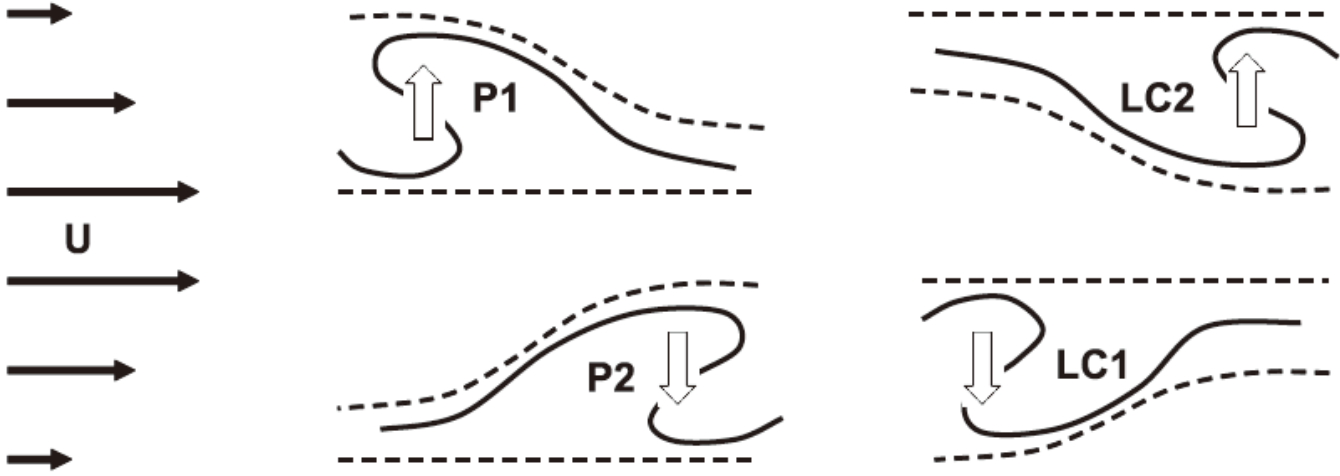
Breaking Rossby waves: overturning of PV contours



Postel and Hitchman 1999
Homeyer et al 2013

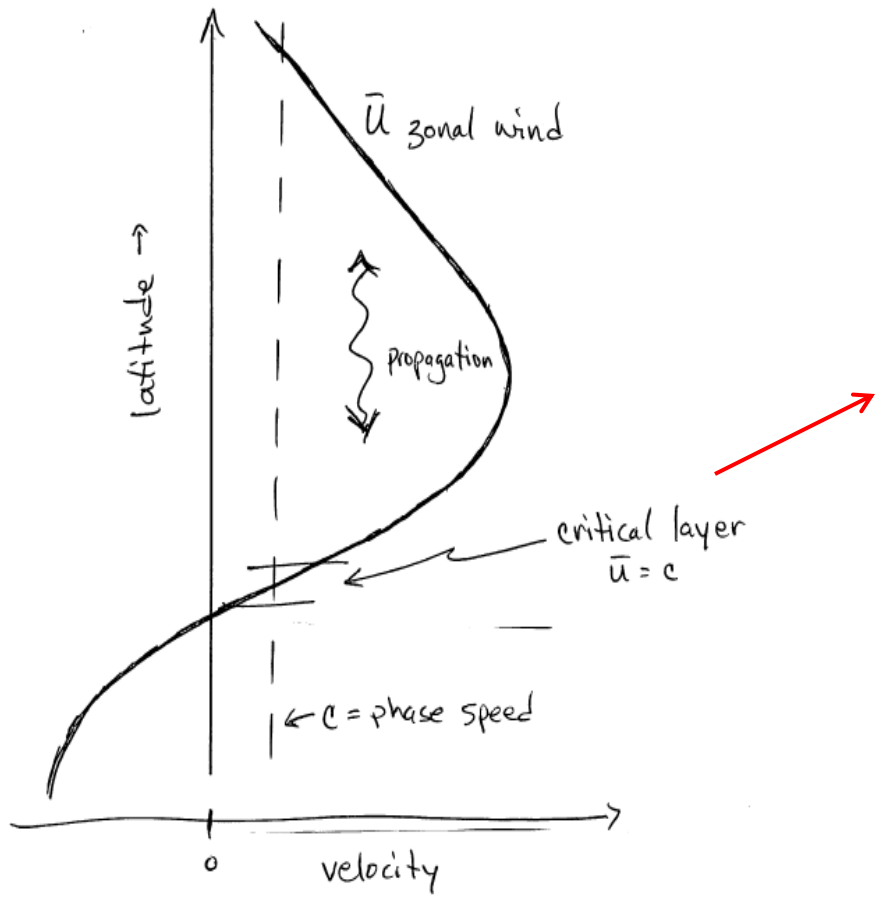
Two types of wave breaking, depending on shear of background winds

poleward breaking (cyclonic)

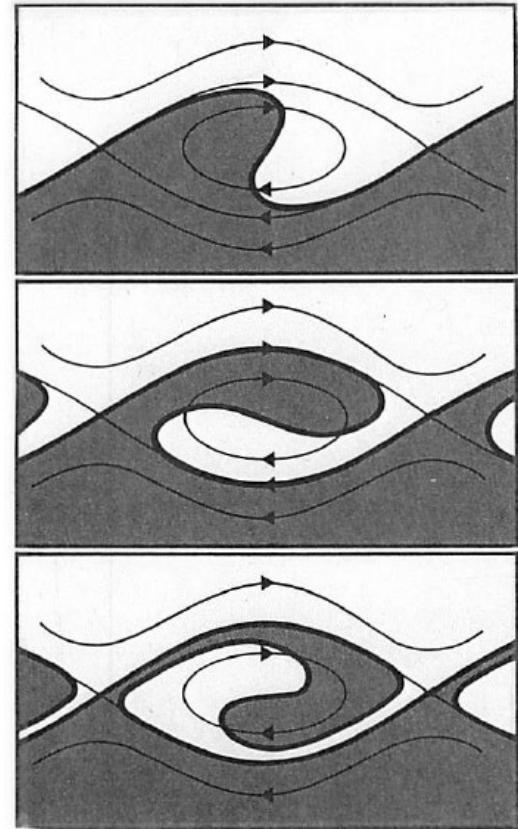


equatorward breaking (anticyclonic)

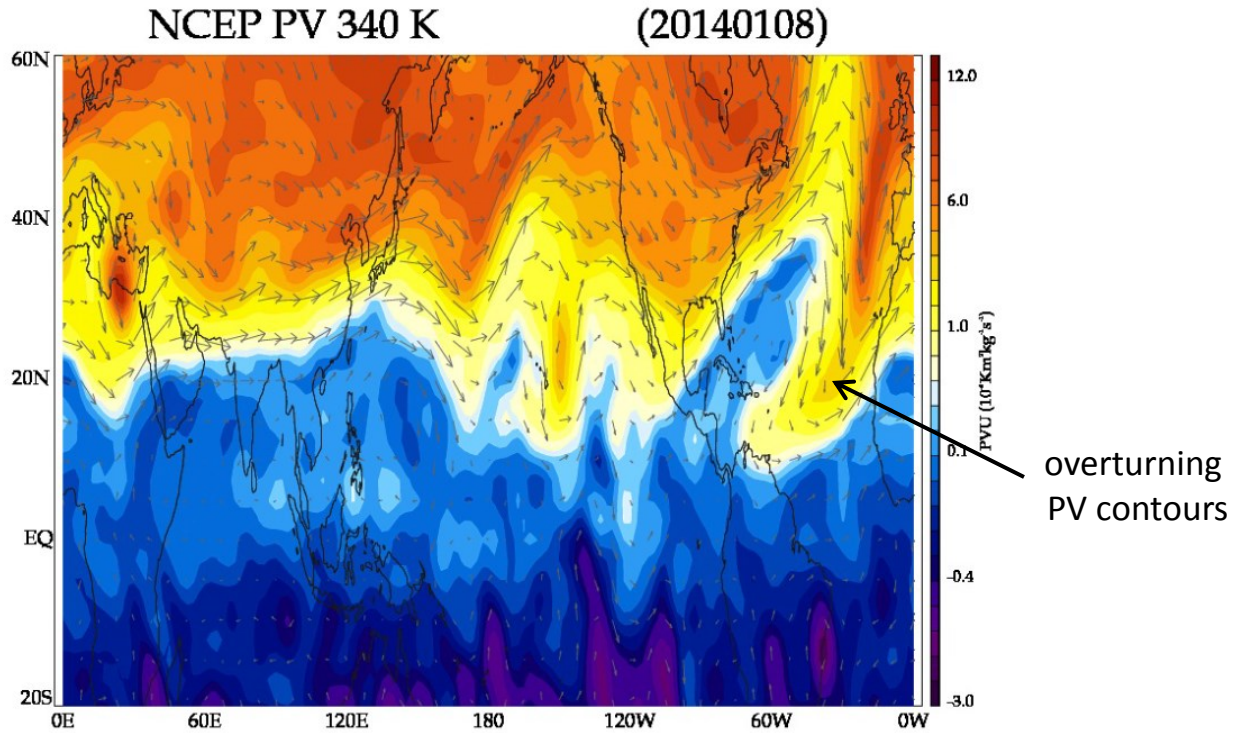
Rossby wave critical layer interactions (critical layer: $U = c$)



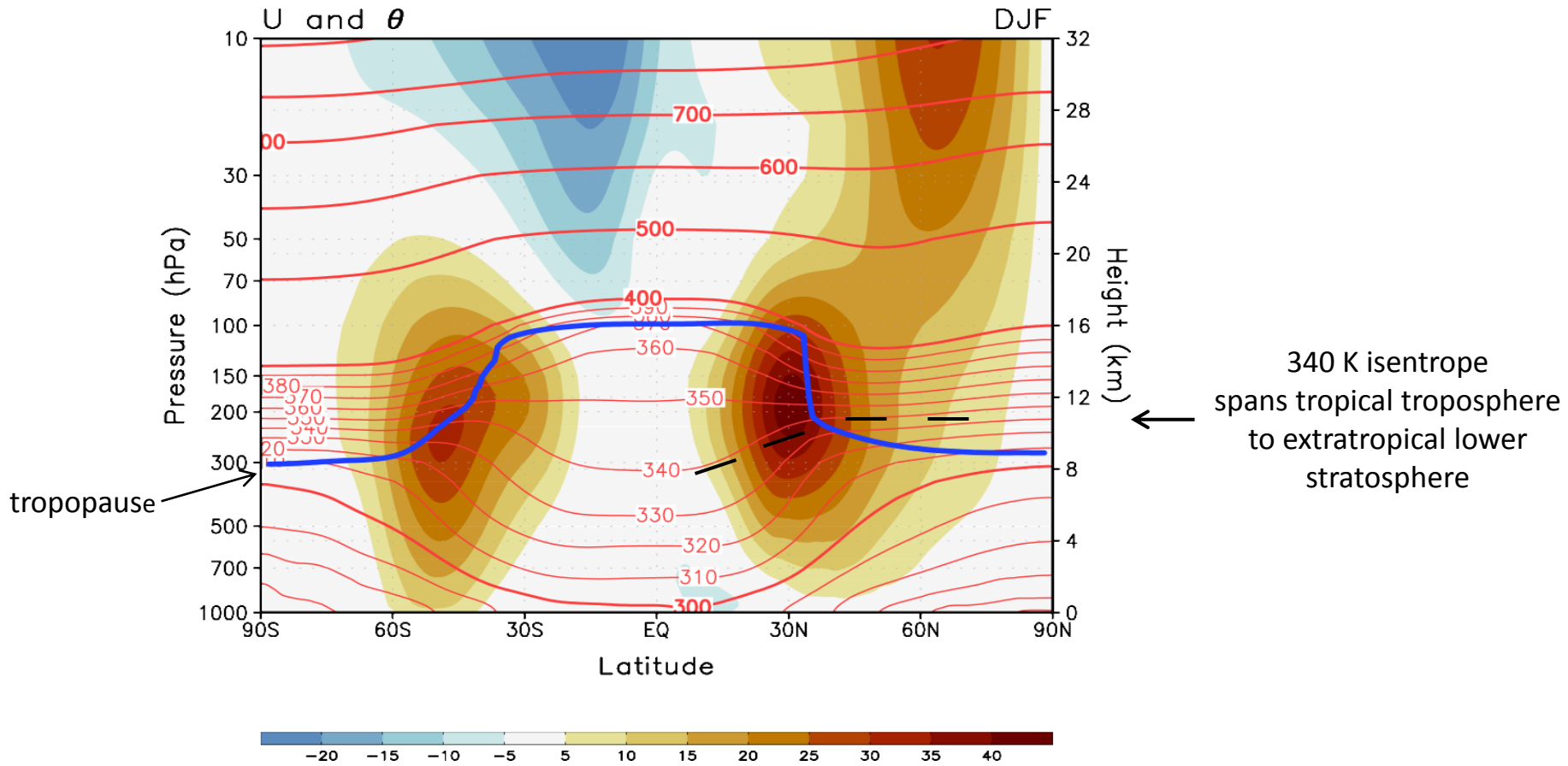
nonlinear overturning at critical layer
(irreversible transport and mixing)



Example of a large-scale breaking Rossby wave

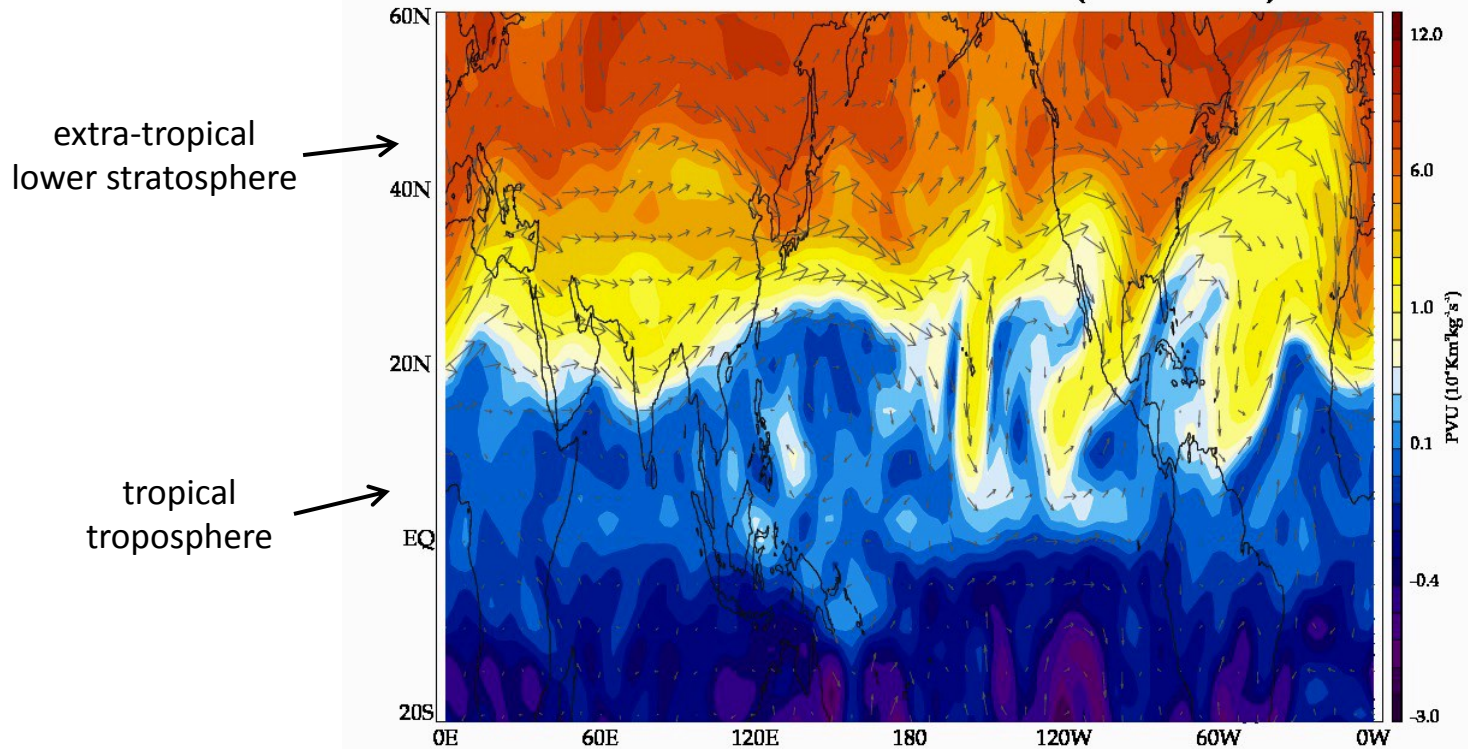


fast, synoptic flow mainly along isentropes:



Rossby waves during January-May at 340 K

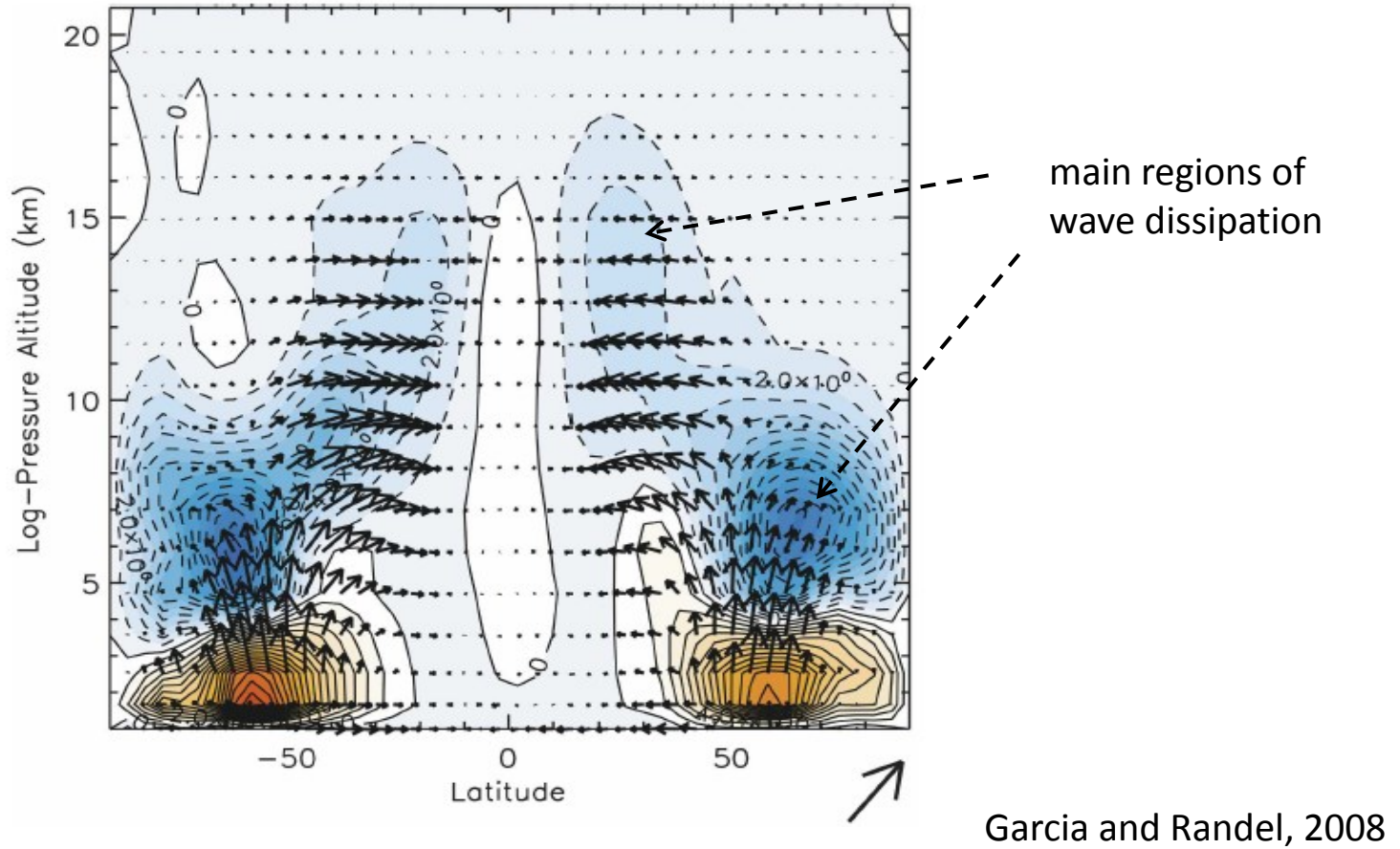
NCEP PV 340 K (20140130)



Key points:

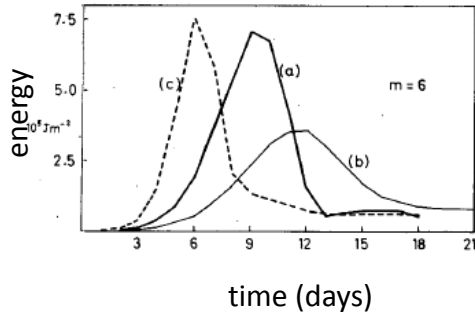
- Rossby waves: conserve PV
- General refraction of Rossby waves towards low latitudes
- Latitudinal or vertical propagation for $U > 0$ (more generally $U > c$)
- Rossby wave breaking near critical lines ($U = c$)
- Poleward or equatorward breaking depending on background U shear
- Key mechanism for dissipation, transporting PV or trace species

Climatological EP fluxes in the troposphere

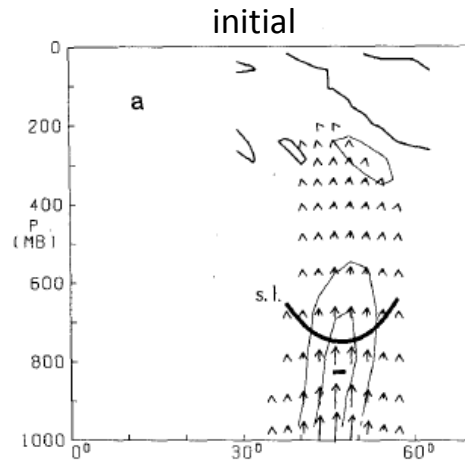


Extratropical EP flux patterns are related to baroclinic wave life cycles

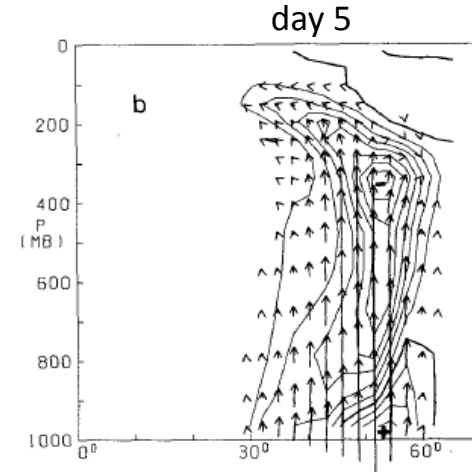
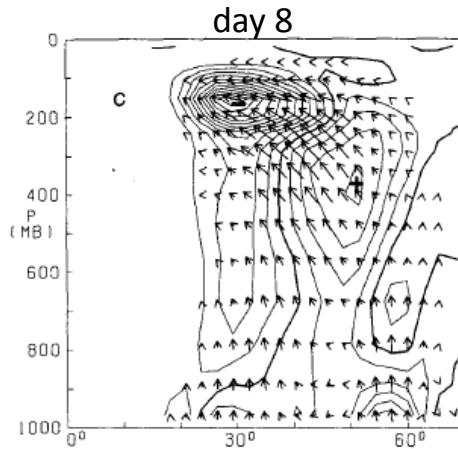
idealized
zonal wave 6
baroclinic eddy
life cycle



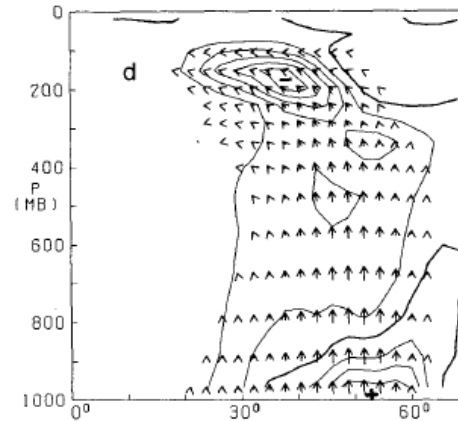
barotropic
decay



TOTAL E-P FLUX DIVERGENCE
DAY -0.00



TOTAL E-P FLUX DIVERGENCE
DAY 5.00
life cycle average



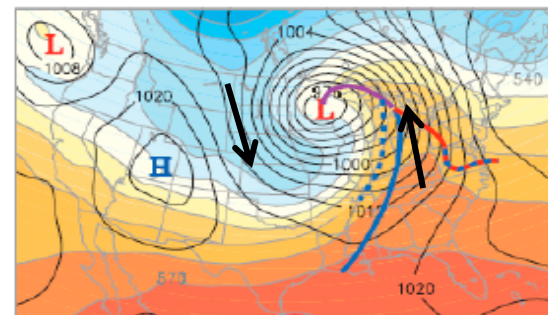
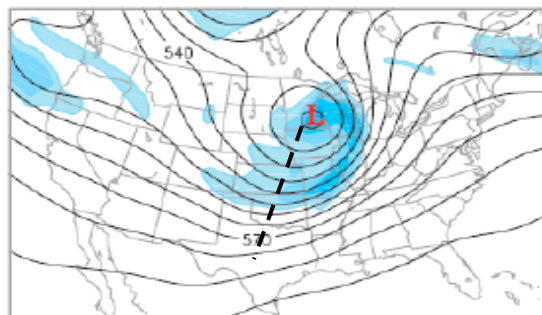
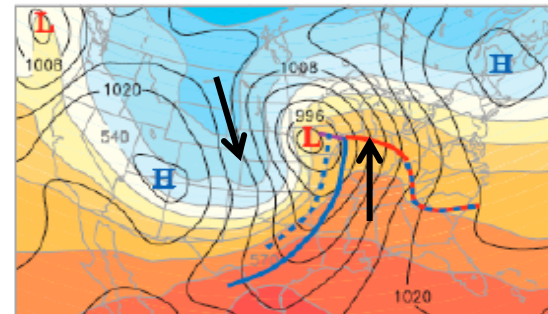
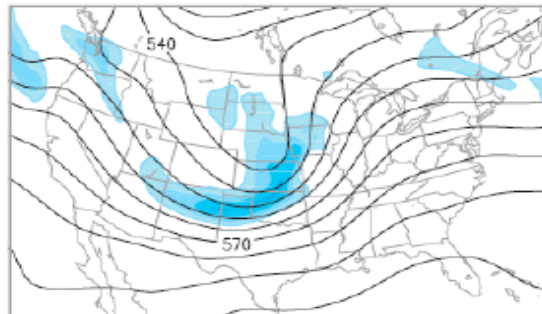
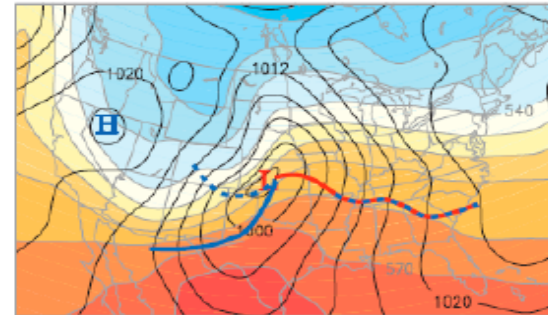
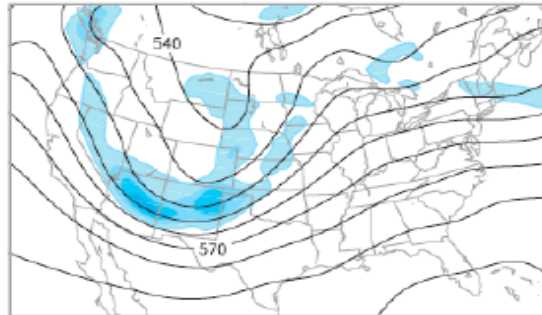
baroclinic
growth

Simmons and Hoskins 1980
Edmon et al 1980

synoptic views
of developing
baroclinic
wave

500 hPa

surface



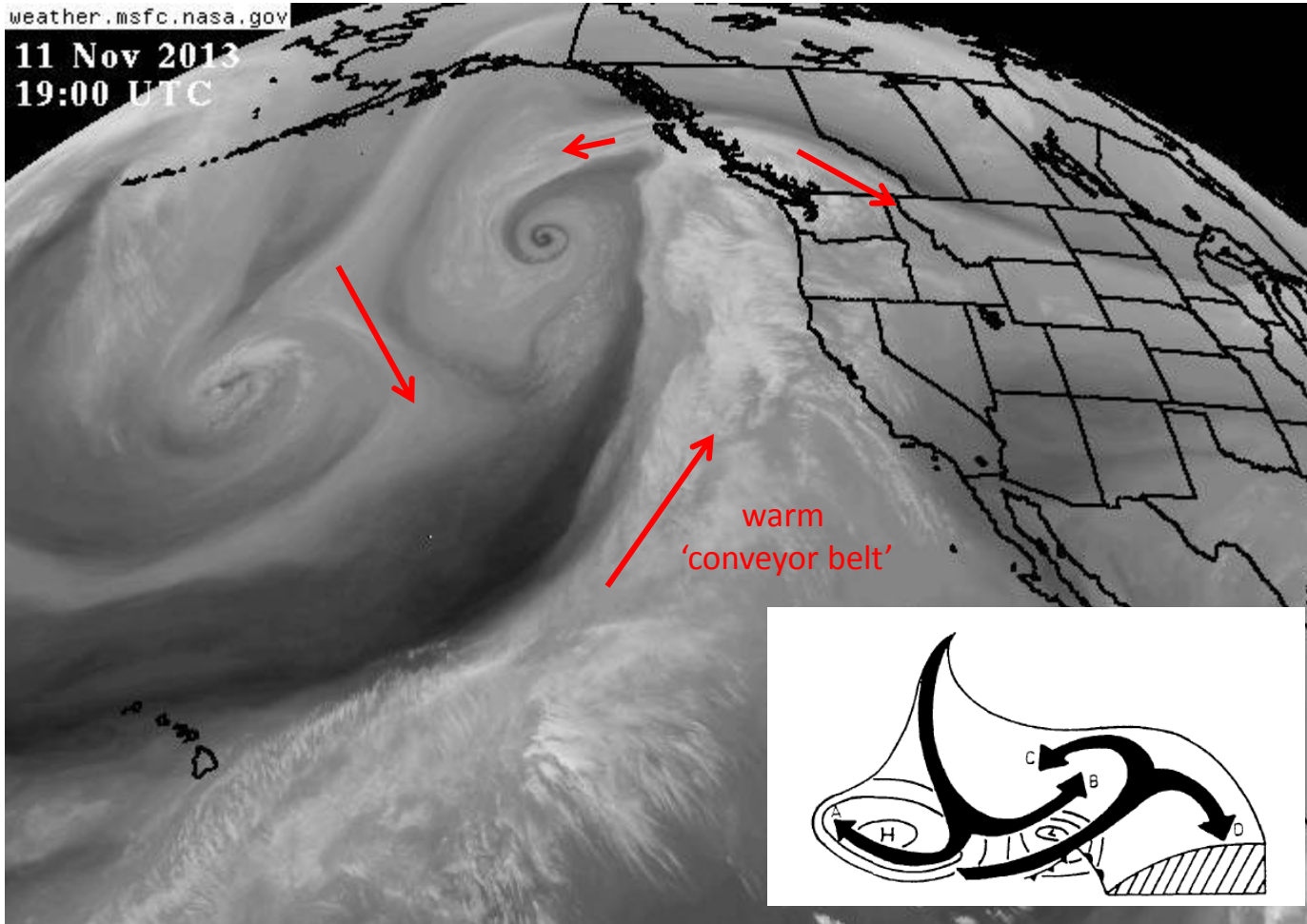
poleward
heat flux

phase tilt:
equatorward
momentum flux



weather.msfc.nasa.gov

11 Nov 2013
19:00 UTC



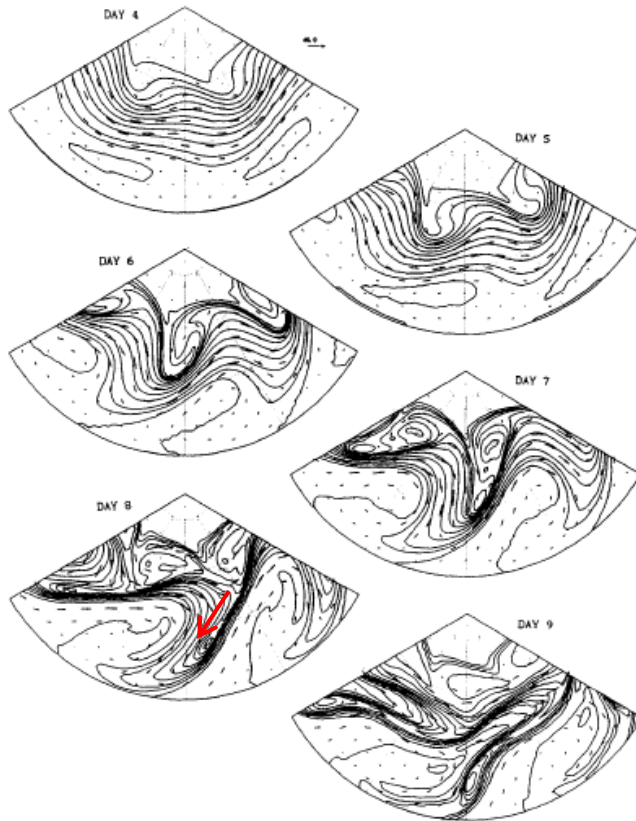
Two paradigms of baroclinic-wave life-cycle behaviour

By C. D. THORNCROFT^{1*}, B. J. HOSKINS¹ and M. E. McINTYRE²

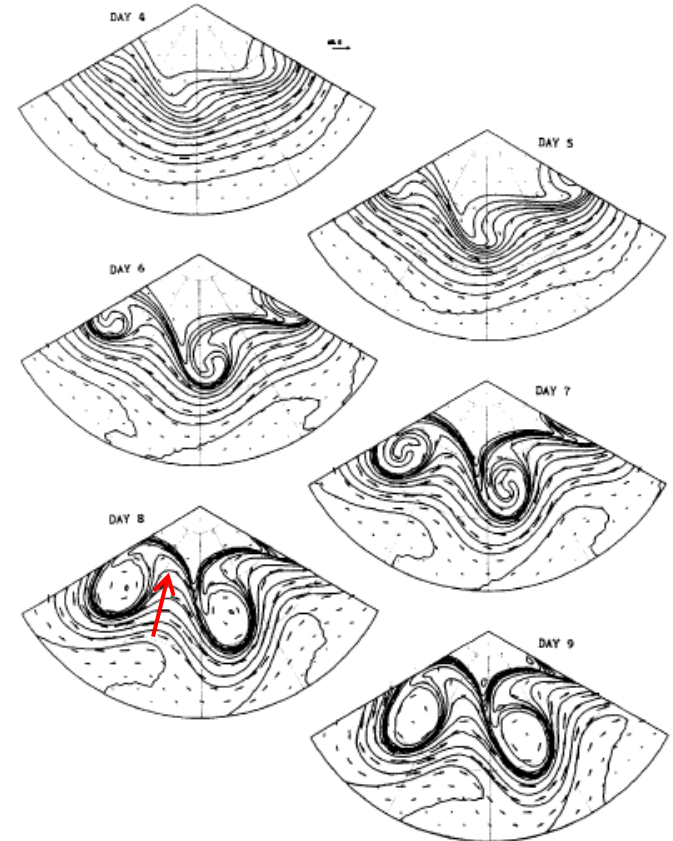
¹Department of Meteorology, University of Reading

²Department of Applied Mathematics and Theoretical Physics, University of Cambridge

LC1 equatorward breaking

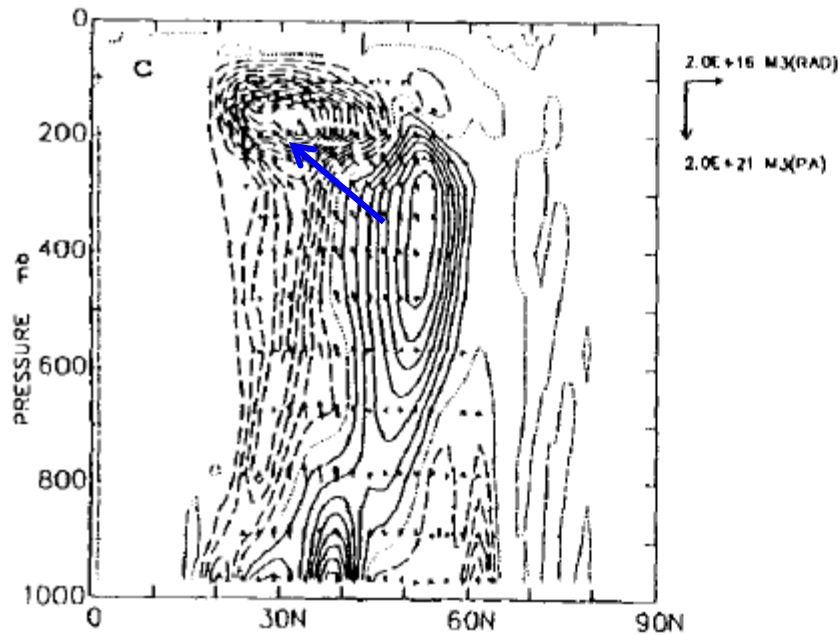


LC2 poleward breaking

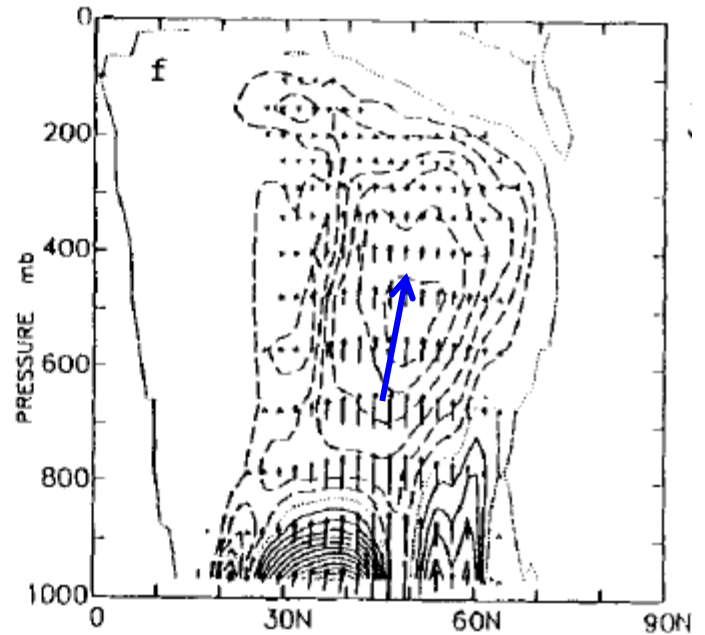


Idealized baroclinic wave life cycles

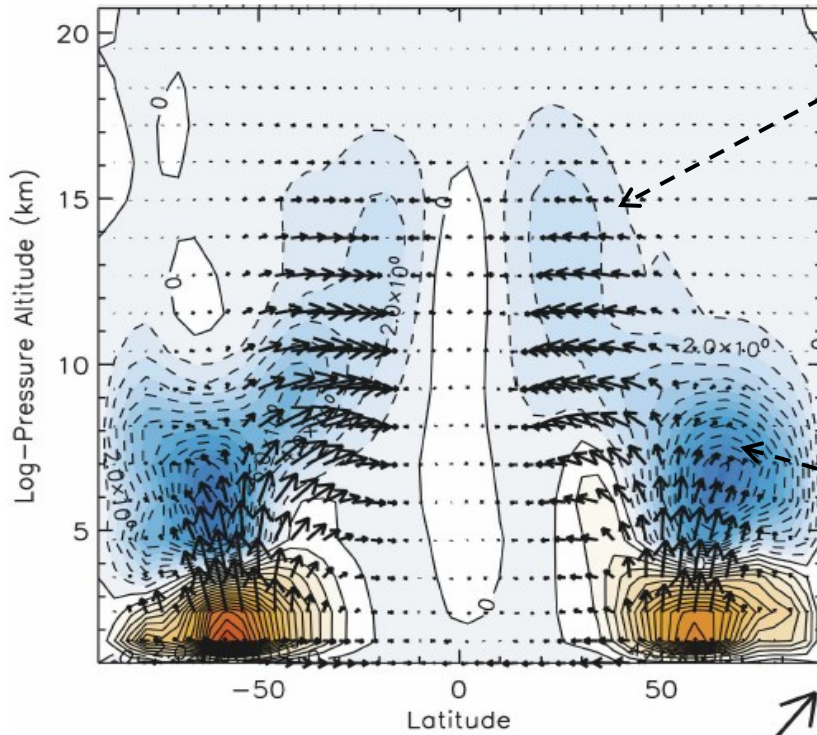
equatorward propagation (LC1)



poleward propagation (LC2)

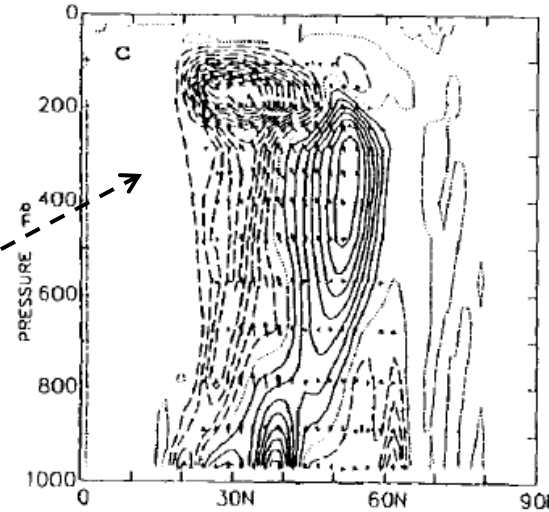


Climatology

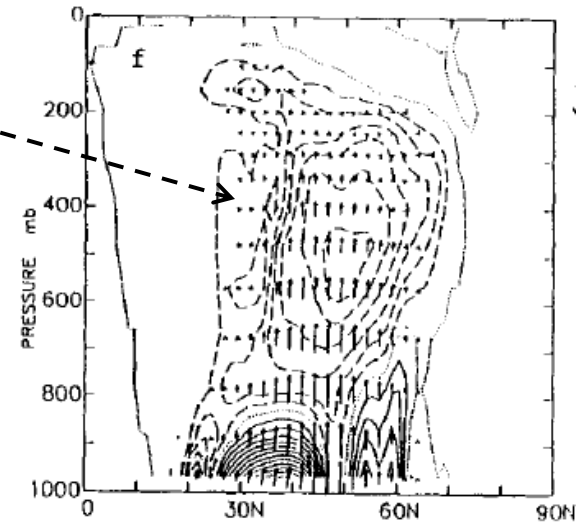


Garcia and Randel, 2008, J. Atmos. Sci.

LC1



LC2



Thorncroft et al., 1993, Q.J.R. Meteorol. Soc.

Using phase speed spectra to diagnose critical layer interactions

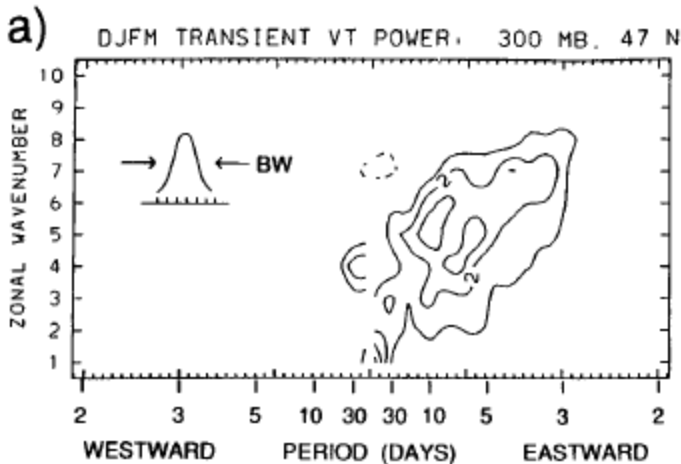
wave flux
co-spectra as a function
of zonal wavenumber
and phase speed

$$K_{n,c} = K_{k,\omega} \cdot \left(\frac{n}{a \cos \phi} \right).$$

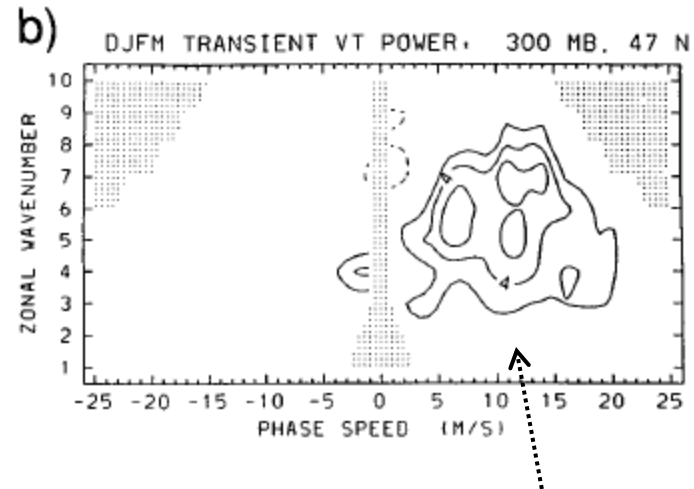
traditional
wavenumber vs.
frequency

Randel and Held 1991

wavenumber vs. frequency



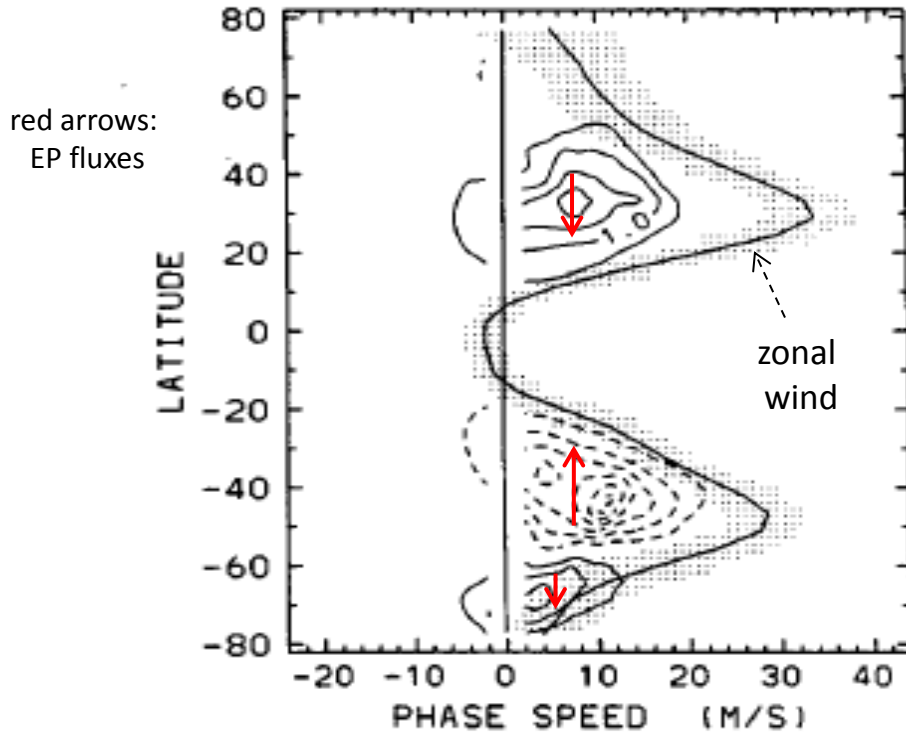
wavenumber vs. phase speed



Rossby waves move eastward
at ~5-15 m/s

Integrate over wavenumber to derive eddy flux phase speed spectra

eddy momentum flux $u'v'$ 300 hPa

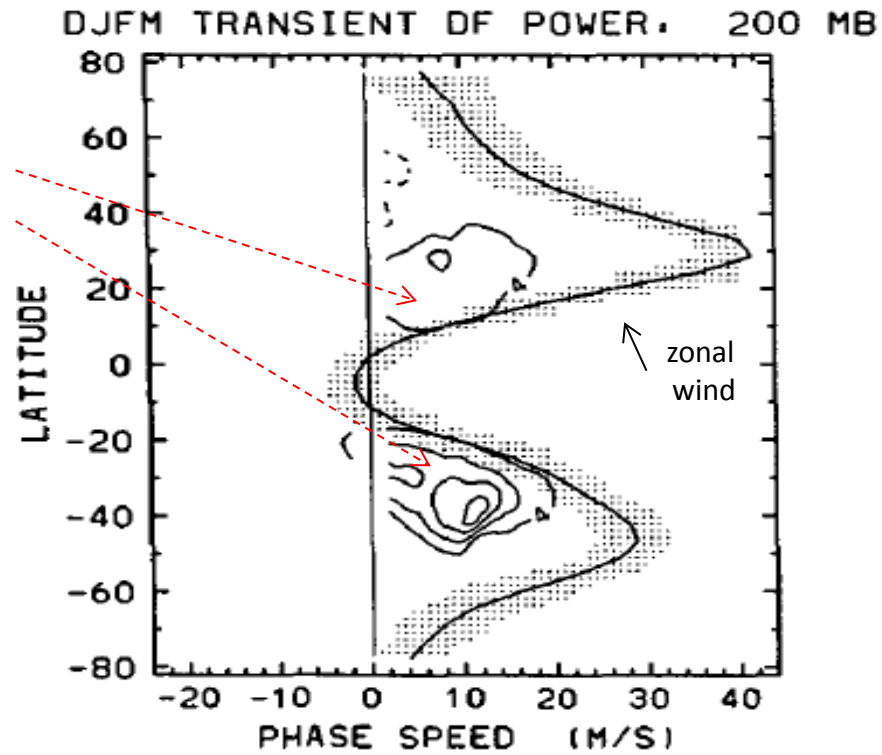


- EP fluxes: propagation to near critical lines ($c = U$)
- evidence for critical layer behavior

Randel and Held 1991

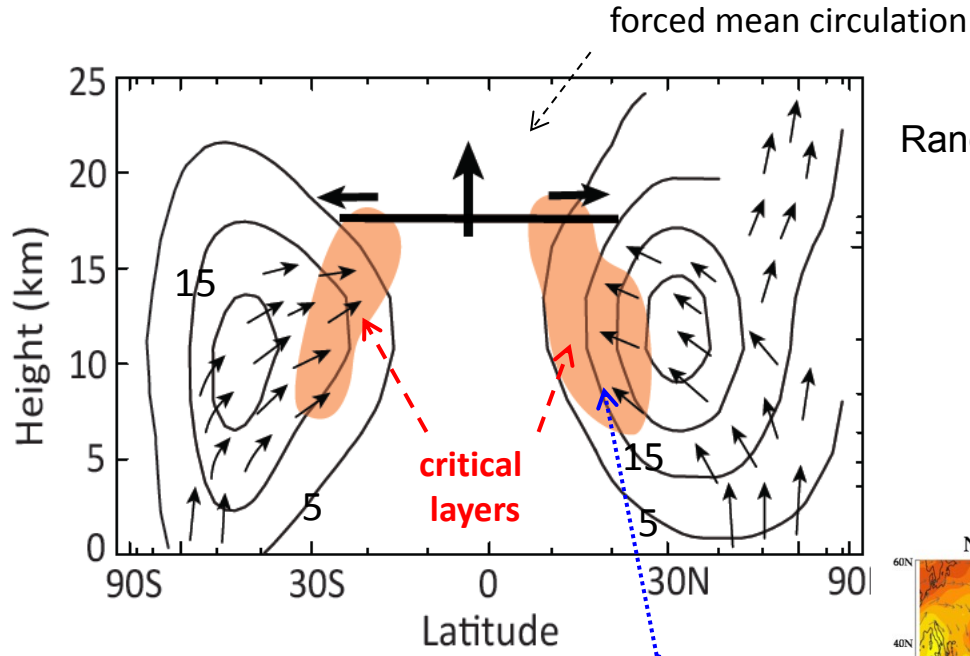
EP flux divergence phase speed spectra

EP flux divergence
poleward of critical lines

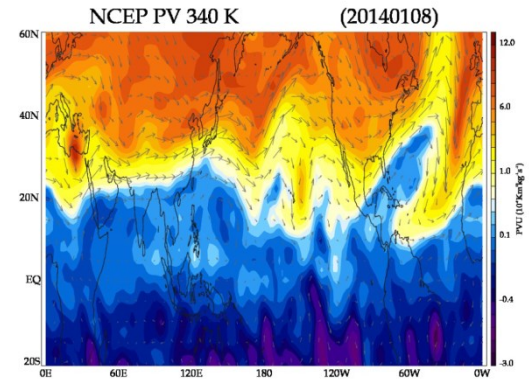


Randel and Held 1991

Subtropical critical layers for Rossby waves with phase speeds $\sim 5-15$ m/s



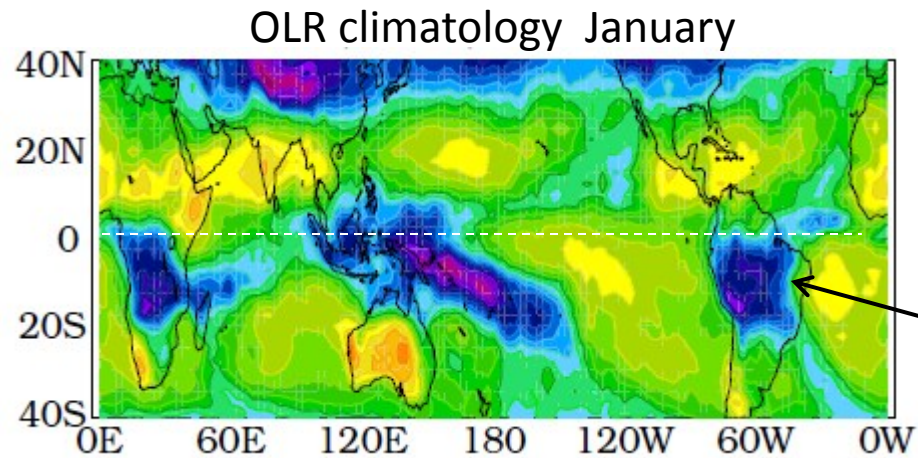
Randel and Jensen, 2013, Nat. Geosci.



Key points:

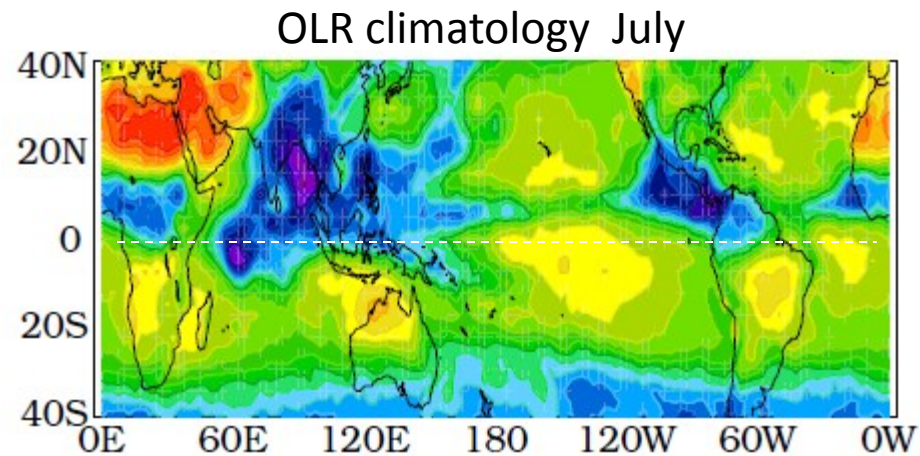
- Baroclinic wave life cycles: baroclinic growth and barotropic decay
- Two idealized types of life cycles: equatorward and poleward wave breaking (LC1 and LC2)
- Consistent with tropospheric EP flux circulation statistics
- Phase speed spectra: clear evidence for critical layers in subtropics (important influence of extratropical waves on tropical circulations)

Large-scale tropical circulations are forced by latent heating from deep convection



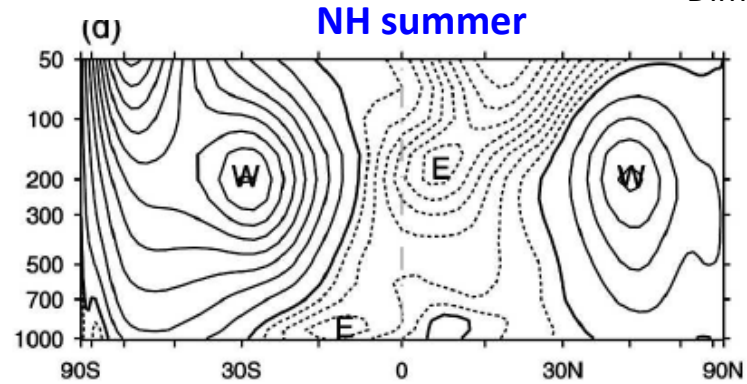
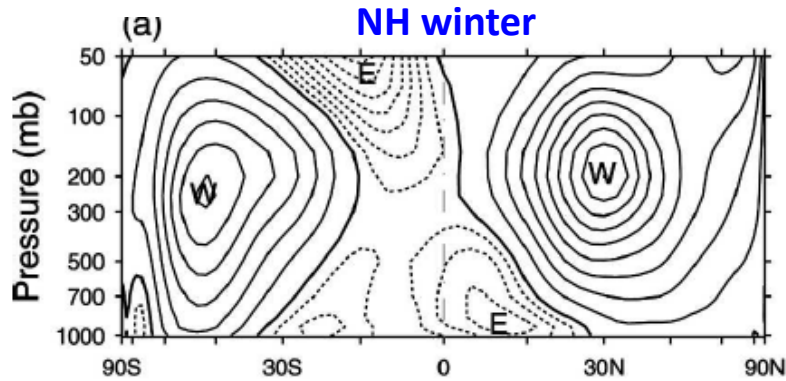
Outgoing Longwave Radiation (OLR) is a useful proxy for deep convection

high clouds ~
deep convection

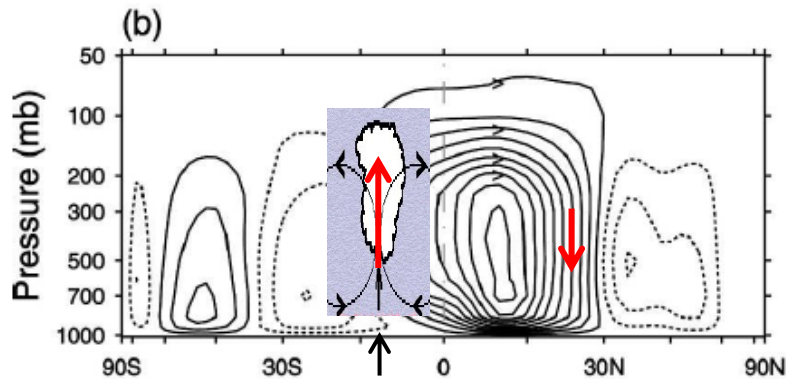


Seasonal variations in tropical overturning circulation (Hadley cell)

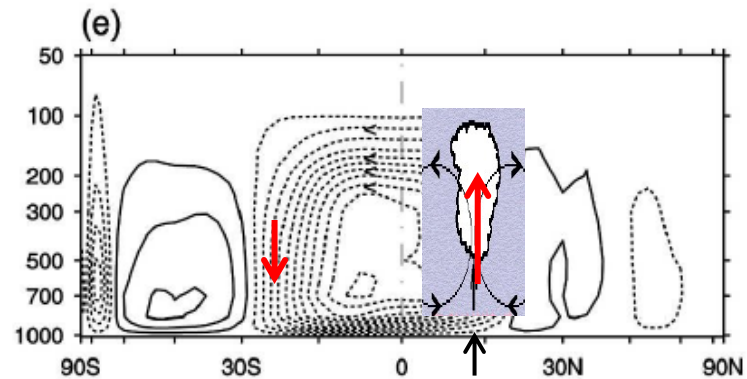
Dima et al 2005



zonal winds



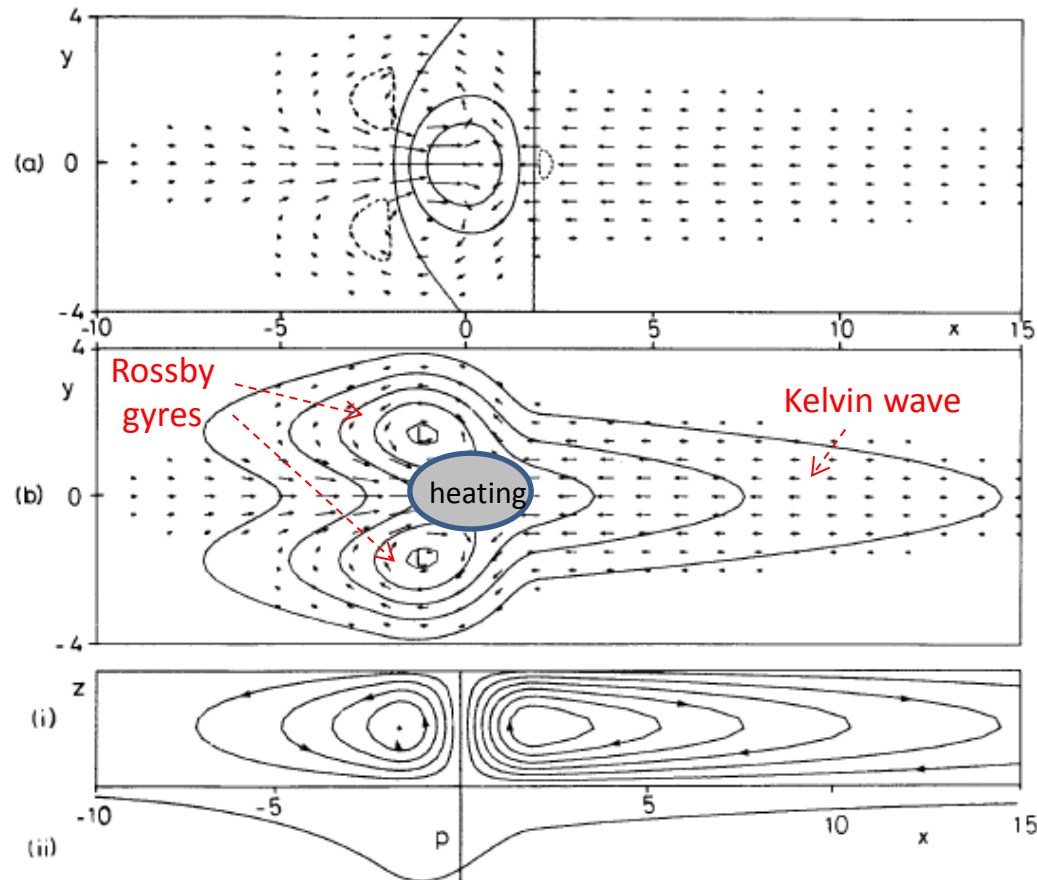
deep
convection



deep
convection

streamfunction

Dynamical response to low frequency convective forcing



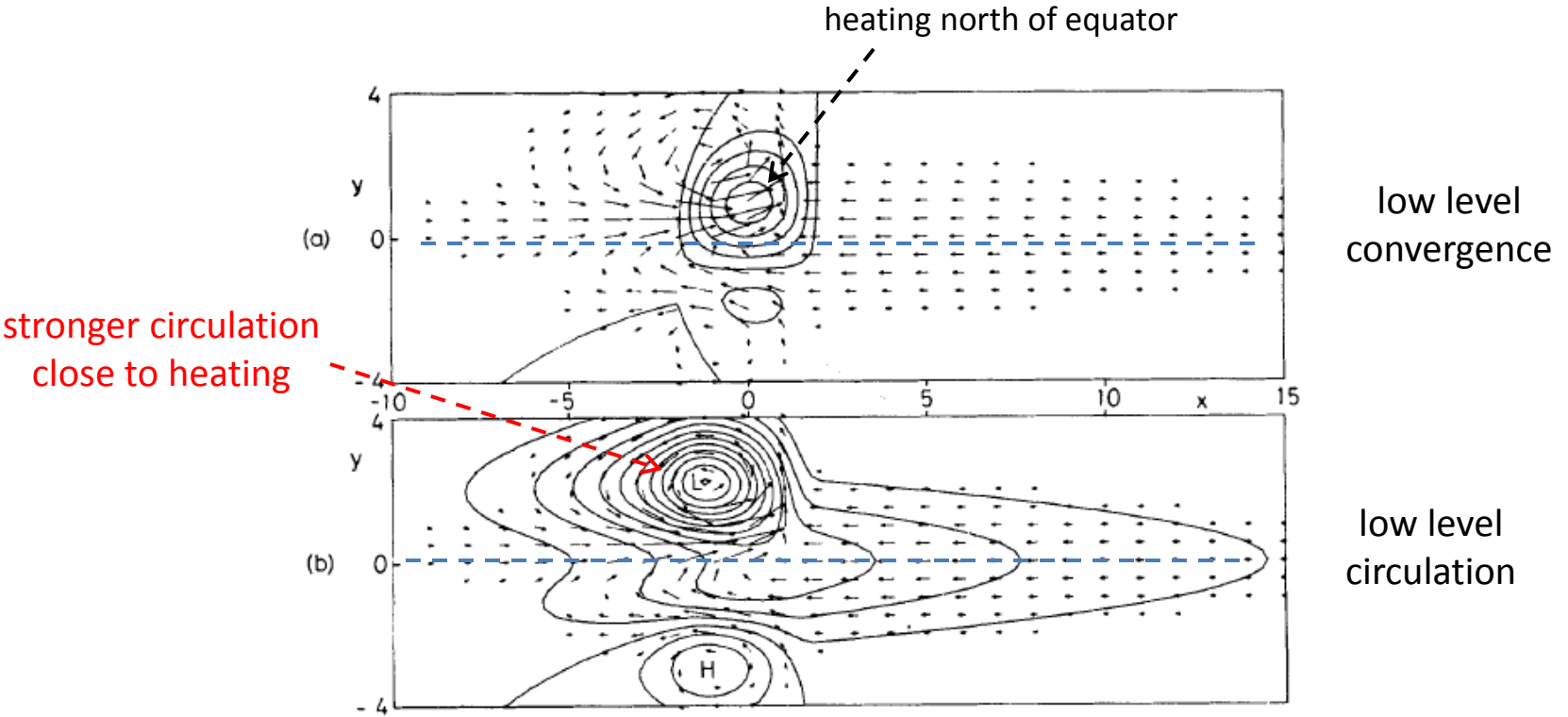
low level
convergence

low level
circulation

vertical structure

Gill, 1980

Dynamical response to heating centered north of equator

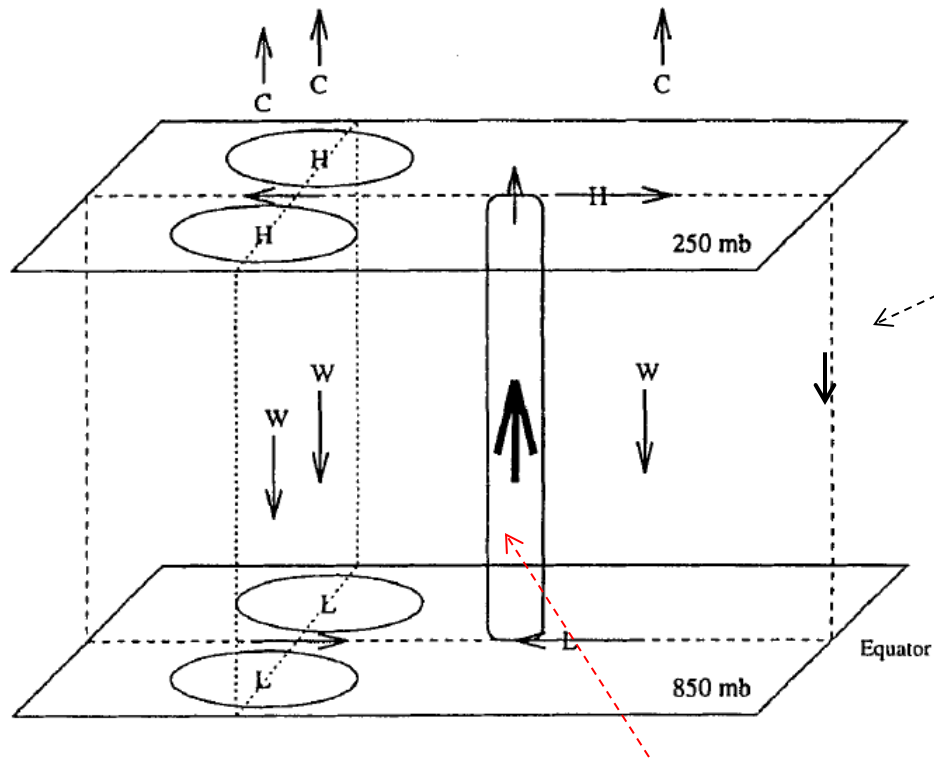


The tropical tropopause

By E. J. HIGHWOOD* and B. J. HOSKINS
University of Reading, UK

vertical structure: out of phase between
lower troposphere and upper troposphere

QJRMS 1998

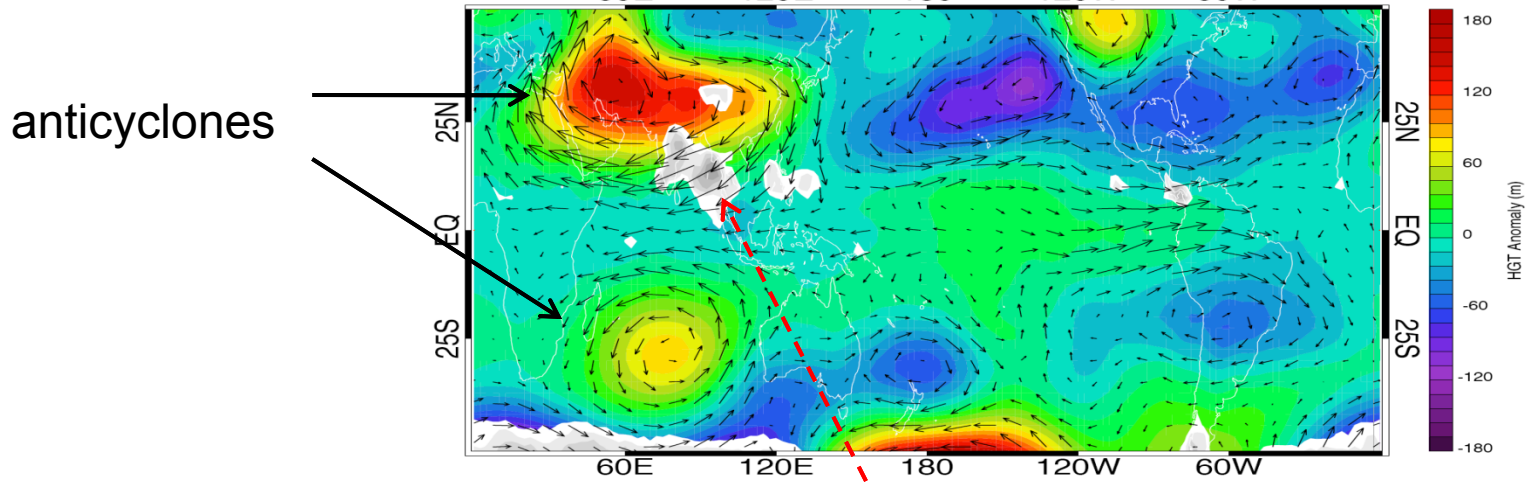


Walker circulation
(east-west)

deep convection / heating

Tropical heating produces subtropical anticyclones in the UTLS

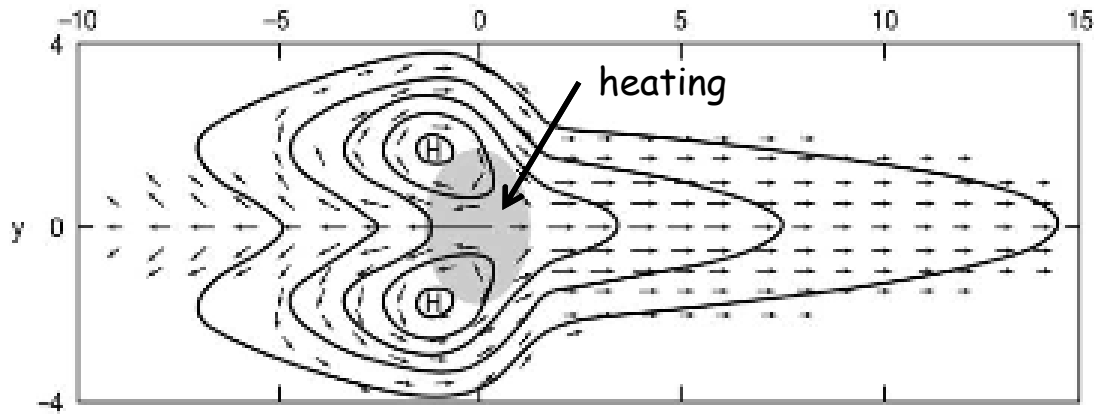
JJA geopotential height and winds 100 hPa



anticyclones

Convection (heating)

Park et al 2007



Matsuno-Gill Solution

Key points:

- Organized deep convection (latent heating) drives large-scale tropical circulations
- Seasonal movement between solstices (SH – NH subtropics)
- Hadley and Walker overturning circulations
- Matsuno-Gill dynamical response to local heating: subtropical Rossby waves and equatorial Kelvin waves
- Subtropical anticyclones in UTLS (especially Asian monsoon during NH summer)

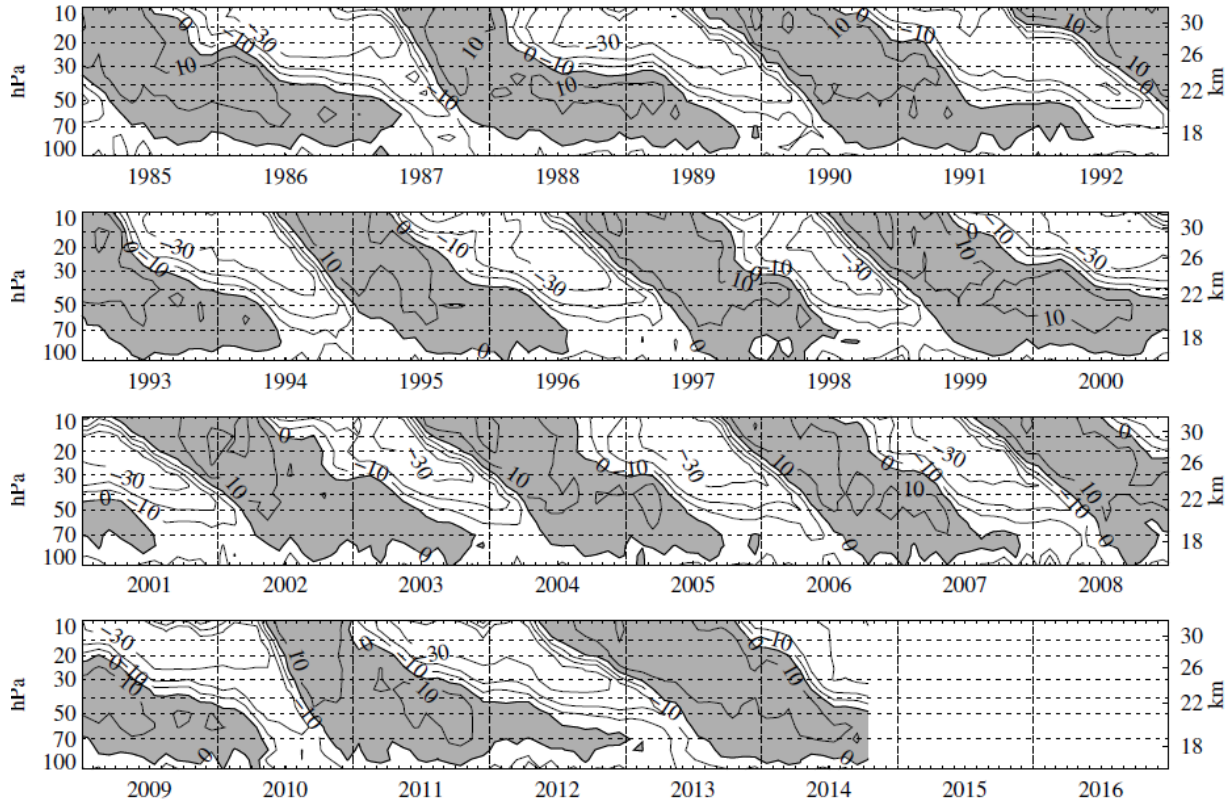
The Quasi-Biennial Oscillation (QBO) and El Niño Southern Oscillation (ENSO) are important modes of circulation in the stratosphere, troposphere and UTLS

- QBO: approximate 28-month oscillation in tropical stratosphere, forced by vertically propagating waves interacting with mean flow
- ENSO: ~2-6 year time scale oscillations of convection, oceanic and atmospheric circulations. Wave effects extend into high latitudes and into tropical lower stratosphere.
- Both QBO and ENSO influence temperature, circulation and transport in UTLS

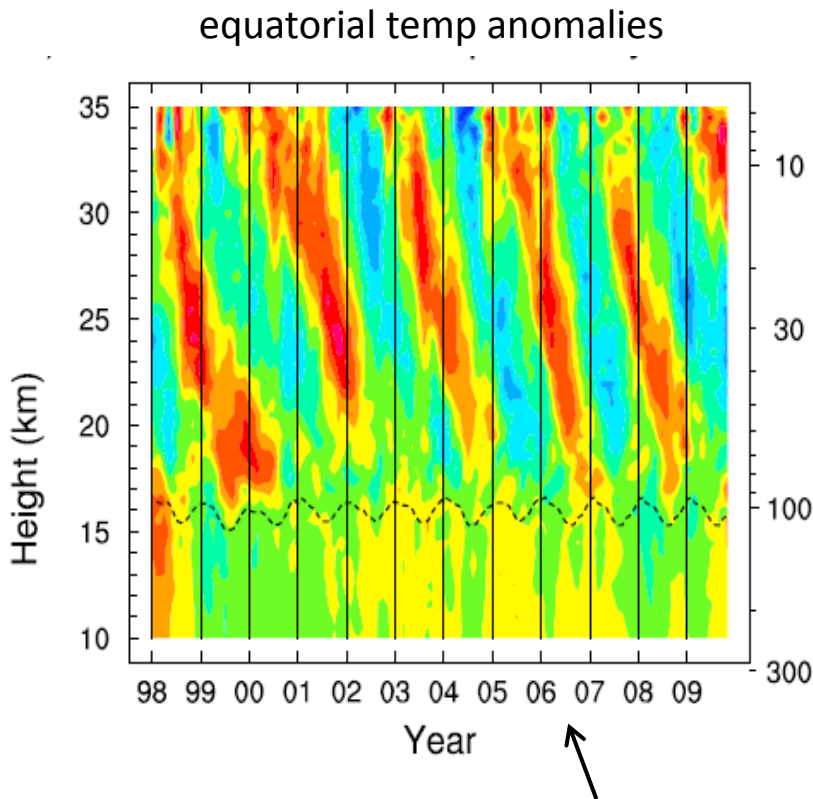
Brief observational characteristics of QBO

easterly descent is regular;
sometimes the westerly phase 'stalls'

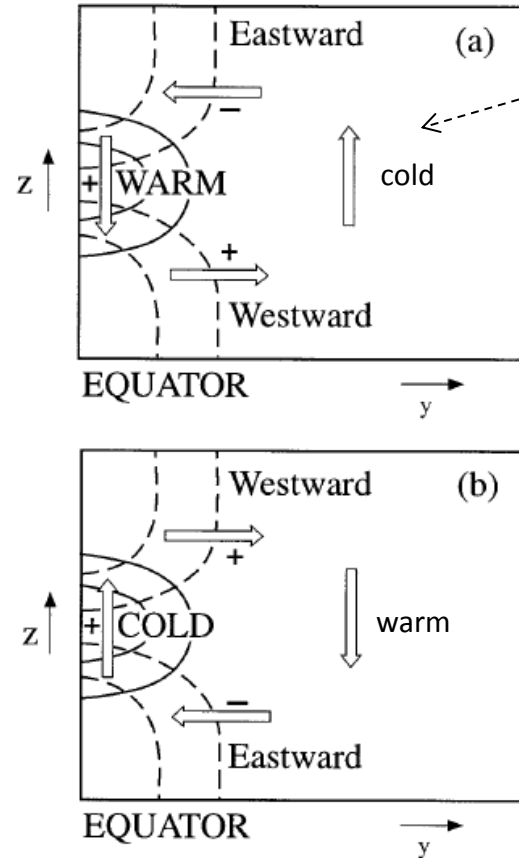
Zonal winds at Singapore (1° N) - standard QBO reference



Meridional circulations and temperatures linked with the QBO



- temp anomalies at equator up to +/- 4 K
- extend downward to near tropopause

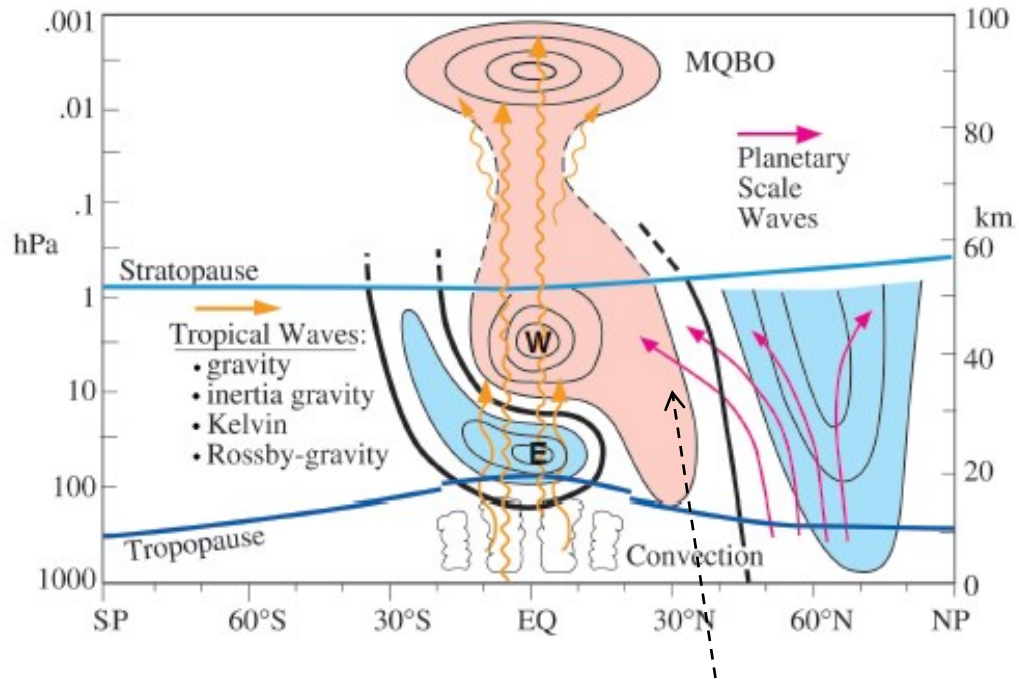


extension to middle latitudes

Plumb and Bell, 1982b, Q.J.R. Meteorol. Soc.

Meridional structure and global influences of the QBO

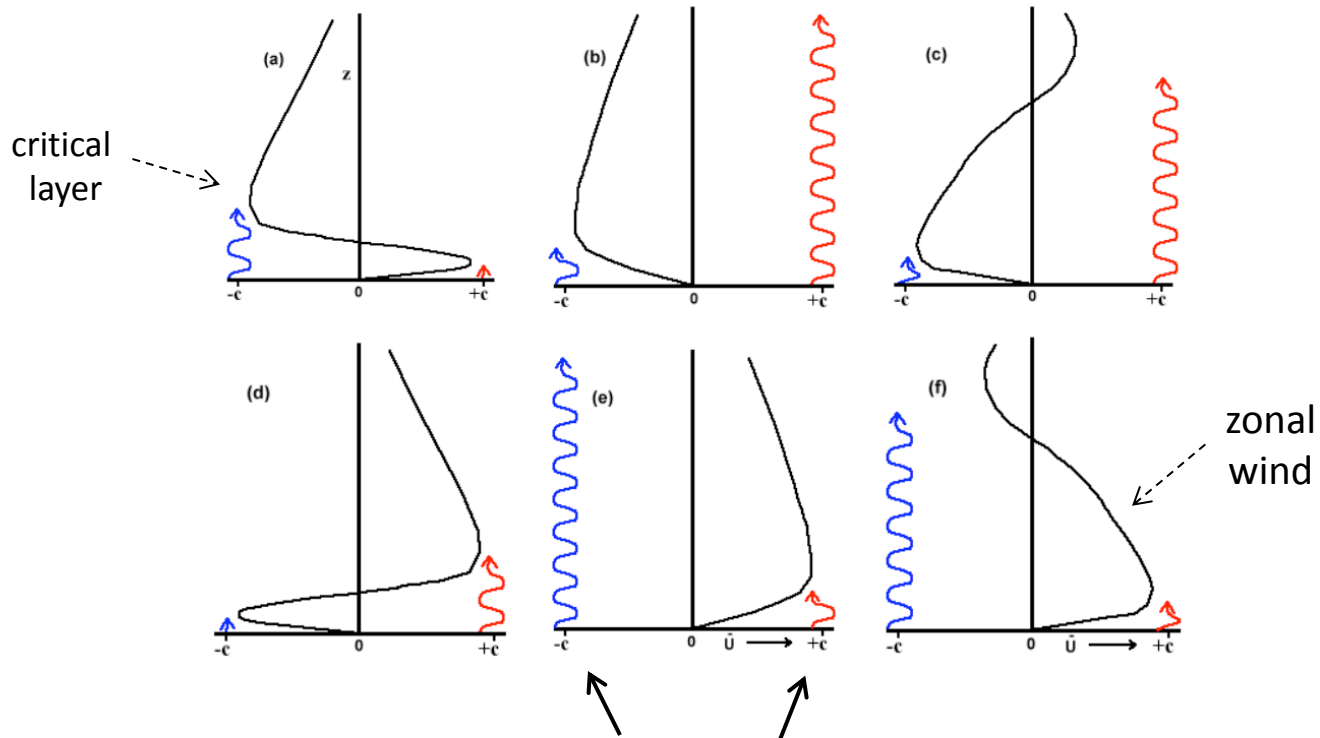
Baldwin et al 1998



influence on high latitudes
via subtropical critical lines

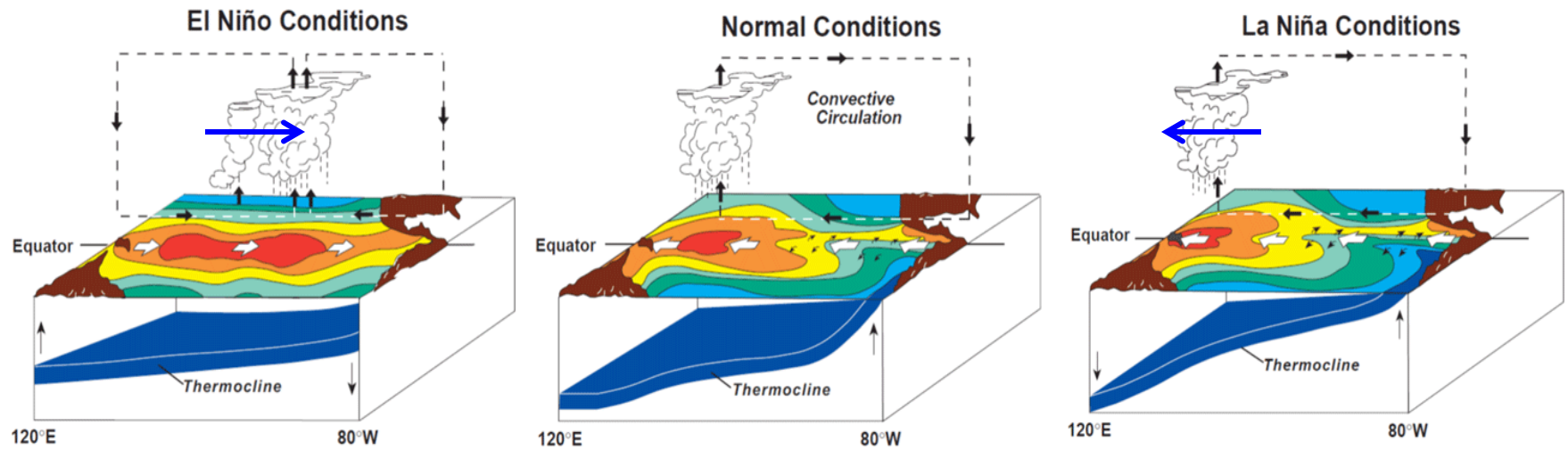
- waves propagate vertically until reaching critical level ($U=c$)
- wave absorption produces zonal acceleration (towards phase speed of wave)

Lindzen-Holton mechanism for the QBO

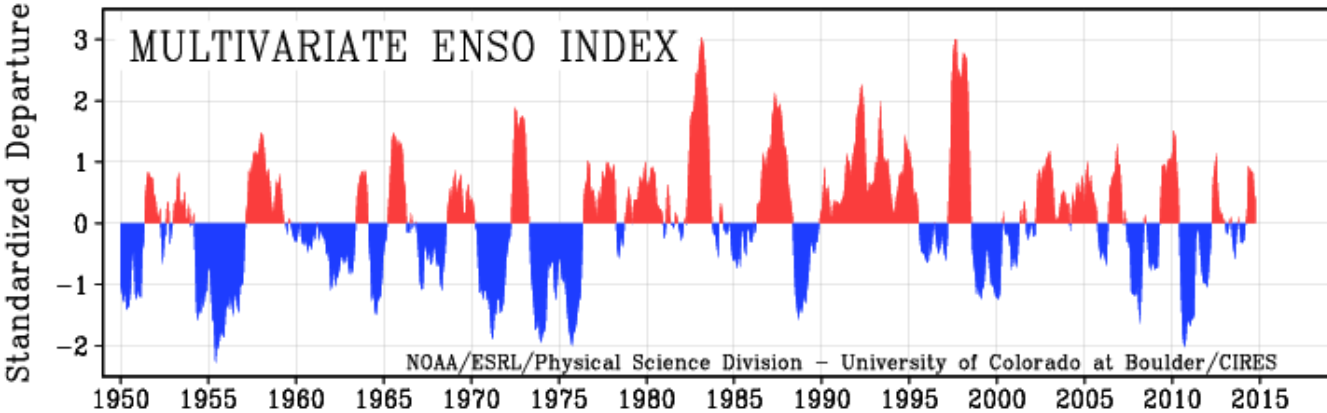


eastward and westward propagating waves forced from troposphere

ENSO – large-scale shifts in tropical convection, winds, temperatures

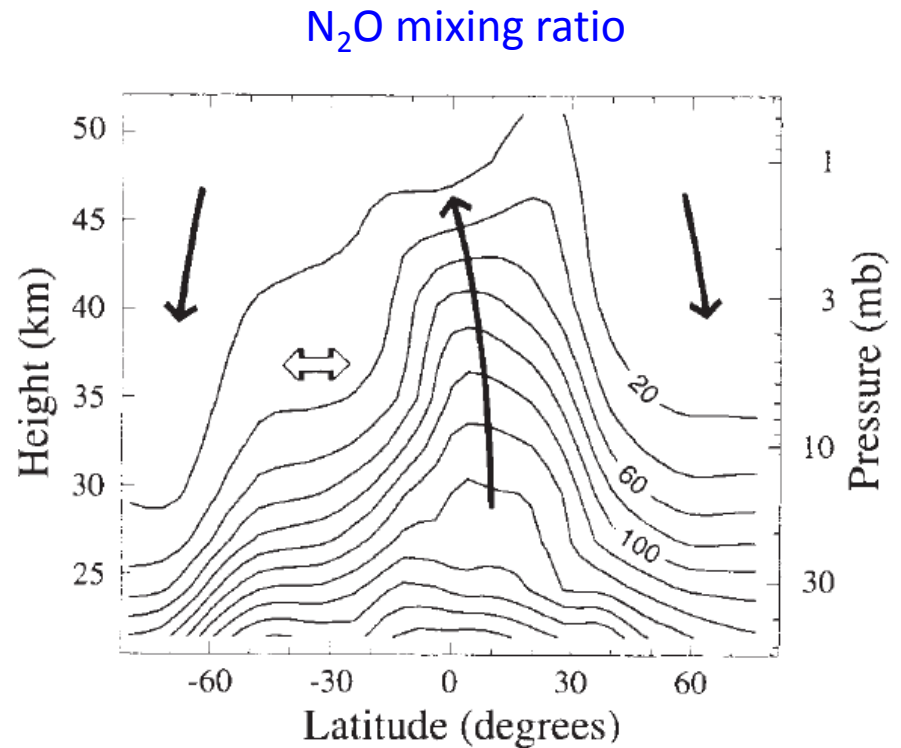


Standard time series for ENSO

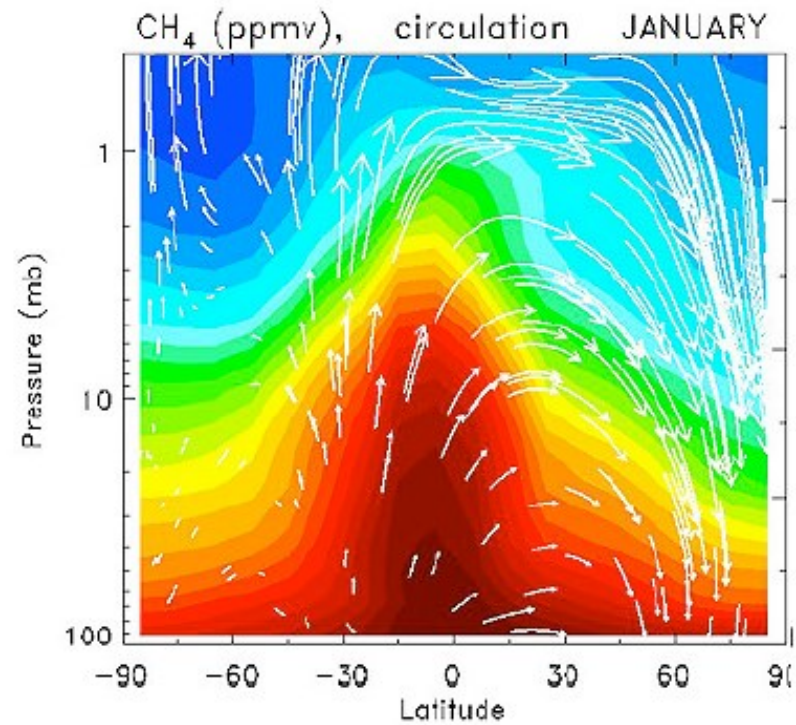


Stratospheric tracer transport: evidence from satellite observations

- N_2O is a 'tropospheric source gas'
- destroyed by photolysis (radiation) in upper stratosphere
- Source of reactive nitrogen (NO , NO_2) in upper stratosphere; important for stratospheric ozone
- Behavior reflects Brewer-Dobson circulation and eddy mixing

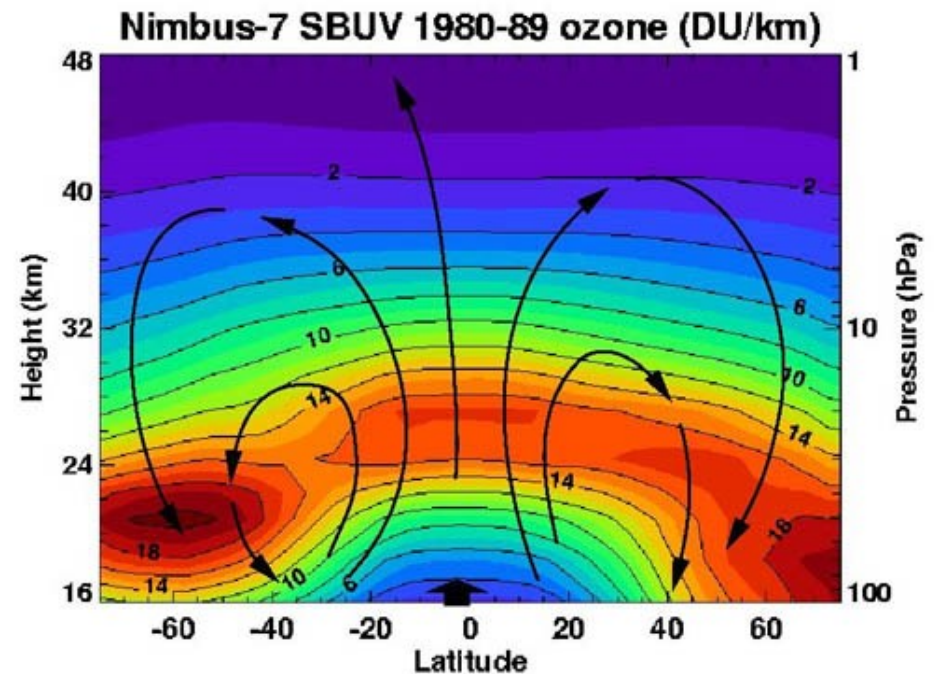


Methane CH_4 is another tropospheric source gas, oxidized to H_2O in upper stratosphere

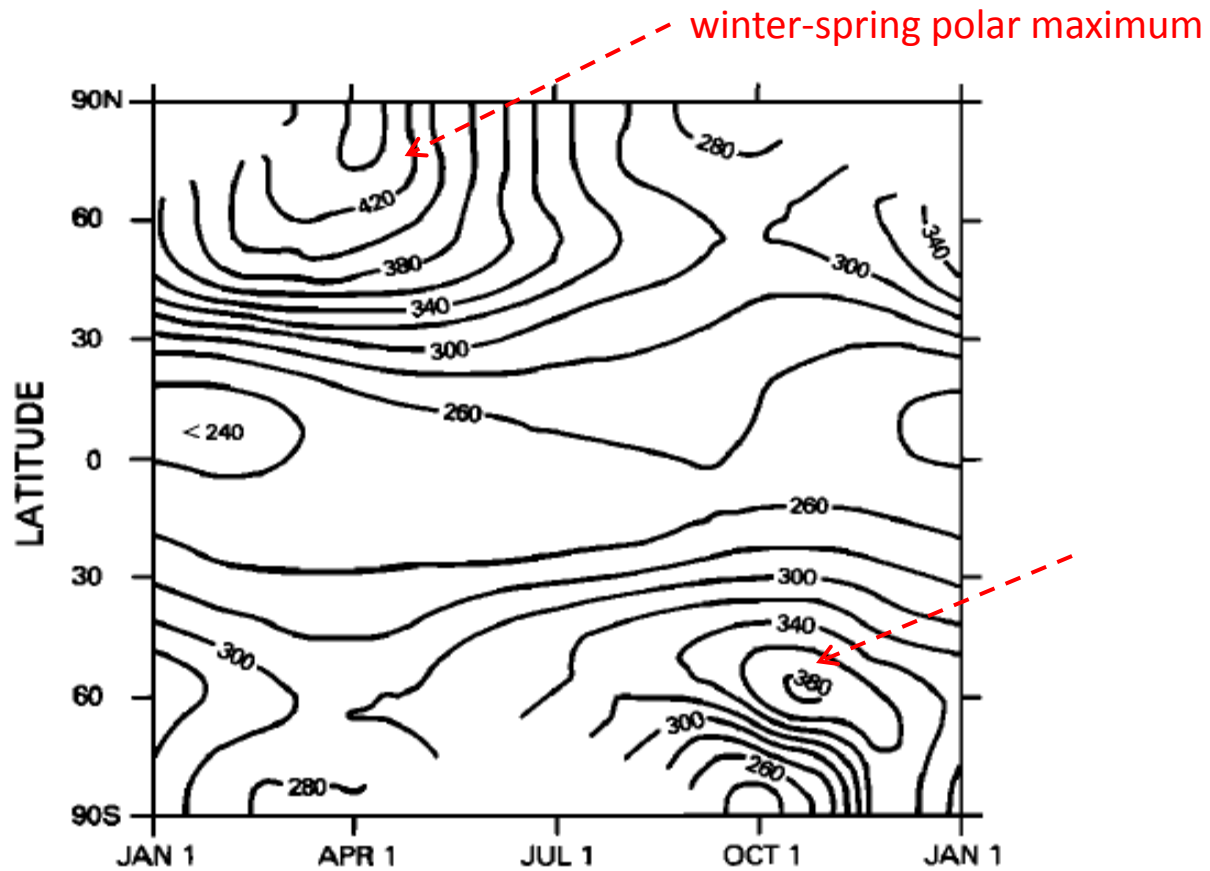


Observed ozone and Brewer-Dobson circulation

- ozone is made in the tropical stratosphere
- Short lifetime in upper stratosphere
- Long lifetime in lower stratosphere
- transport causes high latitude maximum during winter / spring



Seasonal cycle of column ozone reflects Brewer-Dobson circulation

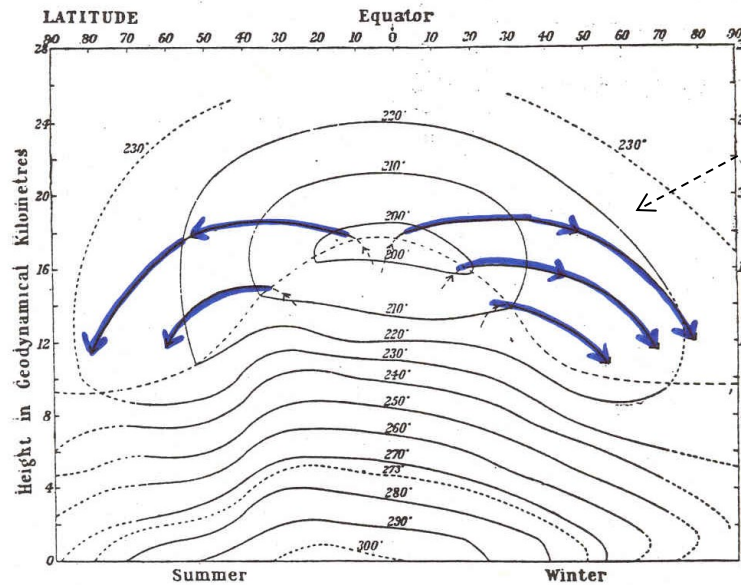


Bowman and Krueger, 1982

EVIDENCE FOR A WORLD CIRCULATION PROVIDED BY MEASUREMENTS OF HELIUM AND WATER VAPOUR DISTRIBUTION IN THE STRATOSPHERE

QJRMS, 1949

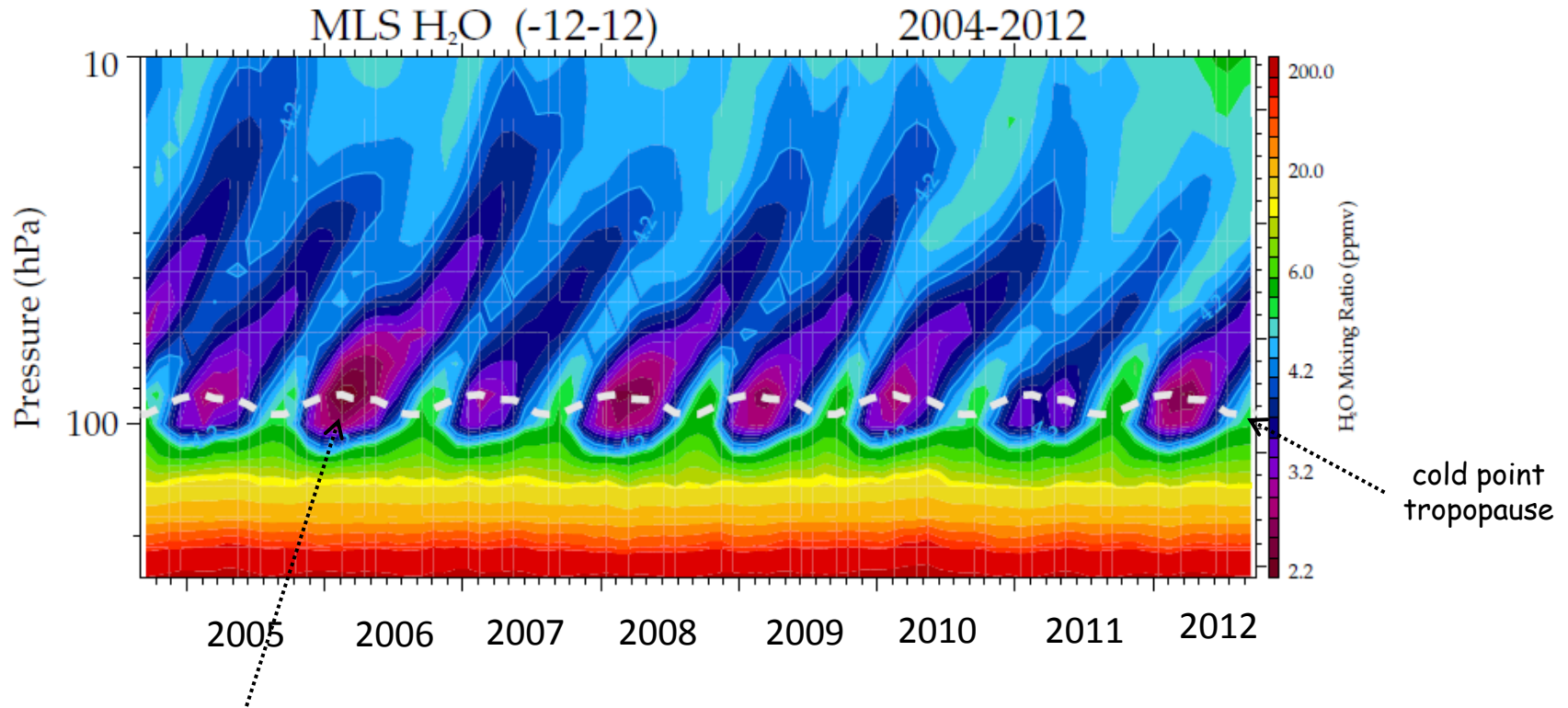
By A. W. BREWER, M.Sc., A.Inst.P.



The stratosphere is extremely dry because air is dehydrated passing the cold tropical tropopause

Isotherms over the Globe
FIG. 5. A supply of dry air is maintained by a slow mean circulation from the equatorial tropopause.

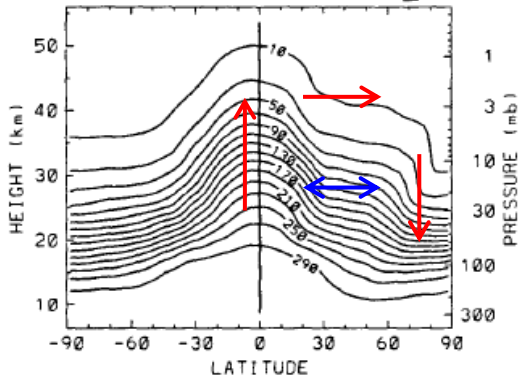
Tropical tape recorder observed by MLS 2004-2012



- annual cycle in tropopause temperature imparts annual cycle in H₂O
- upward propagation with Brewer-Dobson circulation

tracer zonal mean transport budget

model N₂O (ppbv)



$$\frac{\partial \bar{\chi}}{\partial t} = -\overset{\text{mean advection}}{\bar{v}^*} \frac{1}{a} \frac{\partial \bar{\chi}}{\partial \phi} - \overset{\text{mean advection}}{\bar{w}^*} \frac{\partial \bar{\chi}}{\partial z} + \overset{\text{eddy transport}}{\nabla \cdot \mathbf{M}} + P - L$$

$$M_y = -e^{-z/H} \left(\overline{v'\chi'} - \frac{\overline{v'T'}}{S} \bar{\chi}_z \right)$$

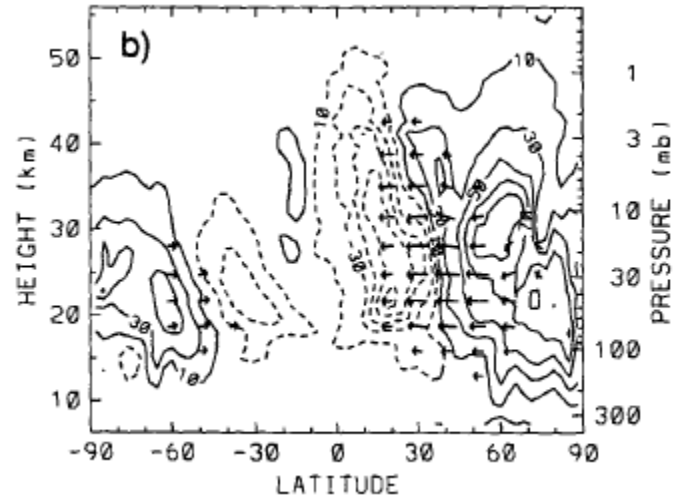
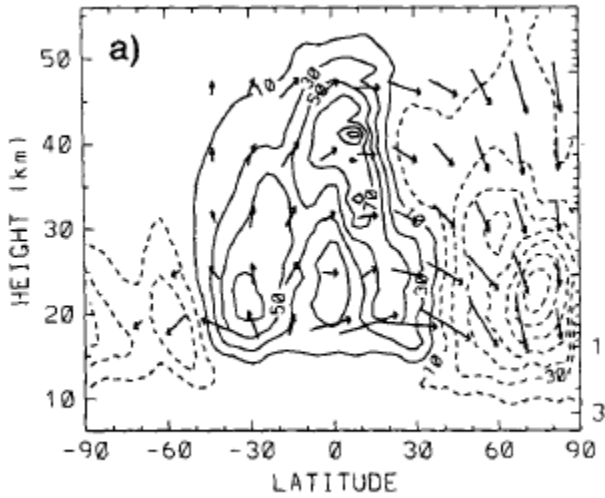
$$M_z = -e^{-z/H} \left(\overline{w'\chi'} + \frac{\overline{v'T'}}{S} \bar{\chi}_y \right)$$

these terms generally balance each other

mean advection

eddy transport

contours:
N₂O
tendency
(ppbv/100 days)



Tracer transport equation similar to thermodynamic equation:

tracer

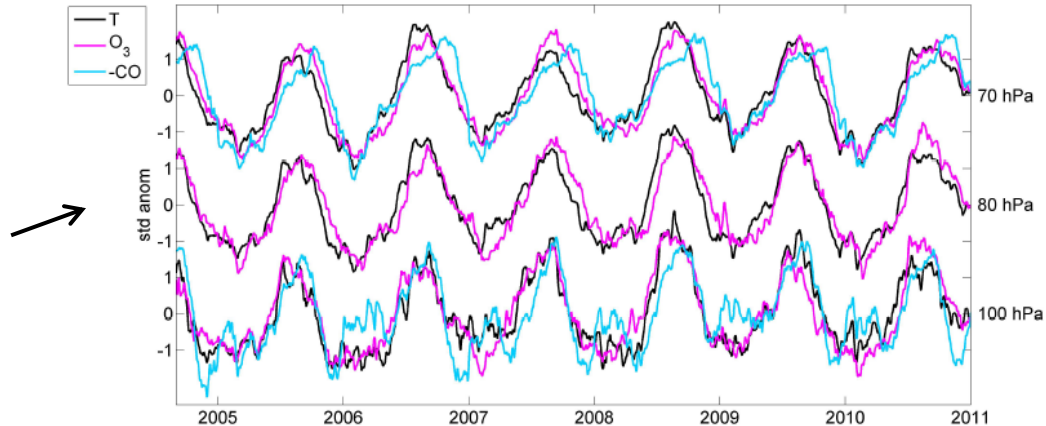
$$\frac{\partial \bar{\chi}}{\partial t} = -\bar{v}^* \frac{1}{a} \frac{\partial \bar{\chi}}{\partial \phi} - \bar{w}^* \frac{\partial \bar{\chi}}{\partial z} + \nabla \cdot \mathbf{M} + P - L$$

temperature

$$\frac{\partial \bar{T}}{\partial t} + \bar{v}^* \frac{1}{a} \frac{\partial \bar{T}}{\partial \phi} + \bar{w}^* S = \bar{Q},$$

This is why temperature and tracers are sometimes highly correlated:

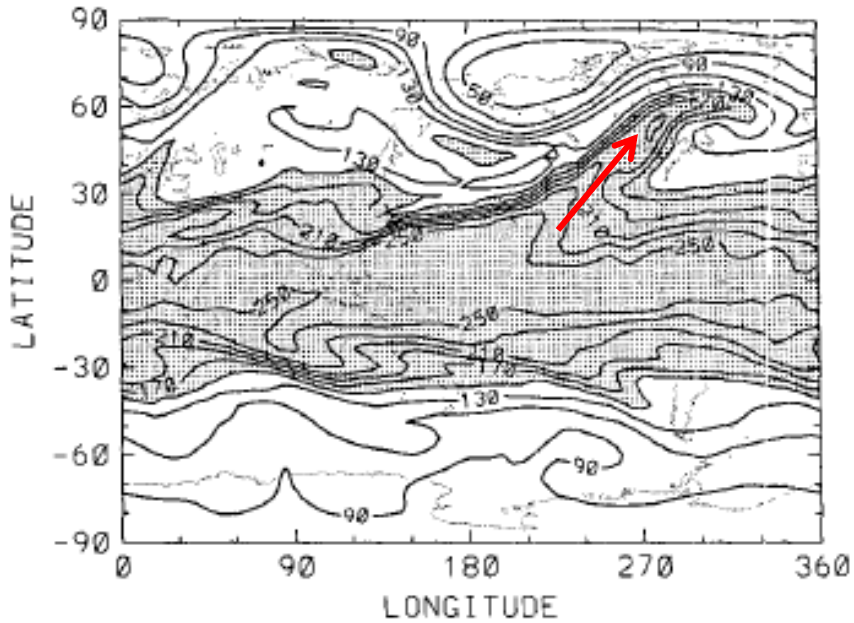
for example,
T, O₃ and CO
in tropical
stratosphere
(Abalos et al 2012)



Examples of stratospheric wave mixing

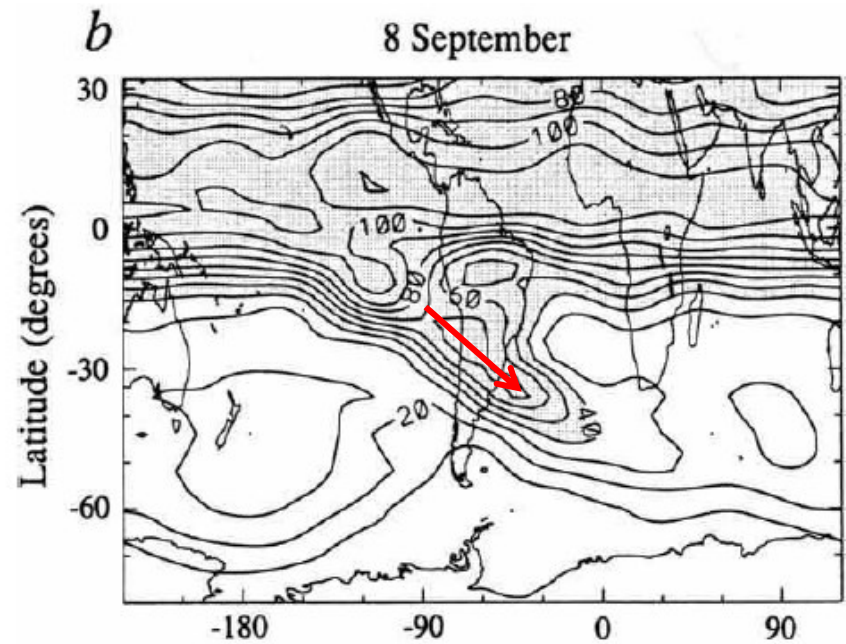
model

CCM2 N₂O February 20 30 mb



N₂O near 35 km from CLAES instrument on UARS

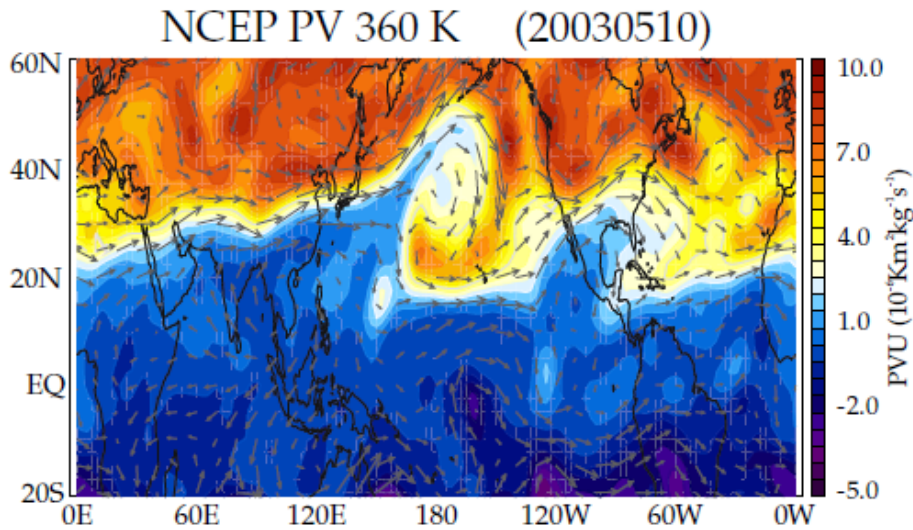
observations



Randel et al 1993

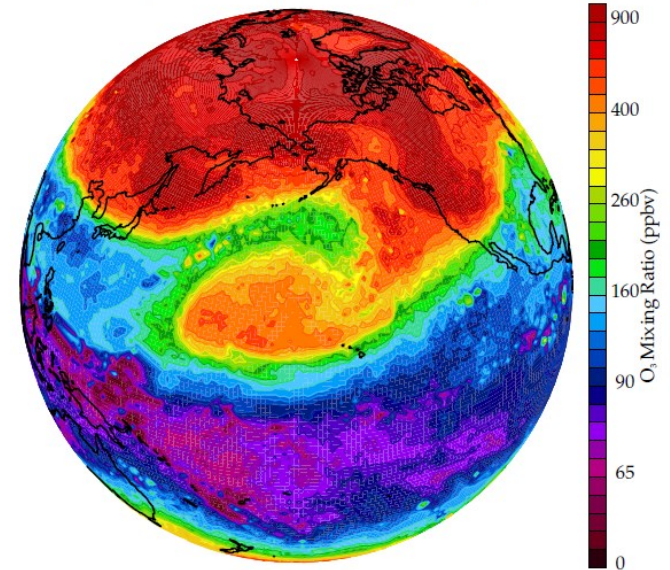
Mixing across tropopause linked to Rossby wave breaking

potential vorticity



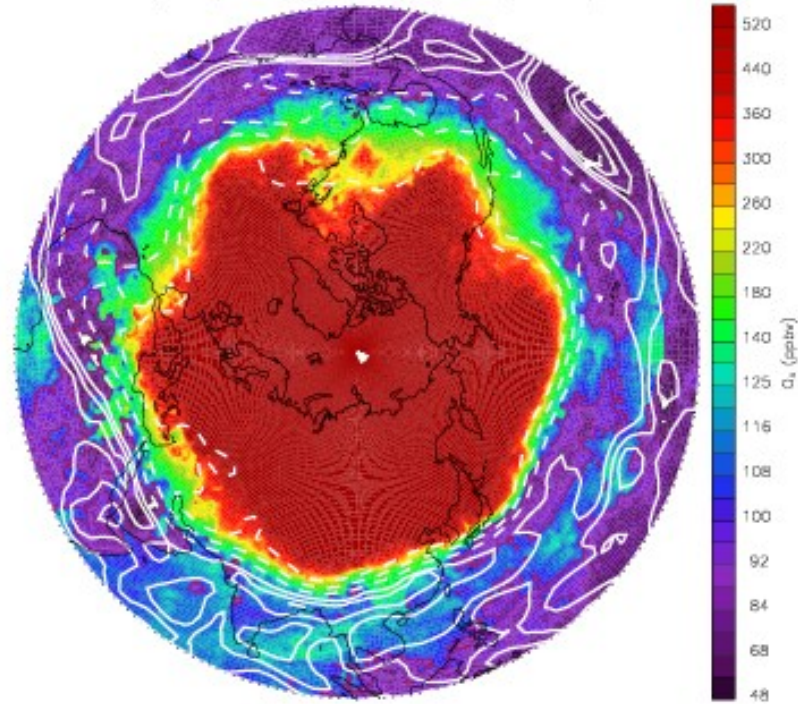
ozone derived from AIRS

AIRS O₃ 360K (MAY/10/2003)



Rossby wave variability reflected in ozone near tropopause

AIRS O₃ (NH) at 360K (JAN/01/2003)



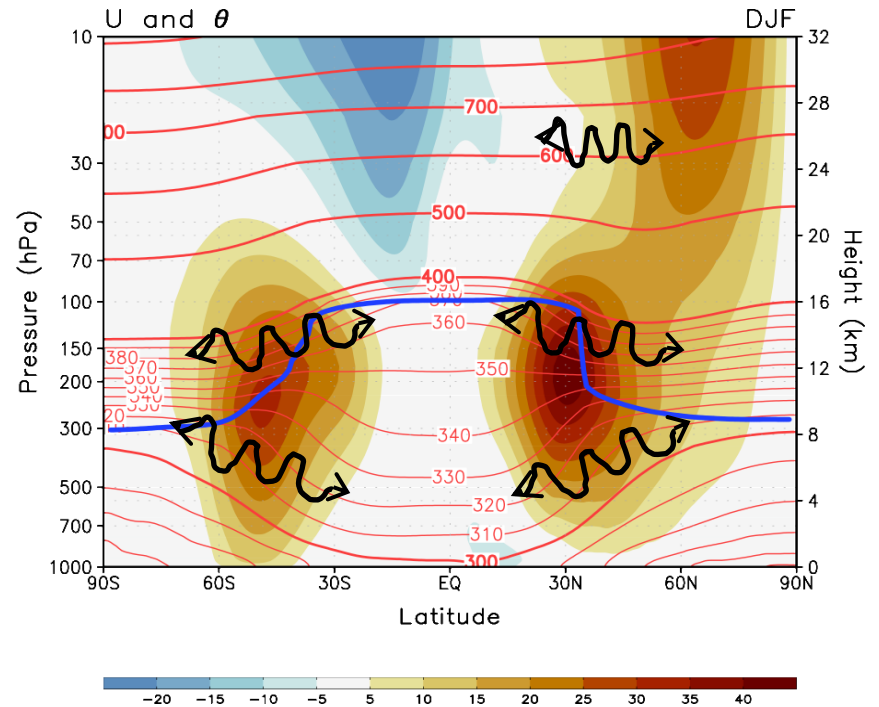
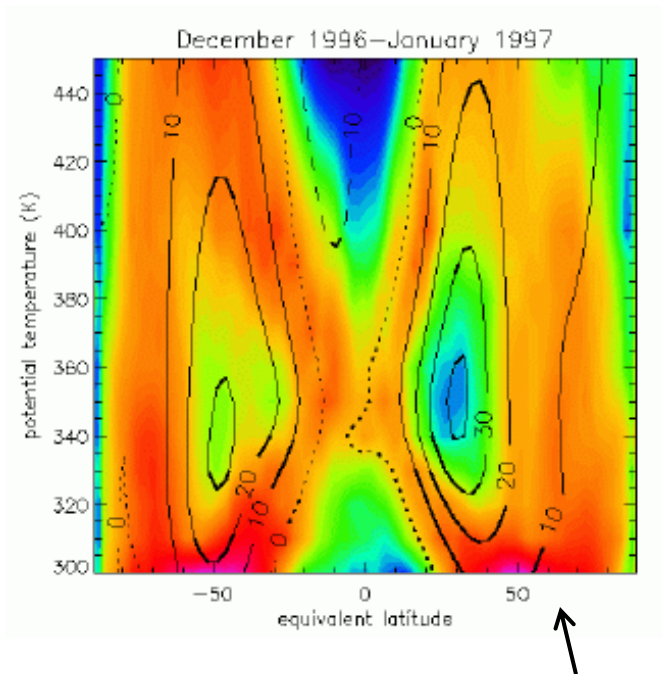
Effective diffusivity as a diagnostic of atmospheric transport

2. Troposphere and lower stratosphere

JGR 2000

Peter Haynes and Emily Shuckburgh

Estimates of mixing based on stretching of PV contours in trajectory calculations



important points:

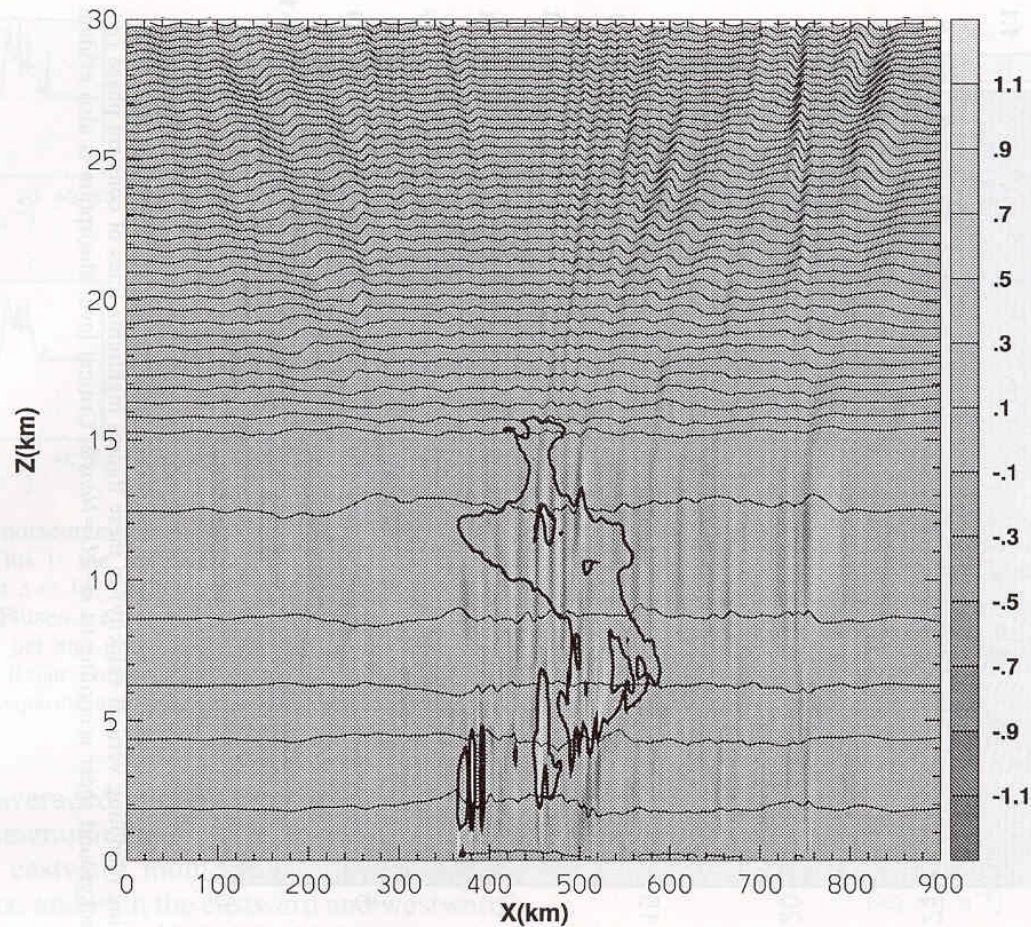
- mixing on flanks of jet (near critical lines for $c \sim 10$ m/s)
- small mixing across jet core (jet cores are mixing barriers)

Key points:

- Stratospheric transport: Brewer-Dobson circulation and wave mixing (clear behavior for tropospheric source gases)
- Stratospheric ozone: produced in tropical stratosphere, transported to high latitudes (reflects seasonal Brewer-Dobson circulation)
- Stratospheric water vapor: dehydration near tropical cold point, strong seasonal cycle ('tape recorder')
- Tracer budgets: mean advection and eddy transports (tied to Rossby waves and critical layers)

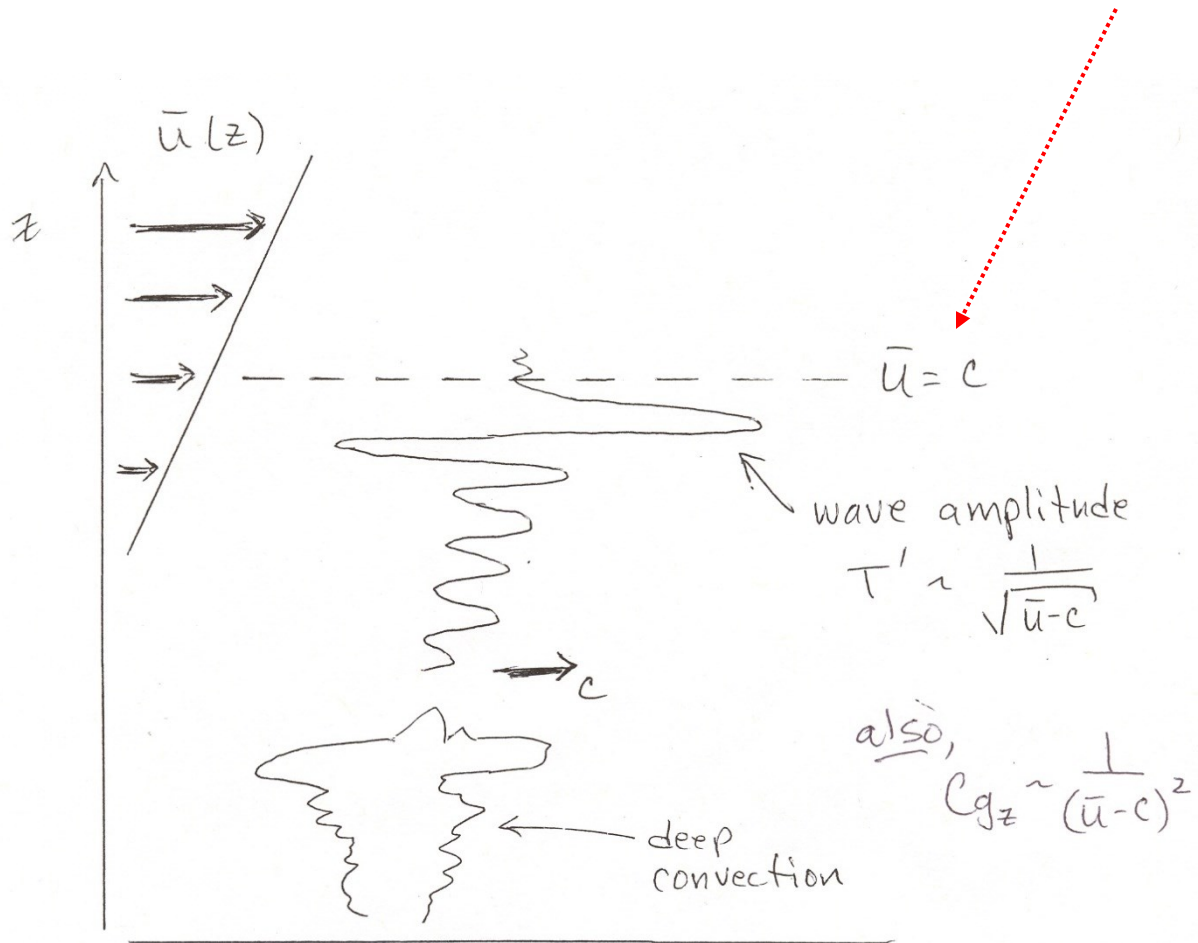
Extra slides

Model simulation of gravity waves forced by deep convection

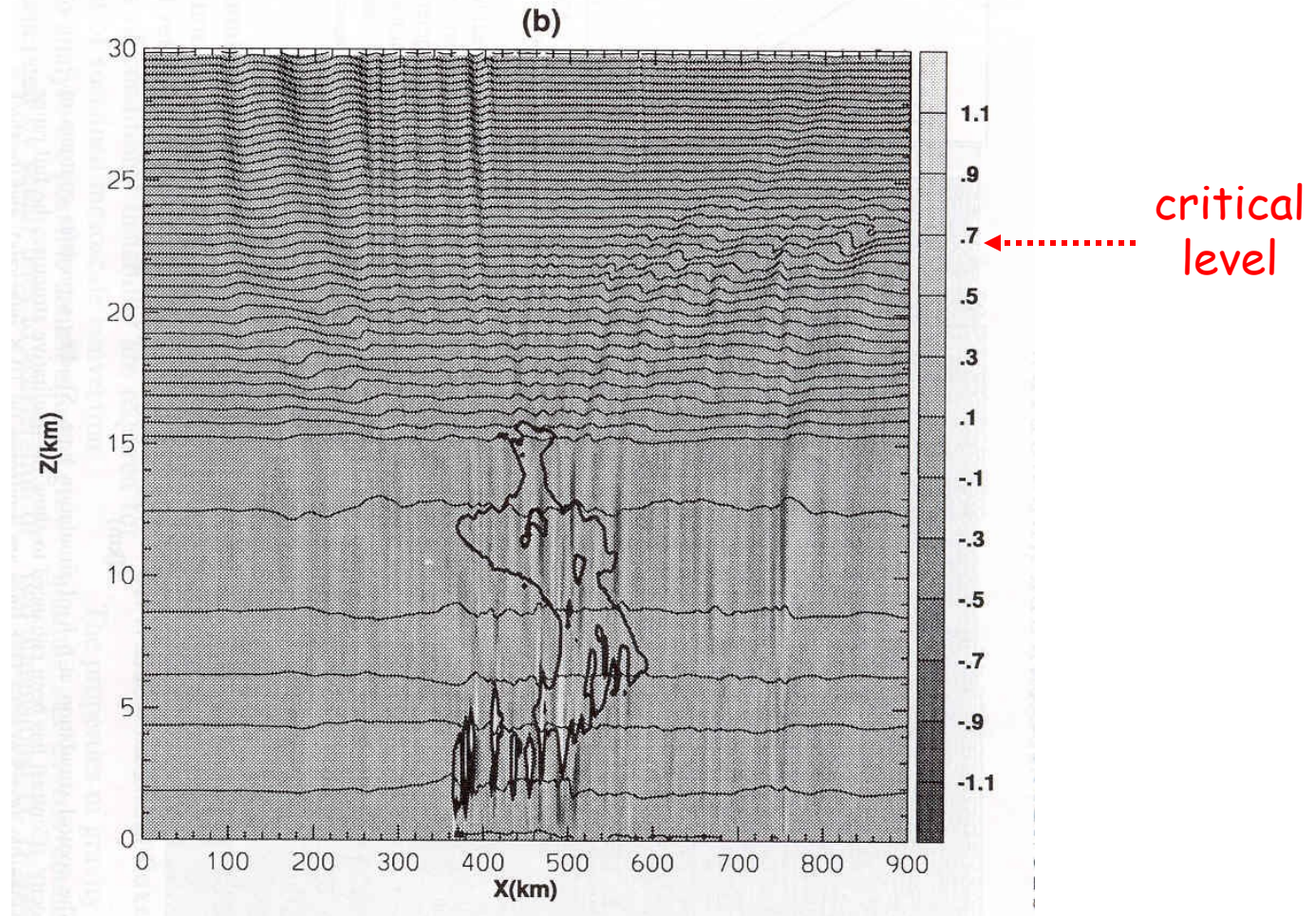


Alexander and Holton, 2000

Gravity waves interacting with a **critical level**



Gravity waves interacting with a critical level



Climatology of Intrusions into the Tropical Upper Troposphere

Darryn W. Waugh

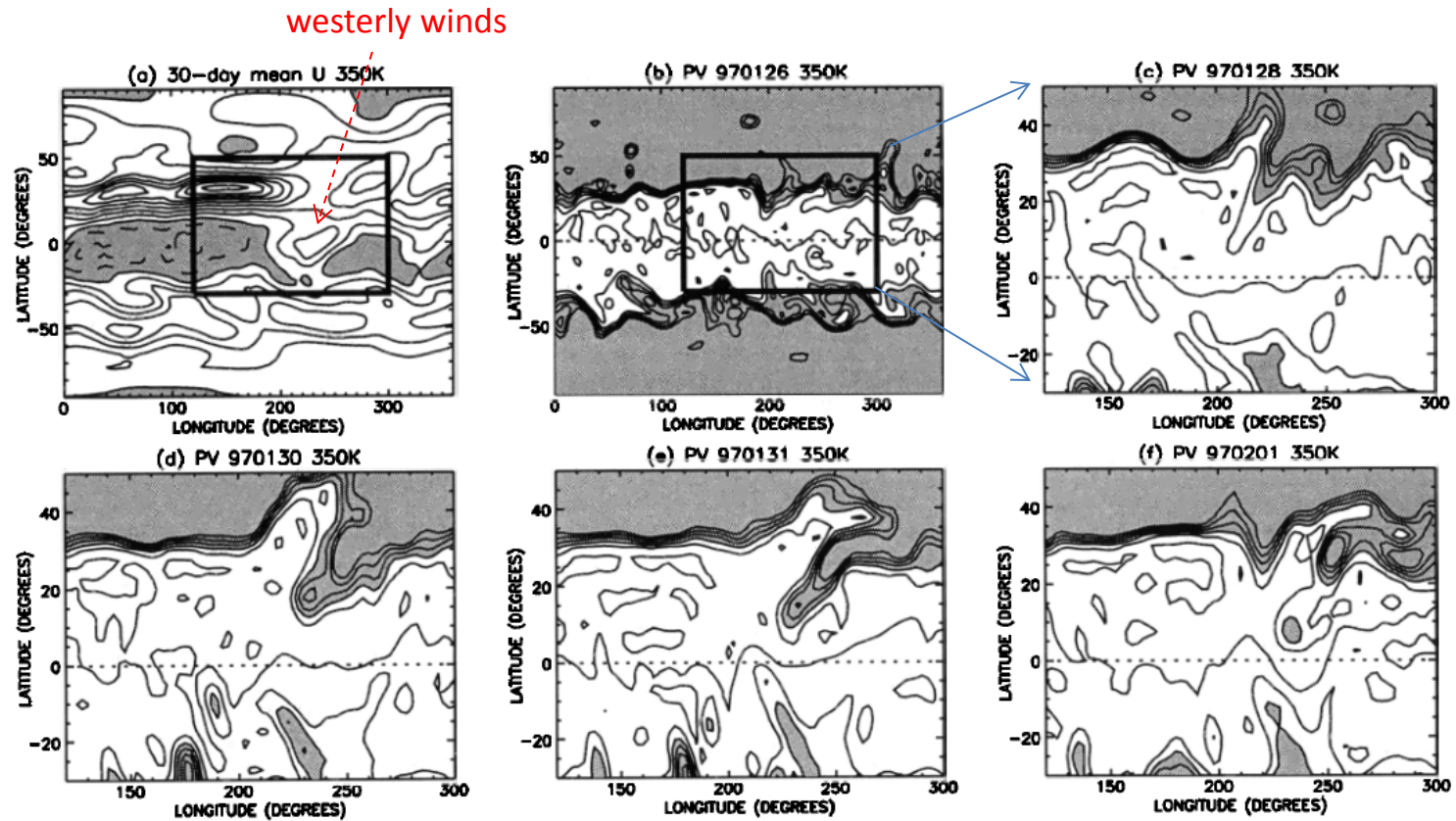
Department of Earth and Planetary Science, Johns Hopkins University, Baltimore, MD.

Lorenzo M. Polvani

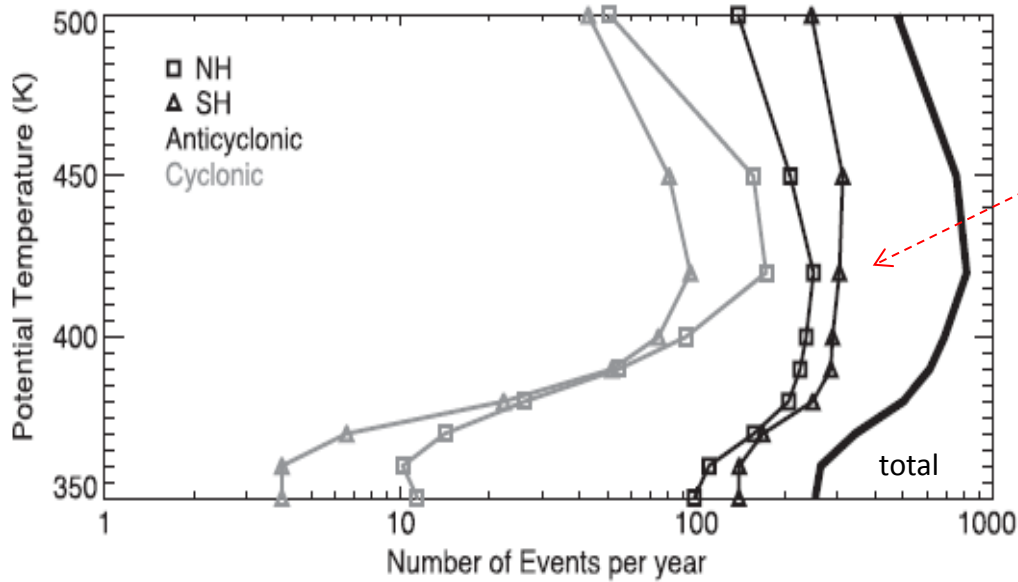
Department of Applied Physics and Applied Mathematics, Columbia University, New York, NY.

Latitudinal propagation depends on background zonal winds:

- Rossby wave propagation through westerly wind 'ducts'



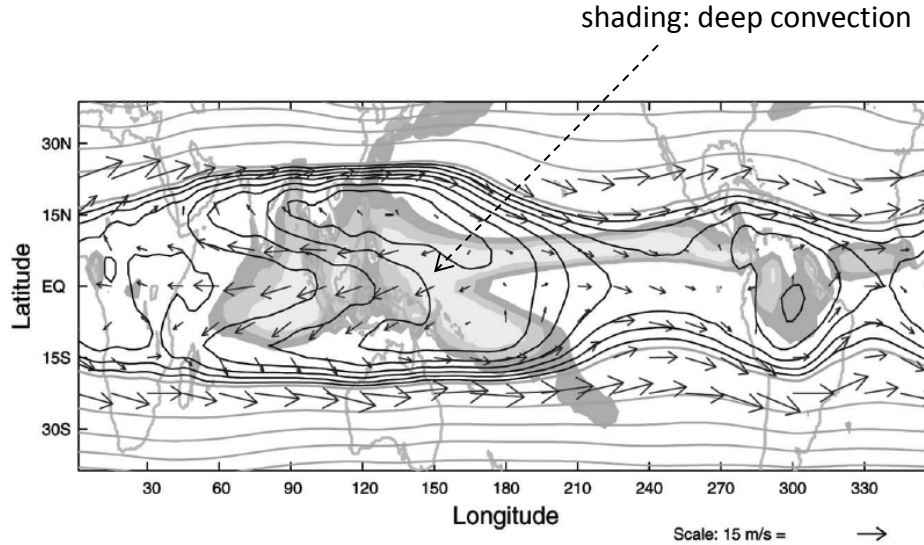
climatology of Rossby wave breaking events



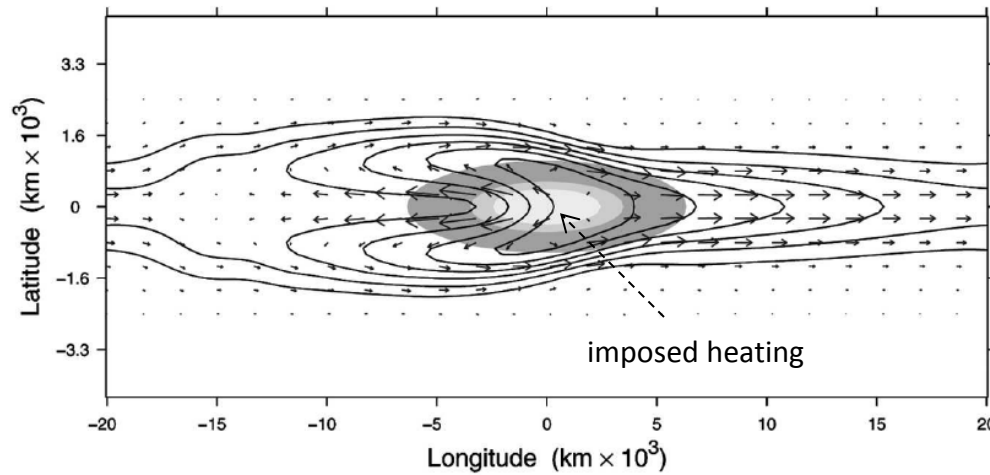
Homeyer and Bowman 2013

observations vs. model

upper troposphere
geopotential height
and winds



annual mean
observations



nonlinear
shallow water
model

Dima et al 2005

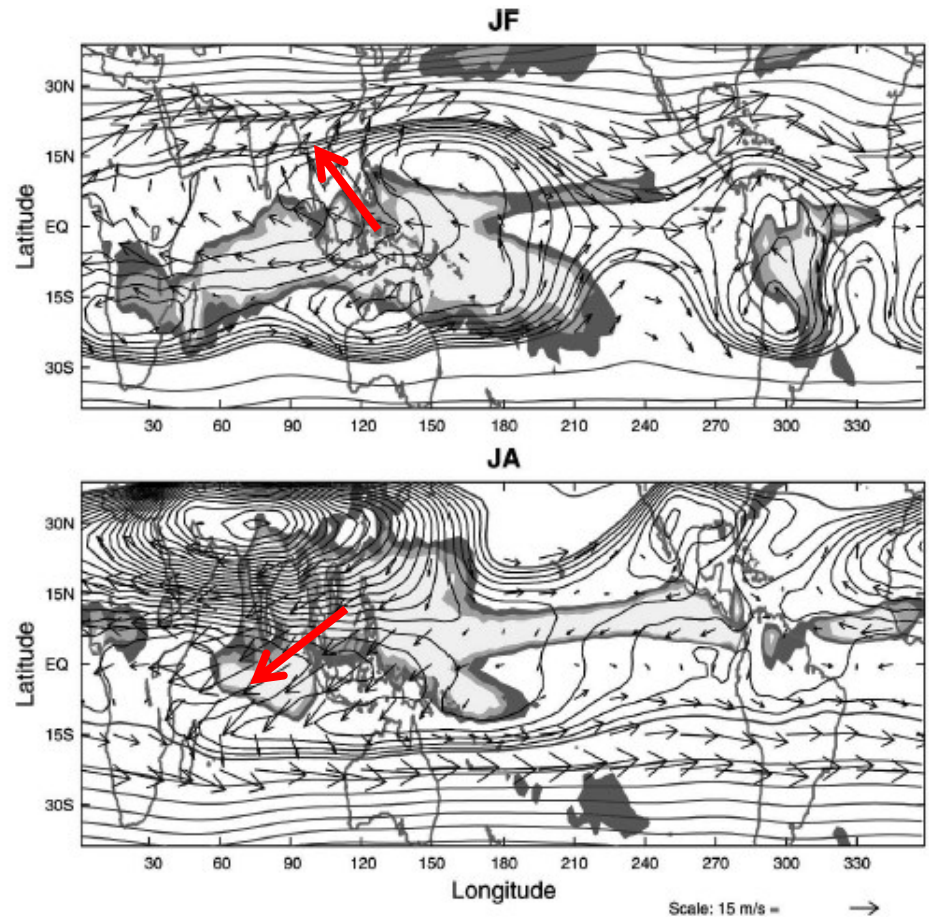
Seasonal variation of tropical waves

tropical zonal mean
momentum balance:

$$\frac{\partial[u]}{\partial t} \equiv [v] \left(f - \frac{1}{\cos\phi} \frac{\partial[u] \cos\phi}{\partial y} \right) - [\omega] \frac{\partial[u]}{\partial p}$$

$$- \frac{1}{\cos^2\phi} \frac{\partial[u^*v^*] \cos^2\phi}{\partial y} - \frac{\partial[u^*\omega^*]}{\partial p} - [F_x].$$

eddy momentum
fluxes balance
Hadley flow

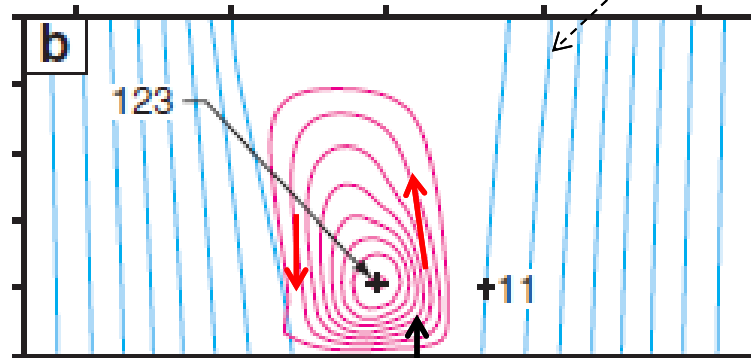


Dima et al 2005

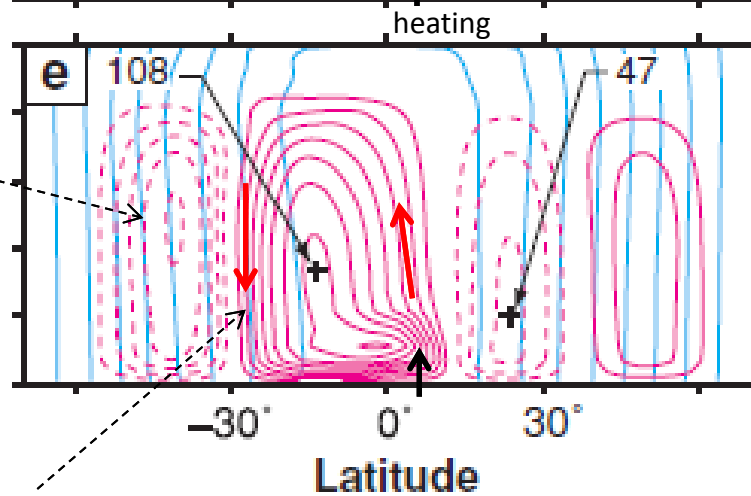
Hadley cell interactions with extratropical eddies: (a complicated subject)

Idealized model simulations:

angular momentum contours



zonally symmetric dynamics



include extratropical eddies

midlatitude
'Ferrel cell' linked
to midlatitude eddies

extension of Hadley
cell to higher latitudes

Schneider 2006

tracer zonal mean transport budget

eddy transport

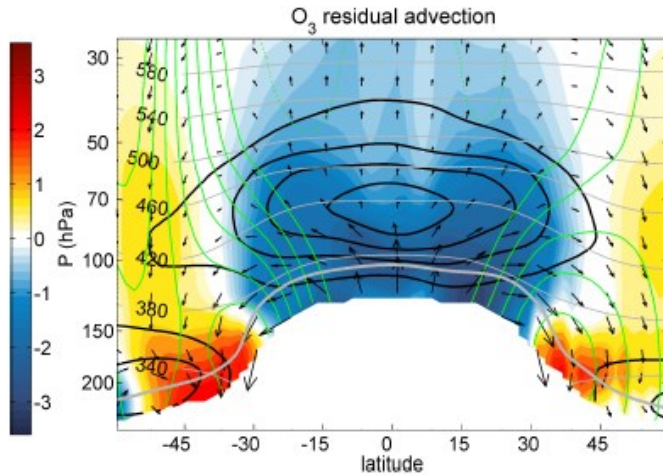
$$\frac{\partial \bar{\chi}}{\partial t} = -\bar{v}^* \frac{1}{a} \frac{\partial \bar{\chi}}{\partial \phi} - \bar{w}^* \frac{\partial \bar{\chi}}{\partial z} + \nabla \cdot \mathbf{M} + P - L$$

tropical ozone budget
from WACCM

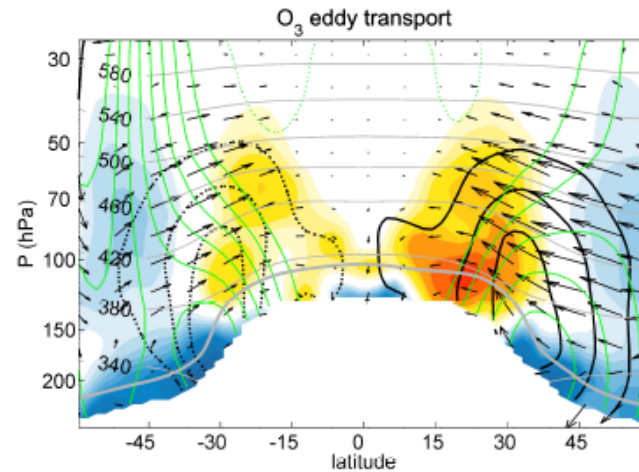
$$M_y = -e^{-z/H} \left(\overline{v' \chi'} - \frac{\overline{v' T'}}{S} \bar{\chi}_z \right)$$

$$M_z = -e^{-z/H} \left(\overline{w' \chi'} + \frac{\overline{v' T'}}{S} \bar{\chi}_y \right)$$

mean advection



eddy transport



ozone
tendency
in % per day

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